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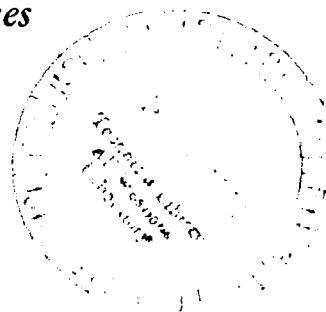
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**ANALYSIS AND TESTING OF TWO-DIMENSIONAL  
VENTED COANDA EJECTORS WITH ASYMMETRIC  
VARIABLE AREA MIXING SECTIONS**

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16. Abstract <p>The analysis of asymmetric, curved (Coanda) ejector flow has been completed using a finite difference technique and a quasi-orthogonal streamline coordinate system. The boundary-layer-type jet mixing analysis accounts for the effect of streamline curvature in pressure gradients normal to the streamlines and on eddy viscosities. The analysis assumed perfect gases, free of pressure discontinuities and flow separation. The analysis treats the three compound flows of supersonic and subsonic streams, those are: (1) primary flow of the driving nozzle, (2) secondary flow between the primary nozzle and the Coanda surface, (3) tertiary flow between the primary nozzle and the other surface of the mixing section.</p> <p>A test program was completed to measure flow parameters and ejector performance in a vented Coanda flow geometry for the verification of the computer analysis. A primary converging nozzle with a discharge geometry of 0.003175 m x 0.2032 m was supplied with 0.283 m<sup>3</sup>/sec of air at about 241.3 kPa absolute stagnation pressure and 82°C stagnation temperature.</p> <p>One mixing section geometry was used with a 0.127 m constant radius Coanda surface. Eight tests were run at spacings between the Coanda surface and primary nozzle 0.01915 m and 0.318 m and at three angles of Coanda turning: 22.5°, 45.0°, and 75.0°.</p> <p>The wall static pressures, the locii of maximum stagnation pressures, and the stagnation pressure profiles agree well between analytical and experimental results.</p>					
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## SUMMARY

The analysis of asymmetric, curved (Coanda) ejector flow has been completed using a finite difference technique and a quasi-orthogonal streamline coordinate system. The boundary-layer-type jet mixing analysis accounts for the effect of streamline curvature in pressure gradients normal to the streamlines and on eddy viscosities. The analysis assumed perfect gases, free of pressure discontinuities and flow separation. The analysis treats the three compound flows of supersonic and subsonic streams, those are: (1) primary flow of the driving nozzle, (2) secondary flow between the primary nozzle and the Coanda surface, (3) tertiary flow between the primary nozzle and the other surface of the mixing section.

A test program was completed to measure flow parameters and ejector performance in a vented Coanda flow geometry for the verification of the computer analysis. A primary converging nozzle with a discharge geometry of 0.003175 m x 0.2032 m was supplied with 0.283 m<sup>3</sup>/sec of air at about 241.3 kPa absolute stagnation pressure and 82°C stagnation temperature.

One mixing section geometry was used with a 0.127 m constant radius Coanda surface. Eight tests were run at spacings between the Coanda surface and primary nozzle 0.01915 m and 0.318 m and at three angles of Coanda turning: 22.5°, 45.0°, and 75.0°.

The wall static pressures, the loci of maximum stagnation pressures, and the stagnation pressure profiles agree well between analytical and experimental results.

### Key Words:

Ejector  
Coanda  
Compressible Flow  
Analysis  
Finite Difference

Computer Program  
Experimental

## Section 1

### INTRODUCTION

#### 1.1 Background

The augmentor wing concept under investigation by NASA for STOL aircraft lift augmentation is powered by an air to air ejector. The wing boundary layer is drawn into the deflected double flap augmentor channel at the trailing edge of the wing and is pressurized by a high velocity slot jet which is oriented at an angle to the augmentor channel. To predict the performance and to optimize the design of the complete augmentor wing, an analytical method is needed to predict the performance of the air ejector which powers the augmentor flap section.

Under contract NAS2-5845 a computer analysis was developed for single nozzle axisymmetric ejectors with variable area mixing sections using integral techniques, Reference 1. The ejectors of primary interest in that program and earlier programs were high-entrainment devices using small amounts of supersonic primary flow to pump large amounts of low-pressure secondary flow. Good agreement was achieved between analytical and experimental results.

The integral analytical techniques used to analyze the axisymmetric ejector configurations are also valid for the analysis of two-dimensional ejectors. However, the augmentor wing configuration may include asymmetric geometries, inlet flow distortions, wall slots, and primary nozzles that are at large angles to the axis of the augmentor mixing section. The integral techniques are not easily adaptable to these more complex flows. Finite difference techniques can be used to analyze these more complex flow geometries at the expense of increased computer time.

Under contract NAS2-6660 a computer program (Reference 2) was developed for two-dimensional, symmetrical mixing sections using finite difference technique and rectangular coordinate system. When Coanda effect is used in two-dimensional ejectors the geometry is not symmetrical and the use of rectilinear coordinates becomes difficult. Flow computations have to account

for the pressure gradient in the direction normal to the streamlines.

## 1.2 Objectives of the Program

The following objectives were defined for this investigation:

1. Develop a computer program for two-dimensional vented Coanda ejectors with non-symmetric variable area mixing sections, with variable Coanda turning, and with variable primary nozzle spacing.
2. Obtain test results with a variable nozzle position vented Coanda ejector configuration for the development and checking of the computer program.

## Section 2

### NOMENCLATURE

$A_N$	Nozzle discharge area ( $m^2$ ; $in^2$ )
$A_{n-1}$	Coefficient appearing in the finite difference equations (-)
$B_{n-1}$	Coefficient appearing in the finite difference equations (-)
$C_P$	Specific heat at constant pressure (J/kg·K; Btu/lbm°R)
$C_N$	Nozzle discharge coefficient (-)
$C_{n-1}$	Coefficient appearing in the finite difference equations (-)
$D_{n-1}$	Coefficient appearing in the finite difference equations (-)
$d$	Nozzle exit height (m; ft)
$E$	Dimensionless eddy viscosity, $v_T/u_o d$ (-)
$k$	Thermal conductivity of fluid (W/mK; Btu/hrft°F)
$K$	Curvature of streamline, $1/R$ (1/m; 1/ft)
$g_o$	Dimensional constant, (32.2 lbm-ft/lbf-sec <sup>2</sup> )
$\ell$	Prandtl mixing length (m; ft)
$L$	Dimensionless mixing length, $\ell/d$ (-)
$m$	Node points along a streamline (-)
$n$	Streamline coordinate (normal to streamlines) (m; ft)
$n$	Streamline designation (-)
$P_b$	Barometric pressure (kPa; psia; inch H <sub>2</sub> O)

$P$	Static pressure (kPa; psig; inch H <sub>2</sub> O)
$P_{01}$	Reference pressure, primary stagnation pressure (kPa; psia)
$P_{rt}$	Turbulent Prandtl number, $\nu_T/\epsilon_H$ (-)
$P_r$	Prandtl number, $\mu C_p/k$ (-)
$q_{eff}$	Effective heat transfer (between streamlines) (W/m <sup>2</sup> °C; Btu/ft <sup>2</sup> sec°F)
$R$	Radius of curvature (streamline; wall) (1/m; 1/ft; 1/in)
$R_g$	Gas constant (Nm/kg°K; lbf-ft/lbm°R)
$R_i$	Richardson number $(2u/R)/(\partial u/\partial n)$ (-)
$S$	Streamwise coordinate (along streamlines) (m; in)
$t$	Nozzle coanda wall spacing (m; inch)
$T_a$	Atmospheric temperature (°C; °F)
$T$	Fluid temperature (°K; °R)
$T_{max}$	Maximum fluid temperature at an $x$ = constant cross section (°K; °R)
$T_{01}$	Reference temperature, primary stagnation temperature (°K; °R)
$u$	Velocity in $s$ direction (m/s; ft/sec)
$u_o$	Reference velocity, $u_o = \sqrt{R_g T_{01}}$ (m/s; ft/sec)
$u_{2,n}$	Unknown velocity at the $n$ th grid point (m/s; ft/sec)
$U$	Velocity in $x$ direction (m/s; ft/sec)
$U_{CL}$	Centerline velocity (m/s; ft/sec)
$U_{max}$	Maximum fluid velocity (m/s; ft/sec)

$U_{sec}$	Secondary velocity (m/sec; ft/sec)
$W_m$	Mixing section total flow (kg/sec m; lbm/sec in)
$W_n$	Nozzle flow rate (kg/sec-m; lbm/sec in)
$W_s$	Secondary flow rate (kg/sec-m; lbm/sec in)
$W_t$	Tertiary flow rate (kg/sec-m; lbm/sec in)
$x$	Space coordinate in the axial direction (m; in)
$X$	Dimensionless space coordinate in the axial direction, $x/d$ (-)
$\Delta X; dX$	Step size in x- direction (-)
$y$	Space coordinate perpendicular to axial direction (m; in)
$Y$	Dimensionless space coordinate perpendicular to axial direction, $y/d$ (-)
$\Delta y$	Dimensionless distance from wall (-)
$\beta$	Nozzle angle (coanda turning) (degrees)
$\gamma$	Ratio of specific heats, $C_p/C_v$ (-)
$\psi$	Stream function (kg m <sup>2</sup> /sec; lbm/ft sec)
$\rho$	Fluid density (kg/m <sup>3</sup> ; lbm/ft <sup>3</sup> )
$\rho_{01}$	Fluid density evaluated at a reference temperature, $T_{01}$ , and pressure, $P_{01}$ (kg/m <sup>3</sup> ; lbm/ft <sup>3</sup> )
$\mu$	Dynamic viscosity (Ns/m <sup>2</sup> ; lbm/ft sec)
$\tau_{eff}$	Total effective shear stress (Pa; psi)
$\tau_w$	Local wall shear stress (Pa; psi)

$\nu_T$	Turbulent eddy conductivity, $\ell^2 \partial u / \partial n$ ( $m^2/s$ ; $ft^2/sec$ )
$\epsilon_H$	Eddy coefficient of heat transfer ( $m^2/s$ ; $ft^2/sec$ )
$\nu$	Kinematic viscosity at local temperature ( $m^2/s$ ; $ft^2/sec$ )
$\delta$	Local wall boundary layer thickness or jet half width (m; in)
$\Phi$	Dissipation function ( $N/m^2 s$ ; $lbm/ft sec^3$ )



## Section 3

### ANALYSIS OF VENTED COANDA FLOWS IN AUGMENTOR DUCTS

#### 3.1 Introduction

This section presents the analysis of asymmetric curved augmentor flows which are steady two-dimensional and compressible and for which the duct geometry is general. The method extends the work presented in reference (2) but a new analysis has been performed and a new program has been developed. The analysis employs the finite-difference technique for representing the equations of motion for compressible flow. It is essentially a boundary-layer-type jet mixing analysis, written in streamline coordinates for ease of computation of curvature effects. The effects of streamline curvature on pressure gradients normal to streamline and on eddy viscosities are computed. Magnitudes of streamline curvature effects are estimated by using a quasi-orthogonal coordinate system and assumed variation of curvature with distance in the normal coordinate direction.

The analysis treats the mixing of three compressible flows of the same perfect gas under the assumption that initial conditions are known and that pressure discontinuities and flow separation are absent. The nozzle exit flow may be supersonic but it is assumed that expansion or recompression outside the nozzle, if needed, will bring the nozzle stream to the local ambient pressure so that shocks and expansion waves at the nozzle exit plane are avoided. Previous work has shown that augmentor performance is little affected by moderate degrees of departure from conditions of correct nozzle expansion. The flows considered include compound flows of supersonic and subsonic streams; however no provision is made for compound choking which may occur with an appropriate transverse distribution of Mach number. Such a condition is amenable to analytical treatment under simplified circumstances, but has not been encountered in experimental tests carried out so far.

To retain the simplicity and speed of the boundary layer approach to augmentor calculation, while incorporating approximate curvature effects, it is necessary to assume an approximate starting line. In the present work the

starting line is comprised of two circular-arcs (See Fig. 1 ) which are tangent to each other, and perpendicular to the nozzle axis at the nozzle exit plane; one arc is normal to the upper wall and one to the lower wall of the duct. In the absence of detailed experimental information on velocity profiles in the wall boundary layers and in the jet shearing layers at the initial plane, the initialization condition has assumed uniform stagnation pressure in the nozzle flow and, separately, for the secondary and tertiary flows. In addition, for the examples worked out in this report, the secondary and tertiary flows have been assumed to have the same stagnation conditions. In computing initial conditions around the circular arcs, the effect of curvature on normal pressure gradient has been taken into account; the initialization satisfies the continuity equation separately for primary, secondary, and tertiary streams under the constraints of local duct width (along the assumed circular arc starting lines), location of the nozzle centre line, and angle of the jet axis with respect to the coordinate system of the duct walls. This initialization is of course approximate but is reasonable to use in the absence of better information on flow starting conditions. If better information is available the initialization procedure adopted in this analysis may readily be replaced to use more detailed or exact information.

### 3.2 Equations of Motion

The momentum, normal pressure gradient, and energy equations in streamwise coordinates are:

$$\rho u \frac{\partial u}{\partial s} = - \frac{\partial P}{\partial s} + \frac{\partial}{\partial n} (\tau_{eff}) \quad (1)$$

$$\rho u^2 K = \frac{\partial P}{\partial n} \quad K = \frac{1}{R} \quad (2)$$

$$\rho u \frac{\partial (C T)}{\partial s} = u \frac{\partial P}{\partial s} + \frac{\partial}{\partial n} (q_{eff}) + \phi \quad (3)$$

in which

$$\tau_{eff} = (\mu + \rho \nu_T) \frac{\partial u}{\partial n}$$

$$q_{\text{eff}} = (k + \rho C_p \epsilon_H) \frac{\partial T}{\partial n}$$

$$\Phi = (\mu + \rho \nu_T) \left( \frac{\partial u}{\partial n} \right)^2$$

In these equations  $s$  and  $n$  measure distance along and normal to the streamlines, respectively, and  $u$  is the velocity component in the stream direction;  $p$  is the static pressure,  $\rho$  the density and  $T$  the temperature of the fluid,  $\tau_{\text{eff}}$  is the total effective shear stress and  $\nu_T$  the eddy viscosity of the fluid with  $\mu$  being the dynamic viscosity. Correspondingly,  $q_{\text{eff}}$  is the effective heat transfer between streamlines with  $\epsilon_H$  be. eddy coefficient of heat transfer and  $k$  the fluid conductivity. Constant values of laminar and turbulent Prandtl numbers have been assumed in the analysis. The term  $\Phi$  is the dissipation function, included in the energy equation. The first order effects of curvature on static pressure are included through the normal pressure gradient equation (2).

### Stream Function

The stream function  $\Psi$  is defined by

$$\frac{\partial \Psi}{\partial n} = \rho u \quad (4)$$

With (4) equations (1), (2), and (3) become:

$$u \frac{\partial u}{\partial s} = - \frac{1}{\rho} \frac{\partial P}{\partial s} + u \frac{\partial}{\partial \Psi} \left[ \rho u (\mu + \rho \nu_T) \frac{\partial u}{\partial \Psi} \right]$$

$$uK = \frac{\partial P}{\partial \Psi}$$

$$u \frac{\partial (C_p T)}{\partial s} = \frac{u}{\rho} \frac{\partial P}{\partial s} + u \frac{\partial}{\partial \Psi} \left[ \rho u (k + C_p \rho \epsilon_H) \frac{\partial T}{\partial \Psi} \right] + \left( \frac{\mu + \rho \nu_T}{\rho} \right) \left[ \rho u \frac{\partial u}{\partial \Psi} \right]^2$$

### 3.3 Dimensionless Parameters

Each variable in the equations of motion is normalized by use of the following reference variables:

$$u^* = \frac{u}{u_o} \quad P^* = \frac{P}{P_{01}} \quad T^* = \frac{T}{T_{01}} \quad \rho^* = \frac{\rho}{\rho_{01}}$$

$$\begin{aligned}
s^* &= \frac{s}{d} & n^* &= \frac{n}{d} & R^* &= \frac{R}{d} & K^* &= \frac{d}{R} \\
\mu^* &= \frac{\mu}{\rho_{01} u_o d} & E &= \frac{v_T}{u_o d} & P_{rt} &= \frac{v_T}{\epsilon_H} & P_r &= \frac{\mu C}{k} \\
\psi^* &= \frac{\psi}{\rho_{01} u_o d} & \gamma &= \frac{C_p}{C_v}
\end{aligned}$$

in which

$$u_o = \sqrt{R T_{01}}$$

$$\rho_{01} = P_{01} / (R T_{01})$$

$d$  = nozzle exit height (the small dimension)

$P_{01}$  = primary stagnation pressure

$T_{01}$  = primary stagnation temperature

Introducing these dimensionless groups into the equation of motion yields the following results:

$$\begin{aligned}
u^* \frac{\partial u^*}{\partial s^*} &= - \frac{1}{\rho^*} \frac{\partial P^*}{\partial s^*} + u^* \frac{\partial}{\partial \psi^*} \left[ \rho^* u^* (\mu^* + \rho^* E) \frac{\partial u^*}{\partial \psi^*} \right] \\
u^* K^* &= \frac{\partial P}{\partial \psi^*} \\
u^* \frac{\partial T^*}{\partial s^*} &= \left( \frac{\gamma - 1}{\gamma} \right) \frac{u^*}{\rho^*} \frac{\partial P^*}{\partial s^*} + u^* \frac{\partial}{\partial \psi^*} \left[ \rho^* u^* \left( \frac{\mu^*}{P_r} + \frac{E}{P_{rt}} \right) \frac{\partial T^*}{\partial \psi^*} \right] \\
&\quad + \left( \frac{\gamma - 1}{\gamma} \right) \left( \frac{\mu^* + \rho^* E}{\rho^*} \right) \left[ \rho^* u^* \frac{\partial u^*}{\partial \psi^*} \right]^2
\end{aligned}$$

From henceforth we omit the superscript \* for convenience so that the following are the equations of motion in dimensionless form:

$$u \frac{\partial u}{\partial s} = - \frac{1}{\rho} \frac{\partial P}{\partial s} + u \frac{\partial}{\partial \psi} \left[ \rho u (\mu + \rho E) \frac{\partial u}{\partial \psi} \right] \quad (5)$$

$$uK = \frac{\partial P}{\partial \Psi} \quad (6)$$

$$u \frac{\partial T}{\partial s} = \left( \frac{\gamma - 1}{\gamma} \right) \frac{u}{\rho} \frac{\partial P}{\partial s} + u \frac{\partial}{\partial \Psi} \left[ \rho u \left( \frac{\mu}{P_r} + \frac{E}{P_{rt}} \right) \frac{\partial T}{\partial \Psi} \right] \\ + \left( \frac{\gamma - 1}{\gamma} \right) \left( \frac{\mu + \rho E}{\rho} \right) \left[ \rho u \frac{\partial u}{\partial \Psi} \right]^2 \quad (7)$$

### 3.4 Evaluation of the Eddy Viscosity

In general the eddy viscosity is defined by:

$$\nu_T = \ell^2 \frac{\partial u}{\partial n} \quad (8)$$

Writing this in dimensionless form, using the dimensionless parameters specified previously and the stream function, the eddy viscosity expression becomes:

$$E = L^2 \rho u \frac{\partial u}{\partial \Psi} \quad (9)$$

in which

$$L = \frac{\ell}{d} \text{ and } E = \frac{\nu_T}{u_o d}$$

The effect of streamline curvature is taken into account by use of the Richardson Number correction in the following approximate form:

$$L = L_o \exp(-3R_i) \quad R_i > 0 \quad (10)$$

$$L = L_o \left[ 2 - \exp(3R_i) \right] \quad R_i < 0 \quad (11)$$

in which  $L_o$  is the dimensionless mixing length in the absence of stream curvature and  $R_i$  is the Richardson Number defined by:

$$R_i = \frac{2u}{R} / \frac{\partial u}{\partial n}$$

$$\text{or} \quad R_i = \frac{2K}{\rho \frac{\partial u}{\partial \Psi}} \quad (12)$$

For small values of  $|R_i|$  the above dependence of  $L$  on  $R_i$  is approximately in accord with the linear relationship derived by Bradshaw (3). For large values of  $|R_i|$  an empirical correlation is not available, and the exponential relationship has been assumed.

In the absence of curvature the mixing lengths  $L_o$  are defined as follows.

#### Boundary Layer

In the inner part of the layer the Van Driest approximation is used.

$$L_o = 0.41 \Delta y [1 - \exp(-y^+/26)] \quad (13)$$

in which  $\Delta y$  is the dimensionless distance from the wall. The variable

$$y^+ = \frac{\Delta y}{\nu} \sqrt{\frac{\tau_w}{\rho}} \quad (14)$$

is evaluated using

$$\tau_w = \frac{\mu u_2}{\Delta y_2} - \frac{\Delta y_2}{2} \frac{\partial P}{\partial s} \quad (15)$$

in which the subscript 2 denotes the streamline coordinate point closest to the wall.

In the outer part of the layer the mixing length is evaluated by

$$L_o = 0.09 \left( \frac{\delta}{d} \right) \quad (16)$$

in which  $\delta$  is the boundary layer 99% thickness and  $d$  is the nozzle exit width.

In the middle part of the layer the smaller of the values of  $L_o$  provided by Equations (13) and (14) is used.

#### Jet Shear Layers

In the shear layer adjacent to the potential-core zone of the primary jet the mixing length is evaluated from

$$L_o = 0.07 \left( \frac{\Delta}{d} \right) \left( 1 + 0.6 \frac{U_{sec}}{U_{CL}} \right) \quad (17)$$

in which  $\Delta$  is the shear layer width (including the zone between 1% and 99% of the total velocity difference between primary and secondary streams) divided by the nozzle width at its exit plane. The term in parentheses takes approximate account of the effects of the secondary velocity,  $U_{sec}$ , of a co-flowing outer stream, on the mixing of a jet whose centerline velocity is  $U_{CL}$ .

For a "fully-rounded" portion of the jet flowing co-axially with a secondary potential stream, the mixing strength has been calculated from

$$L_o = 0.09 \left( \frac{\Delta}{d} \right) \left( 1 + 0.6 \frac{U_{sec}}{U_{CL}} \right) \quad (18)$$

in which  $\Delta$  is the half-width of the jet (evaluated from centerline to the point at which difference between local and secondary velocity is only 1% of the difference between centerline and secondary velocity), divided by the nozzle width at its exit plane.

#### Developing Pipe Flow Region

For the region downstream at the point where the jet spreads to intersect the edge of the boundary layer the mixing is evaluated, as a first approximation only, from

$$L_o = \frac{w}{d} \left( 0.14 - 0.08 \left[ 1 - \frac{\Delta y}{w} \right]^2 - 0.06 \left[ 1 - \frac{\Delta y}{w} \right]^4 \right) \quad (19)$$

in which  $w$  is half the total width, and  $\Delta y$  the distance from the wall. This formula is due to Nikuradse and is cited by Schlichting (3) for fully developed flow in round tubes. Near the wall the mixing length is evaluated by the Van Driest approximation cited earlier, provided the local mixing length so calculated is less than that given by the Nikuradse formula.

The laminar dynamic viscosity is evaluated from (1):

$$\mu = \frac{\mu_{\text{ref}}}{P_{01} \sqrt{R T_{01}}} d \left[ \frac{T_{01} T}{T_{\text{ref}}} \right]^{1/2} \left[ \frac{T_{\text{ref}} + 198.7^\circ \text{R}}{T_{01} T + 198.7^\circ \text{R}} \right] \quad (20)$$

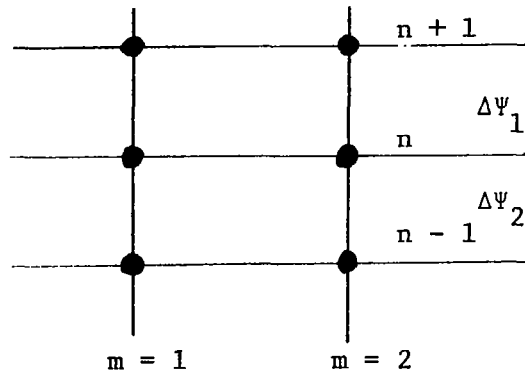
in which  $\mu_{\text{ref}}$  is the laminar dynamic viscosity at a temperature  $T_{\text{ref}}$  and a pressure  $P_{\text{ref}}$  and  $T$  is the dimensionless temperature (normalized by the inlet stagnation temperature of the primary stream  $T_{01}$ ).

The wall shear stress (normalized by the inlet stagnation pressure  $P_{01}$ ) is given by Equation (15). For this equation to be realistic, the grid spacing must be chosen so that  $y^+$  at node point 2 is not greater than 3 - 4. The associated wall friction velocity (normalized by the reference velocity  $u_o$ ) is given by:

$$u^* = \sqrt{\left( \frac{u_2}{\Delta y} - \frac{\Delta y}{2\mu} \frac{dP}{dX} \right) \frac{\mu}{P}} \quad (21)$$

### 3.5 Finite Difference Procedure

By the finite difference technique the derivatives in the differential equations of motion are replaced by differences either along a streamline between two neighboring points  $X$  and  $X + dX$  or normal to it between two neighboring points  $\Psi$  and  $\Psi + d\Psi$ . The finite difference equivalence of the equations of motion are obtained as follows with reference to the following grid lines:





Writing the velocities in terms of a taylor expansion:

$$u_{n+1} = u_n + \left. \frac{\partial u}{\partial \Psi} \right|_n \Delta \Psi_1 + \left. \frac{\partial^2 u}{\partial \Psi^2} \right|_n \frac{\Delta \Psi_1^2}{2!}$$

$$u_{n-1} = u_n - \left. \frac{\partial u}{\partial \Psi} \right|_n \Delta \Psi_2 + \left. \frac{\partial^2 u}{\partial \Psi^2} \right|_n \frac{\Delta \Psi_2^2}{2!}$$

Eliminating  $\left. \frac{\partial^2 u}{\partial \Psi^2} \right|_n$

$$\begin{aligned} \frac{\Delta \Psi_2^2}{2} u_{n+1} - \frac{\Delta \Psi_1^2}{2} u_{n-1} &= \frac{\Delta \Psi_2^2}{2} u_n - \frac{\Delta \Psi_1^2}{2} u_n \\ &+ \left. \frac{\partial u}{\partial \Psi} \right|_n \left[ \Delta \Psi_1 \frac{\Delta \Psi_2^2}{2} + \Delta \Psi_2 \frac{\Delta \Psi_1^2}{2} \right] \end{aligned}$$

Dividing by  $\Delta \Psi_1 \Delta \Psi_2$

$$\frac{\Delta \Psi_2}{\Delta \Psi_1} u_{n+1} - \frac{\Delta \Psi_1}{\Delta \Psi_2} u_{n-1} = \frac{\Delta \Psi_2}{\Delta \Psi_1} u_n - \frac{\Delta \Psi_1}{\Delta \Psi_2} u_n + \left. \frac{\partial u}{\partial \Psi} \right|_n [\Delta \Psi_1 + \Delta \Psi_2]$$

This equation leads to:

$$\left. \frac{\partial u}{\partial \Psi} \right|_n = \frac{\Delta \Psi_2}{\Delta \Psi_1 (\Delta \Psi_1 + \Delta \Psi_2)} (u_{n+1} - u_n) + \frac{\Delta \Psi_1}{\Delta \Psi_2 (\Delta \Psi_1 + \Delta \Psi_2)} (u_n - u_{n-1})$$

$$\text{or } \left. \frac{\partial u}{\partial \Psi} \right|_n = S_5 (u_{n+1} - u_n) + S_4 (u_n - u_{n-1})$$

$$\text{and } \left. \frac{\partial T}{\partial \Psi} \right|_n = S_5 (T_{n+1} - T_n) + S_4 (T_n - T_{n-1})$$

$$\text{in which } S_4 = \frac{\Delta \Psi_1}{\Delta \Psi_2 (\Delta \Psi_1 + \Delta \Psi_2)} \quad \text{and} \quad S_5 = \frac{\Delta \Psi_2}{\Delta \Psi_1 (\Delta \Psi_1 + \Delta \Psi_2)}$$

Using another Taylor Series expansion with a general coefficient:

$$S \left( \frac{\partial u}{\partial \Psi} \right) \Big|_{n+1/2} = S \frac{\partial u}{\partial \Psi} \Big|_n + \frac{\partial}{\partial \Psi} \left( S \frac{\partial u}{\partial \Psi} \right) \Big|_n \frac{\Delta \Psi_1}{2} + \dots$$

$$S \left( \frac{\partial u}{\partial \Psi} \right) \Big|_{n-1/2} = S \frac{\partial u}{\partial \Psi} \Big|_n - \frac{\partial}{\partial \Psi} \left( S \frac{\partial u}{\partial \Psi} \right) \Big|_n \frac{\Delta \Psi_2}{2} + \dots$$

Solving for  $\frac{\partial}{\partial \Psi} \left( S \frac{\partial u}{\partial \Psi} \right) \Big|_n$  we obtain:

$$\begin{aligned} \frac{\partial}{\partial \Psi} \left( S \frac{\partial u}{\partial \Psi} \right) \Big|_n &= \left( \frac{2}{\Delta \Psi_1 + \Delta \Psi_2} \right) \left[ S \left( \frac{\partial u}{\partial \Psi} \right) \Big|_{n+1/2} - S \left( \frac{\partial u}{\partial \Psi} \right) \Big|_{n-1/2} \right] \\ &= \frac{1}{(\Delta \Psi_1 + \Delta \Psi_2)} \left[ (S_{n+1} + S_n) \frac{(u_{n+1} - u_n)}{\Delta \Psi_1} \right. \\ &\quad \left. - (S_n + S_{n-1}) \frac{(u_n - u_{n-1})}{\Delta \Psi_2} \right] \end{aligned}$$

The terms in Equation (5) become:

$$\begin{aligned} u \frac{\partial u}{\partial s} &= u_{1,n} \frac{u_{2,n} - u_{1,n}}{\Delta S_n} \\ - \frac{1}{\rho} \frac{\partial P}{\partial s} &= - \frac{1}{\rho_{1,n}} \left[ \frac{1}{2} \left( \frac{\partial P}{\partial s} \right)_{m=1} + \frac{1}{2} \left( \frac{\partial P}{\partial s} \right)_{m=2} \right] \end{aligned}$$

$$\begin{aligned} \text{and } u \frac{\partial}{\partial \Psi} \left( S \frac{\partial u}{\partial \Psi} \right) &= \frac{u_{1,n}}{\Delta \Psi_1 + \Delta \Psi_2} \left[ \left( \frac{S_{n+1} + S_n}{\Delta \Psi_1} \right) u_{n+1} \right. \\ &\quad \left. - \left( \frac{S_{n+1} + S_n}{\Delta \Psi_1} + \frac{S_n + S_{n-1}}{\Delta \Psi_2} \right) u_n + \left( \frac{S_n + S_{n-1}}{\Delta \Psi_2} \right) u_{n-1} \right] \end{aligned}$$

in which  $S = \rho u(\mu + \rho E)$

With these finite-difference equivalents for the derivative terms Equation (5) may be written in the form:

$$A_{n-1} u_{2,n} + B_{n-1} u_{2,n+1} + C_{n-1} u_{2,n-1} = D_{n-1} \quad (22)$$

in which

$$A_{n-1} = \frac{u_{1,n}}{\Delta S_n} + Y_8 + Y_9 \quad (23)$$

$$B_{n-1} = -Y_8 \quad (24)$$

$$C_{n-1} = -Y_9 \quad (25)$$

$$D_{n-1} = -\frac{1}{\rho_{1,n}} \left( \frac{\partial P}{\partial x} \right)_{m+1/2} + \frac{u_{1,n}^2}{\Delta S_n} \quad (26)$$

$$Y_8 = \frac{u_{1,n} (S_{n+1} + S_n)}{(\Delta \Psi_1 + \Delta \Psi_2) \Delta \Psi_1} \quad (27)$$

$$Y_9 = \frac{u_{1,n} (S_n + S_{n-1})}{(\Delta \Psi_1 + \Delta \Psi_2) \Delta \Psi_2} \quad (28)$$

$$S = \rho u (\mu + \rho E) \quad (29)$$

Similarly, Equation (7) may be written in the form:

$$A_{n-1} T_{2,n} + B_{n-1} T_{2,n-1} + C_{n-1} T_{2,n-1} = D_{n-1} \quad (30)$$

in which

$$A_{n-1} = \frac{u_{1,n}}{\Delta S_n} + Y_8' + Y_9' \quad (31)$$

$$B_{n-1} = -Y_8' \quad (32)$$

$$C_{n-1} = -Y_9' \quad (33)$$

$$\begin{aligned}
D_{n-1} = & \frac{u_{1,n}}{S_n} T_{1,n} + \frac{\gamma-1}{\gamma} \frac{u_{1,n}}{\rho_{1,n}} \left( \frac{\partial P}{\partial x} \right)_{m+1/2} \\
& + \frac{\gamma-1}{\gamma} u_{1,n} S_n \left[ S_5 (u_{1,n+1} - u_{1,n}) \right. \\
& \left. + S_4 (u_{1,n} - u_{1,n-1}) \right]^2
\end{aligned} \tag{34}$$

$$Y_8' = u_{1,n} \frac{S'_{n+1} + S'_n}{\Delta \Psi_1 (\Delta \Psi_1 + \Delta \Psi_2)} \tag{35}$$

$$Y_9' = u_{1,n} \frac{S'_n + S'_{n-1}}{\Delta \Psi_2 (\Delta \Psi_1 + \Delta \Psi_2)} \tag{36}$$

$$S' = \rho u \left( \frac{\mu}{P_r} + \frac{\rho E}{P_{rt}} \right) \tag{37}$$

#### Coefficient Matrix

The general equation

$$A_{n-1} X_n + B_{n-1} X_{n+1} + C_{n-1} X_{n-1} = D_{n-1}$$

with boundary conditions

$$u_1 = 0$$

$$T_1 = T_2$$

$$U_n = 0$$

$$T_n = T_{n-1}$$

can be written for each of the  $n$  grid points with the result:

$$\begin{bmatrix} 1 & -\delta & & & & \\ C_1 & A_1 & B_1 & 0 & 0 & - \\ & C_2 & A_2 & B_2 & 0 & - \\ & & C_{n-3} & A_{n-3} & B_{n-3} & 0 \\ & & & C_{n-2} & A_{n-2} & B_{n-2} \\ & & & & -\delta & 1 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \\ X_3 \\ \vdots \\ X_{n-2} \\ X_{n-1} \\ X_n \end{bmatrix} = \begin{bmatrix} 0 \\ D_1 \\ D_2 \\ \vdots \\ D_{n-3} \\ D_{n-2} \\ 0 \end{bmatrix}$$

where  $\delta = 0$  in momentum equation

$\delta = 1$  in energy equation

The first and second equations are:

$$X_1 - \delta X_2 = 0$$

$$C_1 X_1 + A_1 X_2 + B_1 X_3 = D_1$$

$$A'_1 X_2 + B_1 X_3 = D_1$$

in which

$$A'_1 = C_1 \delta + A_1$$

The last two equations are:

$$C_{n-2} X_{n-2} + A_{n-2} X_{n-1} + B_{n-2} X_n = D_{n-1}$$

$$-\delta X_{n-1} + X_n = 0$$

Combining these

$$C_{n-2} X_{n-2} + A'_{n-2} X_{n-1} = D_{n-1}$$

in which

$$A'_{n-2} = A_{n-2} + B_{n-2} \delta$$

With these two results the order of the matrix can be reduced to  $n - 2$ .

$$\begin{bmatrix} A'_1 & B_1 & 0 & - & - \\ C_2 & A_2 & B_2 & 0 & - \\ & & C_{n-3} & A_{n-3} & B_{n-3} \\ & & & C_{n-2} & A'_{n-2} \end{bmatrix} \begin{bmatrix} X_2 \\ X_3 \\ X_{n-2} \\ X_{n-1} \end{bmatrix} = \begin{bmatrix} D_1 \\ D_2 \\ D_{n-3} \\ D_{n-2} \end{bmatrix}$$

This is solved by the Thomas Algorithm as indicated in References (2) and (5).

### 3.6 Boundary Conditions

The boundary conditions of the walls are:

$$y = y_w(x) \quad (\text{upper and lower walls})$$

$$K = K_w(x) \quad (\text{wall curvature})$$

$$\Psi = \text{const}$$

$$u = 0$$

$$\frac{\partial T}{\partial \Psi} = 0 \quad (\text{adiabatic wall})$$

In the program provision is made for the calculating wall curvatures from wall coordinates  $x, y$  using a least-squares smoothing procedure. Wall coordinates must be specified with sufficient precision to estimate realistic values of wall curvature.

### 3.7 Quasi-Orthogonal Coordinate System

To preserve orthogonality of the coordinate system the step sizes  $\Delta s$  must be adjusted according to:

$$\frac{\partial \Delta s}{\partial n} = \frac{\Delta s}{R}$$

$$\text{or} \quad \frac{\partial \Delta s}{\partial \Psi} = \frac{\Delta s}{\rho u R} \quad (38)$$

Starting with an arbitrarily chosen step size at the nozzle centre line the step size for adjacent streamlines is calculated with the finite-difference equivalent of this formula, using mean values of curvature, density, velocity and step size between the neighboring streamlines.

Streamline curvature ( $K = 1/R$ ) is assumed to decay exponentially with distance from the wall according to:

$$K = K_1 \left( 1 - \frac{\Delta n_1}{w} \right) \exp(-\alpha |K_1| \Delta n_1) + K_2 \frac{\Delta n_2}{w} \exp(-\alpha |K_2| \Delta n_2) \quad (39)$$

in which  $K_1$  (positive) and  $K_2$  (negative) are the curvatures of the lower and upper walls, respectively.  $\Delta n_1$  and  $\Delta n_2$  are the distances along the streamline orthogonal to the lower and upper walls, respectively and  $w$  is the sum of the two. The decay constant  $\alpha$  appears to have a value 2 for potential flow around a cylinder, but in these boundary layer flows a value of 4 seems more appropriate, from comparisons of measured and calculated pressure distributions. The contribution of the first term to streamline curvature in the vicinity of the upper wall is negligible, and vv.

### 3.8 Solution Procedure

The first step in the solution is to determine all flow properties on the initial line. First the radii of curvature of the two circular arcs are determined iteratively to satisfy the conditions of mutual tangency at the nozzle centerline, and orthogonality to the nozzle centerline and to the two walls. The initial stagnation conditions are given and assumed the same for secondary and tertiary flows (See Fig. 1 ). The primary mass flow and the sum of primary and secondary mass flow rates are also given. The initialization then solves iteratively for the split between secondary and tertiary flow rates, and the values of all properties along the initial line under the requirement of mass conservation, and assuming isentropic flow up to the initial line. The solution procedure accommodates substantial effects of streamline curvature, but will not succeed if the total ingested flow rate is so small that the local velocity e.g. at the lower side of the nozzle (Fig. 1 ) is negative. The

initialization procedure also selects the location of streamline node points, with very close spacing required near the walls, and moderately close spacing in the initially-thin shear layers of the jet.

For a set of  $n$  streamlines and known boundary conditions, equations (22) and (30) each provide a set of  $n - 2$  conditions to solve for the unknown velocities and temperatures. Each set of equations can be solved simultaneously if the pressures at downstream node points are known or assumed. For calculation of flow between curved channel walls, the pressure gradient along one streamline is assumed and the downstream pressure on that streamline determined from this gradient and the arbitrary step size. Pressures at corresponding node points on adjacent streamlines are determined by use of the finite-difference form of Equation (6). With the downstream pressures determined, equations (22) and (30) are solved to provide downstream velocities and temperatures, and subsequently all other properties at downstream points. If the calculated value of the outer boundary location does not agree satisfactorily with the actual wall geometry, a new value of the pressure gradient is chosen.



## Section 4

### TEST PROGRAM

A two dimensional experimental rig was designed, fabricated, and installed in our laboratory. The purpose of the experimental work was to obtain test data for verification and adjustment of the computer analysis. The experimental program is described in this section.

#### 4.1 Experimental Apparatus

##### 4.1.1 Two Dimensional Ejector

The two-dimensional ejector consists of a slot type primary nozzle and a two dimensional mixing section. The arrangement of the ejector system for the four nozzle positions tested is shown in Figures 2 and 3 and the positions are listed below.

<u>Test #</u>	<u>Nozzle Angle (<math>\beta</math>)</u>	<u>Spacing (t) (inches)</u>
1 & 2	22.5°	0.80
3 & 4	45.0°	0.80
5 & 6	45.0°	1.31
7 & 8	67.5°	0.70

The nozzle angle ( $\beta$ ) is defined as the angle measured between the vertical and the line running from the center of the coanda arc to the center-line of the nozzle at the throat while the spacing (t) is defined as the perpendicular distance of closest approach between the nozzle and the coanda surface.

A picture of the primary nozzle is shown in Figure 4. The discharge slot is 0.1215"  $\pm$  .0005" by 8.00" with rounded corners. The side walls are made from one quarter inch carbon steel. Four internal supports prevent substantial widening of the discharge slot when the nozzle is pressurized. Dial indicator measurements performed in previous tests revealed that the slot opened up by about 0.0008 inches in the center of the nozzle, about 0.0004" at the

quarter width location, and zero near the ends of the slot. This corresponds to an increase in nozzle slot area of 0.33% when pressurized. Stagnation pressure measurements were made with a kiel probe from side to side at two different axial locations and were found to be uniform across the 8" width of the slot (See Appendix A).

Aluminum pieces on which the nozzle pattern had been cut at the desired angle were bolted to the inside of the side plates to position the nozzle in the inlet to the mixing section (Figs. 6, 7, 8). The nozzle was held firmly in place by threaded rods connecting brackets on each side of the nozzle to brackets attached to busses welded to the side plates of the mixing section (See Figs. 6, 7, 8).

The mixing section, as shown in Figure 2, consists of a rectangular variable area channel formed by two identically contoured aluminum plates, two flat side plates and two dissimilar curved inlet pieces. The upper bellmouth, the coanda surface, is a constant radius ( $r = 5.00$  inch). The nozzle is positioned close to this surface and relies upon the coanda effect to turn the primary flow smoothly into the test section. The lower inlet consists of a bellmouth piece and a straight section. The pictures in Figures 6, 7, and 8 show three views of the mixing section. The two contoured plates were positioned in symmetrical locations about the centerline to form the channel tested (throat height of 1.875"). The width of the mixing section is a constant 8.00 inches along the entire length. The variation of channel height with distance from the nozzle discharge is given in Table 1. Three plexiglass windows were installed along each side of the mixing section so that tufts of wool mounted inside could be observed for indications of flow separations and unsteadiness.

The screened mixing section inlet is shown on Figure 9. Earlier tests (Ref. 2) without the extended inlet showed that highly swirling corner vortices were formed in the four corners of the bellmouth and extended into the test section. The extended inlet eliminated the corner vortices and improved the stability of the ejector flow and static pressures. Four sets of screened inlets were used in the testing to accommodate the four nozzle positions.

#### 4.1.2 Facilities for Ejector Tests

Three subsystems are required for the operation, control and measurement of the air flow through the two dimensional ejector (Figure 9). These three subsystems referred to as the primary flow, the mixed flow, and the boundary layer suction systems are discussed below.

The primary air flow is supplied by a 900 scfm non lubricated screw compressor at 100 psig and an equilibrium operating temperature between 100°F and 140°F. The primary air flow rate and pressure are controlled by an automatic pressure regulator capable of maintaining pressure to within  $\pm .1$  psi of a set value. The mass flow is measured by a standard 3 inch Danial orifice system. A flexible hose connects the primary orifice system to the nozzle.

The mixed flow system consists of a plenum chamber and an 8 inch orifice system. Two different operating flow rates were achieved by the following equipment combinations.

1. Maximum Flow Rate at Atmospheric Discharge - Mixed flow discharges directly into the laboratory.
2. Reduced Flow Rate at Back Pressure - The plenum and orifice are connected to the mixing section discharge.

Mixed orifice flow rates were obtained only for the reduced flow rate conditions. The plenum is shown connected to the mixing section by flexible hose in Figure 10.

The suction system removes the boundary layer flow from each of the four corners of the mixing section to prevent wall boundary layer separation in the ejector. Figure 11 shows six 3/4 inch tubes connected to the top corners of the mixing section. A total of 12 tubes (top and bottom) collect the boundary layer flow from the four corner suction slots which are 0.060 inches wide and are machined into the sides of the contoured plates (See Figs. 12 and 13). The four tubes at one X location are connected to a single large tube under the mounting table (Fig. 14). The three large tubes are each

connected to a large tank plenum through a separate throttle valve. A Roots blower draws the air through the suction system and through a three inch orifice system. The suction system is capable of removing about 1% and 2% of the mixing section flow rate. During the operation of the ejector rig, the boundary layer suction system was used to prevent flow separation in the mixing section diffuser.

The ejector system was operated by starting the primary air flow at low pressure and flow rate. The primary nozzle pressure was increased to 22 psig and the suction was then turned on. Approximately 1/2 hour of warm up time was allowed before testing.

## 4.2 Instrumentation and Data Reduction

### 4.2.1 Instrumentation

The following instrumentation was used in the test facility.

#### Primary Flow System

Flow Rate - Standard 3" orifice system

Nozzle Pressure - Bourdon Pressure Gage accurate to  $\pm .10$  psig

Nozzle Temperature - Copper Constantan thermocouple with digital readout

#### Mixed Flow System

Flow Rate - 8" orifice system for reduced flow rate conditions  
(Tests 2, 3, 6 and 7)

Static Pressures - 77 wall static pressure taps located throughout the mixing section on both top and bottom contoured plates and on both bellmouth pieces (See Figs. 12 and 13). As shown in Figure 10, the static taps were connected to a valving system which permitted easy determination of individual static pressures without the need for a large

manometer bank. Tygon tubing was used to connect the taps to four pressure sampling valves each capable of handling 24 inputs. These pressure valves were in turn connected through a switching network to any of 3 well type manometers which permitted the accurate determination of pressures over the range of + 25 inches of water gage to - 75 inches of water gage.

Traverse Data - Stagnation pressure and temperature profiles were measured at up to 11 axial locations using a 1/8" diameter stem kiel pressure-temperature probe. Both a mercury manometer and a pressure transducer coupled with a direct digital readout were used for pressure measurements. A direct digital readout was used to indicate total temperatures.

#### Suction Flow System

Flow Rate - 3" orifice system

Suction Pressure - Bourdon type pressure gage

#### 4.2.2 Data Reduction Procedures

Three types of data reduction calculations were needed in this program.

1. Standard orifice calculations
2. Velocity profile calculations
3. Integration of velocity profiles to calculate flow rate

All of these calculations were programmed on a time sharing computer. The orifice calculations were programmed as a subroutine to the main data reduction program using standard orifice equations and ASME orifice coefficients. The equations used in the determinations of velocity were included in the main program and are the standard compressible flow relationships which

can be found in most fluid mechanics text books.

The integrated mass flow rate for each traverse was computed by integrating the product of the local velocity and local density over a two dimensional section of unit width. The program also calculated the "mass-momentum" stagnation pressure at each traverse section using the equations presented on pages 52 and 53 of Reference 6. The mass momentum method determines the flow conditions for a uniform velocity profile which has the same integrated values of mass flow rate, momentum, and energy as the non-uniform velocity profile actually present.

#### 4.2.3 Experimental Uncertainty

##### Orifice Calculations

The techniques presented in Reference 7 were applied to the primary flow orifice calculations and the mixed flow orifice calculations. The following uncertainty results were obtained:

<u>Orifice</u>	<u>Pressure (psig)</u>	<u>Uncertainty</u>
Primary Nozzle	22 psig	$\pm 0.8\%$
Mixed	slightly above atmospheric	$\pm 1.3\%$

##### Static Pressure

Uncertainty in the wall static pressures occur mainly because of fluctuations in the manometer liquid columns caused by unsteadiness in the flow. The degree of these fluctuations may therefore be used as an indication of the uncertainty of the pressure readings. For the unrestricted maximum flow rate condition the wall static pressure fluctuation reached a maximum of  $\pm 1.0$  inch of water, while for the reduced flow rate condition the maximum reached only  $\pm 0.4$  inch of water.

##### Integrated Mass Flow Rate

The mass flow rate calculated by integrating the results of the

stagnation pressure and temperature traverses is influenced by many items and is therefore very difficult to estimate. The following items all contribute to the uncertainty in integrated mass flow rate:

1. unsteady wall static pressures
2. unsteady traverse stagnation pressures
3. instrument accuracy of the pressure transducer and digital readout
4. inaccuracies due to the effect of steep velocity gradients on sensed pressure
5. inaccuracies due to probe effect near the mixing section walls
6. inaccuracy in probe position
7. assumptions and inaccuracies associated with the data reduction computer program
8. data recording errors or computer data input errors
9. errors caused by loose connections in the pneumatic sensing tube between the probe and the transducer
10. non-two-dimensional flow distribution across the width of the 8 inch mixing section

All of these effects could combine to give both a  $\pm$  uncertainty band and a fixed error shift.

One measure of the uncertainty due to these effects is obtained from the limits of individual integrated mass flows for each test run. These values are listed on Table 2 for all of the test runs with traverse data. The results presented on Table 2 show an average variation of +4.4 and -3.0 or a total spread of 7.4%. These values only include the effect of variable uncertainty and exclude the uncertainty due to probe errors in steep gradients and near walls and integration assumptions. Both of the excluded errors probably cause the integrated mass flows to be too large because the probe tends to measure too high near the wall and the integration program neglects wall boundary layers.

### 4.3 Test Schedule and Results

A total of eight ejector tests were carried out on one mixing section configuration (1.875" height) and at one nozzle pressure, 22 psig. Four different nozzle positions were tested each at atmospheric discharge and at a back pressure condition. Figure 3 summarizes the conditions for each of the tests. In each test, readings were taken from the wall static pressure taps, orifice system instruments, and from the total pressure and temperature taps on the traversing probe. The number and location of traverses varied from test to test. The traverse locations and a summary of the data presented for each test is given in Table 3.

The data presented in this report falls into the following categories:

Test Conditions and Mass Flows

Static Pressures

Maximum Local Pressures

Velocity Profiles

Richardson Number Coefficient Sensitivity

Streamline Curvature Decay Sensitivity

A summary of the figures and tables used to present data from each test run is presented in Table 3. Discussion of the data will be taken up in the next section.

A sample of the static pressure data as taken is tabulated in Appendix A for Tests 7 and 8.



## Section 5

### COMPARISON OF ANALYTICAL AND TEST RESULTS

#### 5.1 Test Conditions and Mass Flows

Table 5 and Figure 15 show a comparison of

- a) the total flow rate measured by an orifice downstream of the diffuser
- b) total flow rates determined by integration of measured velocity profiles at various sections in the test section
- c) flow rates inferred by use of the computer program with best fit to the static pressure in the throat region

In general the flow rates determined by methods (a) and (c) agreed satisfactorily within approximately 6%. Flow rates estimated by integration of experimental velocity profiles at various stations were consistent with one another within approximately 8% but disagreed with the results of (a) and (c) up to 15%. Previous experience (Ref. 2) also showed that integration of experimental velocity profiles yielded too high a mass flow. In that case the discrepancy was of the order of 6%. The velocities were determined experimentally by the use of a Kiel probe to determine a local stagnation pressure coupled with the assumption that the local static pressure was equal to the wall static pressure. The velocities were determined in this way only for stations downstream of high wall curvature. The disagreement between total flows determined by integrating velocity profiles and those obtained from orifice measurements may be due to the effect of high shear and turbulence level in these flows upon the apparent stagnation pressure reading of the Kiel probe. In view of these discrepancies reliance was placed in these tests upon the mass flows determined by methods (a) and (c).

The accuracy of the primary flow measurements determined by orifice readings for the primary flow are of the order of 3%. The secondary flow determinations by orifice have an apparent uncertainty level of + or - 1.5%.

## 5.2 Wall Static Pressures

Figure 16 shows a comparison between the static pressures predicted using this computer program and those measured with the symmetrical diffuser employed in the NAS-50 program (Reference 2) for which wall curvatures were very small and had little effect upon the axial static pressure profile. This comparison which is included as a check point in this discussion shows that the curvature program predicts the wall static pressures reasonably well when the nozzle is located at the mid-plane of a symmetrical test section.

Figures 17 through 20 show experimental values of the wall static pressures measured for eight experimental cases with the unsymmetrical test section and with the nozzle located at various distances from the curved wall and at various angles with respect to the test section axis. In general there is a considerable difference in static pressure between top and bottom walls with the lowest pressure being near the top wall which had the highest curvature and near which the nozzle was located. These low pressures between the top wall and the nozzle were accompanied by velocities considerably higher than those in the region between the nozzle and the lower wall. Downstream of the region of considerably high wall curvature the measured wall static pressures were nearly the same on top and bottom walls.

Figures 17 through 20 also show the computed static pressure distributions on the top and bottom walls. In general the agreement between analytical measured results is considerably better with high coanda effect, i.e., large turning angle (up to  $67.5^\circ$ ) and small spacing between the nozzle and the wall. The greatest discrepancies between analytical and experimental results are associated with those cases where the calculation method indicates that the flow is on the verge of separation, for example, cases 1 and 2. In cases 5 and 6 the static pressure distribution on the top wall near the nozzle is quite different from the calculated value. Here the spacing between nozzle and wall is large at 1.10 inches. This large venting of the flow between the nozzle and the wall appears to have substantially diminished the Coanda effect lessening the tendency of the flow to cling to the upper wall and increasing the possibility of separation in the region immediately

downstream of the nozzle. In the region of separation the flow calculation becomes somewhat uncertain and experimental details were insufficient to ascertain whether the flow were actually separated in that region. However the degree of agreement between measured and calculated results was substantially poorer for cases 5 and 6 with large venting between the nozzle and the wall.

In general the reasons for differences between the analytical and calculated results are as follows:

1. Flow Separation - The experimental results for cases 1 and 2 and case 6 appear to be on the verge if not actually past the margin of flow separation, as indicated by the calculated values of the wall shear stress or the velocities near the wall for those cases.
2. Initialization Approximations - As explained in the earlier section of the report, the initialization process assumes a circular arc starting line for each of the two regions between the nozzle and the upper and lower walls. The initialization process is assumed to have isentropic flow up to the starting line where the nozzle has zero thickness, so an approximation is used for estimating the decay of streamline curvature with distance away from the wall. In the initialization process it was found that the calculated wall static pressure distribution immediately downstream of the nozzle was very sensitive to the ratio of the flows between nozzle and upper and lower walls respectively. This flow ratio was not available experimentally and was determined in the initialization process by requiring smooth continuity of static pressure along the circular arc starting lines. The initialization process was also sensitively dependent upon decay of wall curvature away from the upper and lower walls, especially in the jet zone.
3. The Use of a Quasi-Orthogonal Coordinate System - In the calculation method curvatures were estimated by the use of

equation 39. In principle this estimation could have been used for a first approximation to determine the entire velocity field then subsequent iterations could have utilized the calculated velocities to determine a second approximation for streamline curvature. However this approach was considered excessively time-consuming for the present problem and not sufficiently justified by the requirements of computing vented Coanda flows where the nozzle spacing is not large and the wall curvature is substantial. The major curvature effects are experienced in a region close to the wall itself.

4. The Effect of Streamline Curvature on Jet Mixing Turbulent Shear Stresses - As pointed out earlier, the first-order effects of curvature on turbulent mixing length have been estimated by Bradshaw. An approximate expression for the dependence of mixing length upon Richardson number is included in the calculation method. However this approximation is not well validated by experimental data and includes curvature effects significantly larger than those considered by Bradshaw, hence, this adds another element of uncertainty to the flow calculation.

Figure 21 shows the effect of varying the Richardson number coefficient (from 3 to 10) upon a computed static pressure distribution along the wall for case 6. This range of Richardson coefficient may be thought to represent the uncertainty in the magnitude of the effect but suggests very little alteration on computed wall static pressure distributions. Increasing the value of the Richardson number coefficient tends to decrease the turbulent shear stresses in the upper part of the jet mixing zone and to increase them on the lower side.

Figure 22 shows the sensitivity of the calculation of wall static pressures for case 6 upon the assumed value of the streamline curvature decay coefficient used in equation 39. The effects of variation of this coefficient are naturally unimportant in the downstream region where curvatures are small

but can be quite significant in the region for  $\alpha$  less than 0, i.e., close to the zones of high wall curvature.

### 5.3 Locus of Maximum Stagnation Pressure

Figures 23 through 30 show the computed location of the line of maximum stagnation pressure from the nozzle to a point far downstream in the test section. Also shown are test data points taken from the maximum stagnation pressure in the upstream zone and from the maximum velocity, i.e., maximum stagnation pressure for the downstream region in which curvature effects are negligible.

As with the static pressure comparisons the best agreement between calculated and measured values of the locations of maximum stagnation pressure correspond to those experimental cases in which there was the largest degree of Coanda turning and the smallest spacing between the nozzle and the wall. This is shown particularly by the comparison for cases 7 and 8. In other cases, for example, cases 5 and 6 (with large venting between the large spacing between the nozzle and the wall) the computed location of maximum stagnation pressure shows a substantial deviation from the experimental results. These two cases as pointed out earlier appear to show a substantially diminished Coanda effect. The jet clearly does not cling as closely to the wall as the computer model predicts. In general the differences between computed and experimental results may be ascribed to the reasons mentioned earlier for the static pressure discrepancies.

Figures 31 and 32 show the total pressure profiles across the mixing section. The analytical prediction shown with continuous lines shows a small deviation from the experimental results in the tertiary flow that is near the bottom wall.

### 5.4 Velocity Profiles

Figures 33 through 40 show comparisons of non-dimensional velocity profiles at various locations throughout the test section for each of the eight cases investigated. Owing to the uncertainty in velocity determination as

evidenced by the integrated mass flow of discrepancy being up to 15% the large uncertainty level must be attached to each velocity determination. Hence, the rather substantial discrepancies between experimental and computed velocities are not conclusive indications of the degree of reliability of the analytical method. The experimental velocity profiles for Runs 5 and 6 show that the maximum velocity region has been shifted towards the bottom wall indicating a reduced coanda effect at large nozzle spacing.

## Section 6

### GENERAL CONCLUSIONS

1. An approximate method has been developed for calculation of vented Coanda flows in ducts. A method has been confirmed by experimental data with angles of turning up to  $67.5^\circ$  and for close spacing between the nozzle and the curved wall. At larger spacing the model indicated flow separation which limits the availability of the model to represent the flow profile in the downstream zone.
2. A quasi-orthogonal method of computation, which is more rapid than an iterative solution of the elliptic boundary value problem, appears best suited to ducted Coanda flows with low venting and large curvature. It requires approximate specification of streamline curvature decay with distance from the wall, and thus is best suited to cases in which the jet sheet is located close to the wall. It is desirable to extend the use of the method to non-vented Coanda flows.
3. Though the effects of streamline curvature on mixing length are known only for small curvature, and perhaps uncertain within a factor of 3, a simple correction for mixing length in terms of Richardson number appears to provide a reasonable estimate for jet curvatures  $d/R$  of the order of 0.02.
4. The flow model developed provides good agreement with secondary and primary mass flows measured with orifice plates. Integration of velocity profiles failed to provide satisfactory agreement with orifice measurements of mass flow apparently due to the effects of a high turbulence and high shear in the mixing zone on the stagnation pressure readings of a Kiel probe.
5. The flow model predictions were in good agreement with measured wall static pressures except in the immediate region of the nozzle apparently due to upstream boundary layer and nozzle thickness effects, and due to incipient flow separation in certain of the tests.

6. A sensitive measure of the degree of agreement between flow model and experimental results is the location of the maximum stagnation pressure line in these highly curved flows.



# APPENDIX A

## Sample Tabulation of Static Pressures (Tests 7 and 8)

Tap Position inch	Wall Static Pressure - inches of water gage			
	Run 7		Run 8	
	Top Wall	Bottom Wall	Top Wall	Bottom Wall
- 5.50	-	- 1.15	-	- 2.10
- 5.00	-	- 1.35	-	- 2.60
- 4.50	-11.60	- 1.50	-12.30	- 2.95
- 4.00	-12.90	- 1.80	-13.70	- 3.65
- 3.50	-15.05	- 2.00	-16.10	- 4.15
- 3.00	-17.45	- 2.05	-18.95	- 4.30
- 2.50	-19.10	- 2.12	-21.30	- 4.63
- 2.00	-19.90	- 2.20	-22.95	- 4.90
- 1.50	-20.15	- 2.50	-24.30	- 5.85
- 1.00	-19.85	- 2.90	-25.50	- 6.95
- 0.50	-21.75	- 3.35	-28.97	- 8.45
0.00	-26.93	- 4.00	-36.99	-10.45
0.50	-22.63	- 4.85	-33.76	-13.15
1.00	-18.35	- 5.60	-31.62	-15.80
1.50	-15.90	- 6.85	-30.80	-19.80
2.00	-15.20	- 7.85	-32.30	-23.30
2.50	-13.50	- 8.55	-31.96	-25.70
3.00	-11.05	- 8.50	-30.06	-26.59
3.50	-	- 8.90	-	-28.29
4.00	-10.90	- 8.65	-32.44	-29.17
4.50	-	- 8.40	-	-30.23
5.00	- 9.40	- 8.30	-32.30	-30.94
5.50	- 8.90	- 8.30	-32.16	-31.62
6.50	- 8.82	- 8.07	-33.15	-32.91
7.50	-10.25	- 9.25	-37.06	-35.43
8.50	- 9.80	- 9.35	-37.06	-36.04
10.00	-	- 9.55	-	-36.96
11.50	- 5.75	- 6.00	-31.48	-31.42
12.50	- 2.85	- 2.55	-26.59	-25.88
14.50	+ 2.85	+ 2.85	-17.85	-17.60
16.50	+ 7.00	+ 7.00	-	-11.18
18.50	+ 9.75	+ 9.75	- 7.65	- 7.05
20.50	+11.90	+11.90	- 4.10	- 3.75
22.50	+13.65	+13.65	- 1.95	- 1.30

— denotes average pressure for 2 or 3 static taps located across width of test section (see Figs. 12 and 13).

## APPENDIX B

### Finite Difference Equations

This Appendix provides the detailed derivations of the finite difference equivalents of the momentum and energy conservation equations (5) and (7) respectively. For convenience the following definitions are introduced:

$$S = \rho u(\mu + \rho E)$$

and

$$S' = \rho u \left( \frac{\mu}{P_r} + \frac{\rho E}{P_{rt}} \right)$$

These definitions permit the momentum and energy equations to be expressed as:

$$u \frac{\partial u}{\partial X} = - \frac{1}{2\rho} \frac{dP}{dX} + u \frac{\partial}{\partial \psi} \left[ S \frac{\partial u}{\partial \psi} \right] \quad (B-1)$$

$$u \frac{\partial T}{\partial X} = \frac{\gamma - 1}{2\rho\gamma} u \frac{dP}{dX} + \frac{\gamma - 1}{\gamma} u S \left[ \frac{\partial u}{\partial \psi} \right]^2 + u \frac{\partial}{\partial \psi} \left[ Q \frac{\partial T}{\partial \psi} \right] \quad (B-2)$$

Before approximating these equations with finite difference relations a system of grid lines parallel to the  $X$  and  $\psi$  axes must be introduced. As illustrated in Figure B-1, a nodal point coincides with each intersection of these lines. Lines parallel to the  $\psi$  axis are termed  $m$ -lines and those parallel to  $X$  axis  $n$ -lines. Each node is given a double subscript, the first being the number of the  $m$ -line passing through it, and the second the  $n$ -line number.

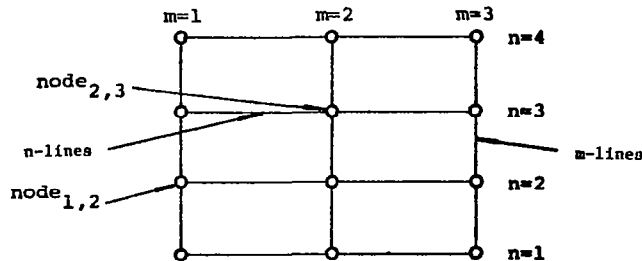


Figure B-1 Definition of Grid Lines for Finite Difference Solution

The values of the variables on the  $m = 1$  line are the known initial conditions. The conservation equations express for each node on the  $m = 2$  line its interrelation with other nodes on the  $m = 2$  line and nodes on the  $m = 1$  line. If  $m = 2$  line nodes are only related to nodes which lie on the  $m = 1$  line, the finite difference scheme is termed explicit. If an  $m = 2$  node is also related to a number of other  $m = 2$  nodes, the scheme is termed implicit (See Figure B-2).

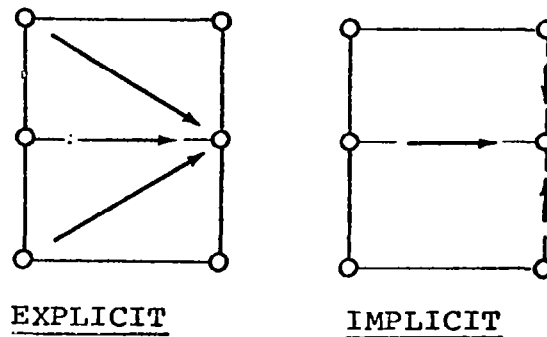


Figure B-2

#### Diagrams of Explicit and Implicit Solutions

The implicit form of finite difference schemes leads to a series of  $N$  simultaneous algebraic equations relating the known initial conditions on the  $m = 1$  line and the unknown variables on each of the  $N$  nodes on the  $m = 2$  line. After solution of these simultaneous equations, the variables on the  $m = 3$  line are expressed in terms of the known values on the  $m = 2$  line. Proceeding in this manner, a solution to the complete flow field is marched out. Although simpler to program, the explicit scheme shows unstable characteristics if the  $m$ -lines are widely spaced relative to the  $n$ -line spacing. Implicit schemes show much more stable characteristics and therefore allow much larger  $m$ -line spacings, thus reducing computation times. The computer procedure presented in this report employs a system of implicit finite difference approximations which are defined using the notation described in Figure B-3.

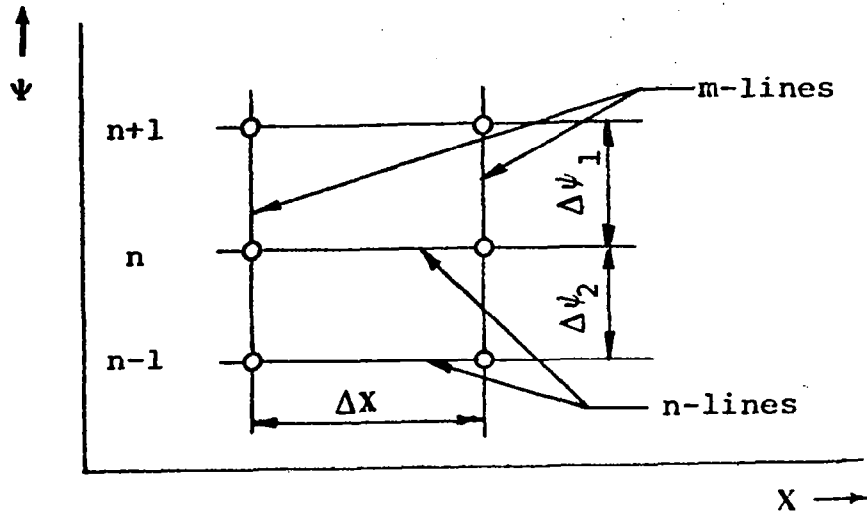


Figure B-3

#### Implicit Finite Difference Term Definition

The velocity at nodes  $n + 1$  and  $n - 1$  can be expressed in terms of a Taylor Series expanded about node  $n$ , on the same  $m$ -line,

$$u_{n+1} = u_n + \Delta\psi_1 \left. \frac{\partial u}{\partial \psi} \right|_n + \frac{(\Delta\psi_1)^2}{2} \left. \frac{\partial^2 u}{\partial \psi^2} \right|_n + \text{higher order terms} \quad (\text{B-3})$$

$$u_{n-1} = u_n - \Delta\psi_2 \left. \frac{\partial u}{\partial \psi} \right|_n + \frac{(\Delta\psi_2)^2}{2} \left. \frac{\partial^2 u}{\partial \psi^2} \right|_n + \text{higher order terms} \quad (\text{B-4})$$

Combining these equations to eliminate  $\left. \frac{\partial^2 u}{\partial \psi^2} \right|_n$  yields,

$$\frac{(\Delta\psi_2)^2}{2} u_{n+1} - \frac{(\Delta\psi_1)^2}{2} u_{n-1} = \frac{u_n}{2} (\Delta\psi_2^2 - \Delta\psi_1^2) + \left. \frac{\partial u}{\partial \psi} \right|_n \frac{1}{2} (\Delta\psi_1 \Delta\psi_2^2 + \Delta\psi_2 \Delta\psi_1^2) + \text{higher order terms}$$

Neglecting terms of the order  $(\Delta\psi)^3$  and higher, yields

$$\left. \frac{\partial u}{\partial \psi} \right|_n = \frac{\left( \frac{\Delta\psi_2}{\Delta\psi_1} \right) u_{n+1} - \left( \frac{\Delta\psi_1}{\Delta\psi_2} \right) u_{n-1} - \left( \frac{\Delta\psi_2}{\Delta\psi_1} - \frac{\Delta\psi_1}{\Delta\psi_2} \right) u_n}{\Delta\psi_2 + \Delta\psi_1}$$

Defining  $S_5 = \frac{\Lambda\psi_1}{\Lambda\psi_2(\Lambda\psi_2 + \Lambda\psi_1)}$

and  $S_4 = \frac{\Lambda\psi_2}{\Lambda\psi_1(\Lambda\psi_2 + \Lambda\psi_1)}$

yields

$$\left. \frac{\partial u}{\partial \psi} \right|_n = S_4 (u_{n+1} - u_n) + S_5 (u_n - u_{n-1}) \quad (B-5)$$

Similarly,

$$\left. \frac{\partial T}{\partial \psi} \right|_n = S_4 (T_{n+1} - T_n) + S_5 (T_n - T_{n-1}) \quad (B-6)$$

The second derivative term in the momentum equation is approximated using the following Taylor Series expansions,

$$\begin{aligned} \left( S \frac{\partial u}{\partial \psi} \right)_{n+1/2} &= S \left( \frac{\partial u}{\partial \psi} \right)_n + \frac{\Lambda\psi_1}{2} \frac{\partial}{\partial \psi} \left[ \left( S \frac{\partial u}{\partial \psi} \right)_n \right] + \frac{\Lambda\psi_1^2}{4} \frac{\partial^2}{\partial \psi^2} \left[ \left( S \frac{\partial u}{\partial \psi} \right)_n \right] \\ &\quad + \text{higher order terms} \end{aligned} \quad (B-7)$$

$$\begin{aligned} \left( S \frac{\partial u}{\partial \psi} \right)_{n-1/2} &= \left( S \frac{\partial u}{\partial \psi} \right)_n - \frac{\Delta\psi_2}{2} \frac{\partial}{\partial \psi} \left[ \left( S \frac{\partial u}{\partial \psi} \right)_n \right] + \frac{\Delta\psi_1^2}{4} \frac{\partial^2}{\partial \psi^2} \left[ \left( S \frac{\partial u}{\partial \psi} \right)_n \right] \\ &\quad + \text{higher order terms} \end{aligned} \quad (B-8)$$

Neglecting terms of the order of  $\frac{\Delta\psi^2}{4}$  and higher yields,

$$\begin{aligned}\frac{\partial}{\partial\psi} \left[ S \frac{\partial u}{\partial\psi} \right]_n &= \left[ \left( S \frac{\partial u}{\partial\psi} \right)_{n+1/2} - \left( S \frac{\partial u}{\partial\psi} \right)_{n-1/2} \right] \left[ \frac{2}{\Delta\psi_1 + \Delta\psi_2} \right] \\ &= \frac{1}{\Delta\psi_1 + \Delta\psi_2} \left[ \frac{(S_{n+1} + S_n)(u_{n+1} - u_n)}{\Delta\psi_1} \right. \\ &\quad \left. - \frac{(S_n + S_{n-1})(u_n - u_{n-1})}{\Delta\psi_2} \right]\end{aligned}\tag{B-9}$$

Similarly,

$$\begin{aligned}\frac{\partial}{\partial\psi} \left[ S' \frac{\partial T}{\partial\psi} \right]_n &= \frac{1}{\Delta\psi_1 + \Delta\psi_2} \left[ \frac{(S'_{n+1} + S'_n)(T_{n+1} - T_n)}{\Delta\psi_1} \right. \\ &\quad \left. - \frac{(S'_n + S'_{n-1})(T_n - T_{n-1})}{\Delta\psi_2} \right]\end{aligned}\tag{B-10}$$

The velocity at a node located at the intersection of the downstream m-line and any n-line  $u_{2,n}$  can be expressed in terms of the following Taylor Series,

$$u_{2,n} = u_{1,n} + \left. \frac{\partial u}{\partial X} \right|_n \Delta X + \left. \frac{\partial^2 u}{\partial X^2} \right|_n (\Delta X)^2 + \text{higher order terms}\tag{B-11}$$

Use of the boundary layer equations implies that gradients in the X- direction are much smaller than those in the  $\psi$ - direction. Therefore it is permissible to use a simpler approximation of the X- direction derivatives.

Neglecting terms of  $(\Delta X)^2$  and higher yields,

$$\left. \frac{\partial u}{\partial X} \right|_n = \frac{u_{2,n} - u_{1,n}}{\Delta X}\tag{B-12}$$

This approximation is termed "backward-difference".

Similarly,

$$\left. \frac{\partial T}{\partial X} \right|_n = \frac{T_{2,n} - T_{1,n}}{\Delta X} \quad (B-13)$$

The only terms in the energy and momentum equations which cannot be approximated using the preceding equations are those containing the pressure gradient  $dP/dX$ . Assuming this gradient varies linearly throughout the  $\Delta X$  interval yields:

$$\frac{dP}{dX} = \frac{1}{2} \left( \left. \frac{dP}{dX} \right|_{m=1} + \left. \frac{dP}{dX} \right|_{m=2} \right) \quad (B-14)$$

#### Momentum Equation

Combining equations (B-1), (B-9), (B-12) and (B-14) yields:

$$u_{1,n} \frac{(u_{2,n} - u_{1,n})}{\Delta X} = - \frac{1}{4\rho_{1,n}} \left[ \left. \frac{dP}{dX} \right|_{m=1} + \left. \frac{dP}{dX} \right|_{m=2} \right] + \frac{u_{1,n}}{2\psi_n} \left( \frac{1}{\Delta\psi_1} + \frac{1}{\Delta\psi_2} \right) \\ \left[ \frac{(S_{n+1} + S_n)(u_{2,n+1} - u_{2,n})}{\Delta\psi_1} - \frac{(S_n + S_{n-1})(u_{2,n} - u_{2,n-1})}{\Delta\psi_2} \right] \quad (B-15)$$

This equation can be expressed in the form

$$A_{n-1} u_{2,n} + B_{n-1} u_{2,n+1} + C_{n-1} u_{2,n-1} = D_{n-1} \quad (B-16)$$

in which the coefficients are defined by equations (23) through (28) of the main text.

#### Energy Equation

Combining equations (B-2), (B-5), (B-10), (B-13) and (B-14) yields

$$\begin{aligned}
\frac{u_{1, n} (T_{2, n} - T_{1, n})}{\Delta X} &= \frac{\gamma - 1}{\gamma} u_{1, n} s_{1, n} \left[ s_4 (u_{2, n+1} - u_{2, n}) \right. \\
&\quad \left. + s_5 (u_{2, n} - u_{2, n-1}) \right]^2 \\
&\quad + u_{1, n} \left[ \frac{1}{\Delta \psi_1 + \Delta \psi_2} \right] \left[ \frac{(s'_{n+1} + s'_n) (T_{2, n+1} - T_{2, n})}{\Delta \psi_1} \right] \\
&\quad - \frac{(s'_n + s'_{n-1}) (T_{2, n} - T_{2, n-1})}{\Delta \psi_2} \left] + \left( \frac{\gamma - 1}{\gamma} \frac{u_{1, n}}{4\rho_{1, n}} \right) \left[ \frac{dP}{dX} \Big|_{m=1} + \frac{dP}{dX} \Big|_{m=2} \right]
\end{aligned}
\tag{B-17}$$

This equation can be expressed in the form:

$$A_{n-1} \cdot T_{2, n} + B_{n-1} \cdot T_{2, n+1} + C_{n-1} \cdot T_{2, n-1} = D_{n-1}
\tag{B-18}$$

in which the coefficients are defined by equations (31) through (36) of the main text.



## APPENDIX C

### Solution Procedure

The calculation procedure starts at the upstream flow boundary, where the values of all flow variables must be known or assumed. Specification of the velocity and temperature distribution, dimensionless eddy viscosity, duct and nozzle inlet dimensions, and working fluid, defines all initial conditions.

The known initial conditions,  $m = 1$  line, are related to the unknown conditions,  $m = 2$  line, by the previously derived equations, and assumed boundary conditions. These inter-relations form a set of  $n - 2$  simultaneous algebraic equations, where  $n$  is the number of  $n$ -lines, and the equations are shown in Appendix B. The resultant matrix of coefficients is tridiagonal in form except for the initial and final rows which only contain two terms. Rapid, exact solutions to this type of matrix are obtained using the Thomas Algorithm, a successive elimination technique, which is described in this Appendix.

The solution for the variables on the  $m = 2$  line is iterative, because of the presence of the unknown pressure in the momentum equation. The procedure adopted was to estimate the pressure gradient, and solve the equations, using the algorithm. The equations automatically satisfy conservation of mass, momentum, and energy, but only one pressure gradient yields the correct wall geometry. The duct dimension corresponding to the estimated pressure gradient was calculated from the  $m = 2$  line variables. The pressure gradient was then incremented by a small percentage of its initial estimated value, and the calculation process repeated for a new duct dimension. A third estimate of the pressure gradient was obtained by interpolation between the two calculated, and the actual duct dimension. In almost all the calculations performed to date, this value has been acceptably close, within 0.001%, to the actual duct dimension. If this criterion is not met, a further iteration is applied, and a fourth solution obtained.

The now known variables on the  $m = 2$  line become the new  $m = 1$  line variables and the procedure is repeated for another set of  $m = 2$  line variables. Thus a solution to the complete flow field is marched out.

The difference form of the momentum and energy equation is:

$$A_{n-1}X_n + B_{n-1}X_{n+1} + C_{n-1}X_{n-1} = D_{n-1} \quad (C-1)$$

where  $X$  is either  $u$  or  $T$ . If the number of  $n$ -lines is  $n$ , there are  $n - 2$  equations of the form (1) and two equations expressing the boundary conditions. The first and the last equations represent the boundary conditions, which are:

$$u_1 = u_n = 0 \quad (C-2)$$

and

$$\left. \frac{\partial T}{\partial \psi} \right|_1 = \left. \frac{\partial T}{\partial \psi} \right|_n = 0 \quad (C-3)$$

Equation (C-2) can be written in terms of  $X$  as follows:

$$X_1 = X_2 = 0 \quad (C-4)$$

Equation (C-3) correspondingly becomes:

$$X_1 = X_2 \quad (C-5)$$

and

$$X_{n-1} = X_n \quad (C-6)$$

Equations (C-4), (C-5) and (C-6) can be written in terms of  $X$  as follows:

$$X_n = KX_{n-1} \quad (C-7)$$

$$X_1 = KX_2 \quad (C-8)$$

where  $K$  is 0 for the momentum equation and unity for the energy equation. Thus, the matrix form of the equation (C-1) is shown on the following page (Table C-1).

Table C-1

Matrix Form of Equation C-1 Designated as Equation C-1

$$\begin{bmatrix}
 1 & -K & 0 & 0 & 0 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 C_1 & A_1 & B_1 & 0 & 0 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 0 & C_2 & A_2 & B_2 & 0 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 0 & 0 & C_3 & A_3 & B_3 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 - & - & - & - & - & - & - & - & - & - & - & - & - \\
 0 & 0 & 0 & 0 & 0 & - & C_{n-1} & A_{n-1} & B_{n-1} & - & 0 & 0 & 0 \\
 - & - & - & - & - & - & - & - & - & - & - & - & - \\
 0 & 0 & 0 & 0 & 0 & - & 0 & 0 & 0 & - & C_{n-2} & A_{n-2} & B_{n-2} \\
 0 & 0 & 0 & 0 & 0 & - & 0 & 0 & 0 & - & 0 & -K & 1
 \end{bmatrix}
 \begin{bmatrix}
 X_1 \\
 X_2 \\
 X_3 \\
 X_4 \\
 X_{n-1} \\
 X_n \\
 X_{n+1} \\
 X_{n-1} \\
 X_n
 \end{bmatrix}
 =
 \begin{bmatrix}
 0 \\
 D_1 \\
 D_2 \\
 D_3 \\
 D_{n-2} \\
 D_{n-1} \\
 D_n \\
 D_{n-2} \\
 0
 \end{bmatrix}$$

The second equation is:

$$C_1 X_1 + A_1 X_2 + B_1 X_3 = D_1 \quad (C-9)$$

Substituting equation (C-4) into this equation yields:

$$A'_1 X_2 + B_1 X_3 = D_1 \quad (C-10)$$

where  $A'_1 = C_1 + A_1$

The  $n^{\text{th}} - 1$  equation is:

$$C_{n-2} X_{n-2} + A_{n-2} X_{n-1} + B_{n-2} X_n = D_{n-2} \quad (C-11)$$

Substituting equation (C-7) into this equation yields:

$$C_{n-2} X_{n-2} + A'_{n-2} X_{n-1} = D_{n-2} \quad (C-12)$$

where  $A'_{n-2} = A_{n-2} + KB_{n-2}$

Thus the  $n$  equations (C-8) can be reduced to the  $n - 2$  equations shown on Table C-2.

#### The Thomas Algorithm

Starting with the first equation,  $X_2$  can be expressed in terms of  $X_3$ . The second equation gives  $X_3$  in terms of  $X_4$ . Continuing through all the equations until the  $n^{\text{th}} - 3$  equation gives  $X_{n-2}$  in terms of  $X_{n-1}$ . Combining this with the last equation gives  $X_{n-1}$ . Working backwards through the equations then allows the remaining unknowns to be found. This procedure is most easily applied by defining the following:

$$\begin{aligned} W_1 &= A'_1 & g_1 &= \frac{D_1}{W_1} \\ Q_{n-1} &= \frac{B_{n-1}}{W_{n-1}} & n &= 2, 3, \dots, (n-2) \\ W_n &= A_n - C_n Q_{n-1} & n &= 2, 3, \dots, (n-2) \\ g_n &= \frac{D_n - C_n g_{n-1}}{W_n} & n &= 2, 3, \dots, (n-2) \end{aligned} \quad (C-14)$$

Table C-2

Matrix Form of Equation C-8 with Simplified Terms Designated as Equation C-13

$$\begin{bmatrix}
 A_1 & B_1 & 0 & 0 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 C_2 & A_2 & B_2 & 0 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 0 & C_3 & A_3 & B_3 & - & 0 & 0 & 0 & - & 0 & 0 & 0 \\
 - & - & - & - & - & - & - & - & - & - & - & - \\
 0 & 0 & 0 & 0 & - & C_{n-1} & A_{n-1} & B_{n-1} & - & 0 & 0 & 0 \\
 - & - & - & - & - & - & - & - & - & - & - & - \\
 0 & 0 & 0 & 0 & - & 0 & 0 & 0 & - & C_{n-3} & A_{n-3} & B_{n-3} \\
 0 & 0 & 0 & 0 & - & 0 & 0 & 0 & - & 0 & C_{n-2} & A'_{n-2}
 \end{bmatrix}
 \begin{bmatrix}
 X_2 \\
 X_3 \\
 X_4 \\
 X_{n-1} \\
 X_n \\
 X_{n+1} \\
 X_{n-2} \\
 X_{n-1}
 \end{bmatrix}
 =
 \begin{bmatrix}
 D_1 \\
 D_2 \\
 D_3 \\
 D_{n-2} \\
 D_{n-1} \\
 D_n \\
 D_{n-3} \\
 D_{n-2}
 \end{bmatrix}$$

Equations (C-13) then reduce to:

$$X_{n-1} = g_{n-2} \text{ and } X_n = g_{n-1} - Q_{n-1}X_{n+1} \quad n = (n-2), (n-3), \dots, 2 \quad (\text{C-15})$$

If the values of W, Q and g are calculated in order of increasing n using equations (C-14), then equations (C-15) can be used to calculate the values of X in order of decreasing X starting with  $X_{n-1}$ . To clarify this procedure, the method is now used to solve the following four simultaneous equations:

$$\begin{bmatrix} A'_1 & B_1 & 0 & 0 \\ C_2 & A_2 & B_2 & 0 \\ 0 & C_3 & A_3 & B_3 \\ 0 & 0 & C_4 & A'_4 \end{bmatrix} \begin{bmatrix} X_2 \\ X_3 \\ X_4 \\ X_5 \end{bmatrix} = \begin{bmatrix} D_1 \\ D_2 \\ D_3 \\ D_4 \end{bmatrix}$$

$$A'_1 X_2 + B_1 X_3 = D_1$$

$$W_1 = A'_1$$

$$Q_1 = \frac{B_1}{W_1}$$

$$g_1 = \frac{D_1}{W_1}$$

$$\text{hence } X_2 = g_1 - Q_1 X_3$$

$$A_2 X_3 + B_2 X_4 + C_2 X_2 = D_2$$

$$W_2 = A_2 - C_2 Q_1$$

$$Q_2 = \frac{B_2}{W_2}$$

$$g_2 = \frac{D_2 - C_2 g_1}{W_2}$$

$$\text{hence } X_3 = g_2 - Q_2 X_4$$

(C-16)

$$A_3 X_4 + B_3 X_5 + C_3 X_3 = D_3$$

$$W_2 = A_3 - C_3 Q_2$$

$$Q_3 = \frac{B_3}{W_3}$$

$$g_3 = \frac{D_3 - C_3 g_2}{W_3}$$

$$\text{hence } X_4 = g_3 - Q_3 X_5 \quad (C-17)$$

$$A'_4 X_5 + C_4 X_4 = D_4$$

$$W_4 = A_4 - C_4 Q_3$$

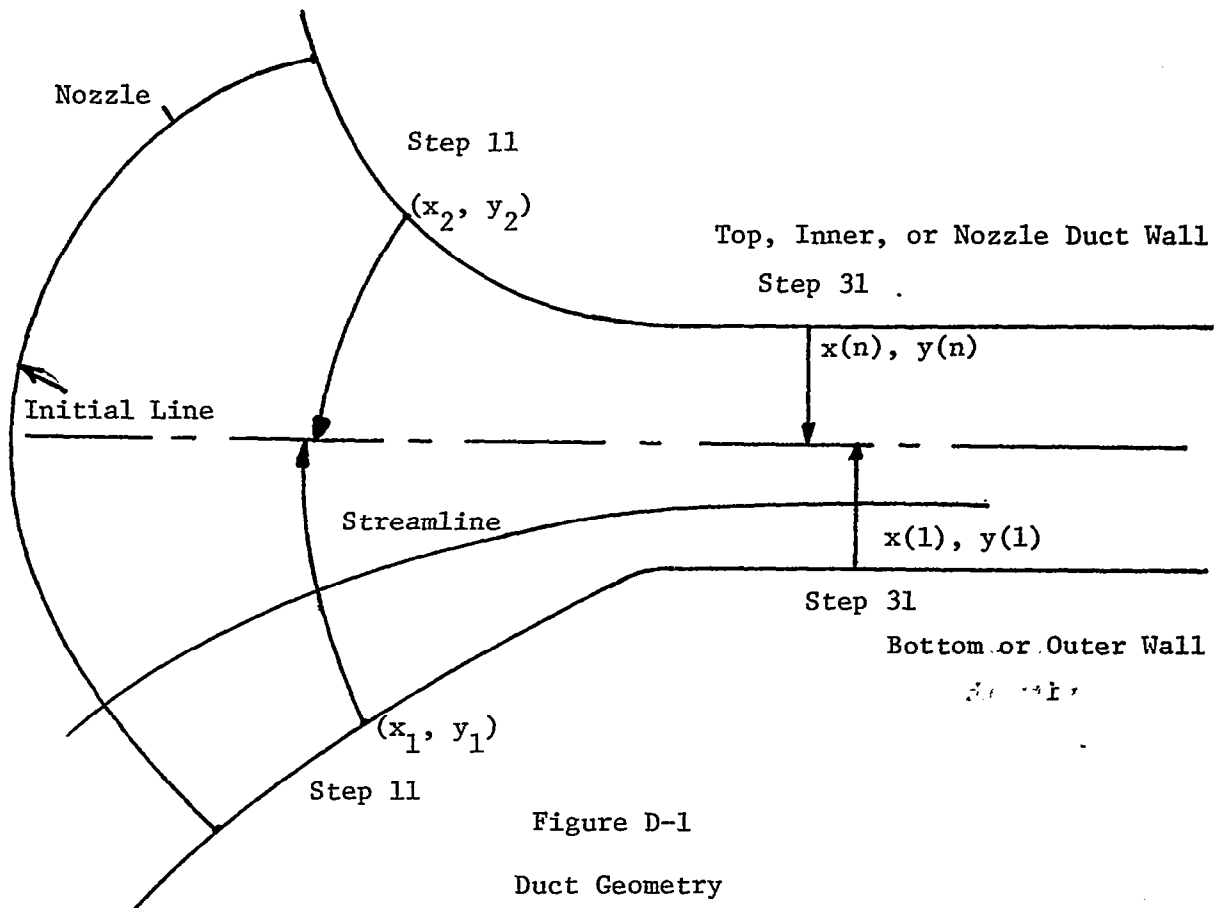
$$g_4 = \frac{D_4 - C_4 g_3}{W_4}$$

$$\text{hence } X_5 = g_4 \quad (C-18)$$

Substituting in equation (C-16) yields  $X_3$ . Equations (C-17) and (C-18) are special forms of equations (C-15) for  $n = 6$  and  $n = 4$ .

Appendix D  
COMPUTER PROGRAM

The computer program is designed to analyze the flow in a duct where geometry is shown in Figure D-1.



The computation proceeds from the initial line by moving along the lower and upper wall a specified distance and then defining two arcs from the new wall locations to the duct midpoint. Since only a rough approximation is available for the wall slopes the two arcs do not necessarily meet resulting in the wall points becoming slightly unsynchronized as the computation proceeds.



The program organization consists of a main program (NAS) and sixteen subroutines and function subroutines. The functions of the main program and subroutines are:

#### PROGRAM NAS

The main program is divided as follows:

##### Input and Data Initialization: Cards \$NA10 to \$NA2200

This section initializes the constants of the program, reads and prints the computation conditions and duct geometry, and puts this data in nondimensional form.

##### Initial Conditions: Cards \$NA2210 to \$NA2810

The subroutine INCOND is called to define the starting conditions. The initial flow conditions are then put in dimensional form and printed.

##### Main Body of Program: Cards \$NA2820 to \$NA4810

The computation proceeds down the duct in a sequence of steps. Values of pressure, temperature, velocity, density etc. are computed which are consistent with the previous step values and the geometry of the duct. The process stops when the end of the duct is reached.

##### Eddy Viscosity: Card \$NA3010

The dimensionless eddy viscosity is calculated in subroutine EDDY using data from the preceding step.

##### Streamline Step Size: Cards \$NA3100 to \$NA3400

An appropriate step size along the duct is determined for the streamline of maximum velocity. Consistent step sizes are then determined for all other streamlines.

Pressure Gradient Approximation: Cards \$NA3730 to \$NA4540

The pressure gradient is determined by selecting a value such that the computed duct widths minus the actual duct width equals  $0 \pm .0001$ .

Flow Boundaries: Cards \$NA4550 to \$NA4700

At each step the boundaries of the flow regions are checked to determine when shear layers vanish.

Output Section: Cards \$NA4820 to \$NA5750

The flow variables are presented in dimensional form at preselected intervals.

Figure D-2 is a flow chart of the main program NAS.

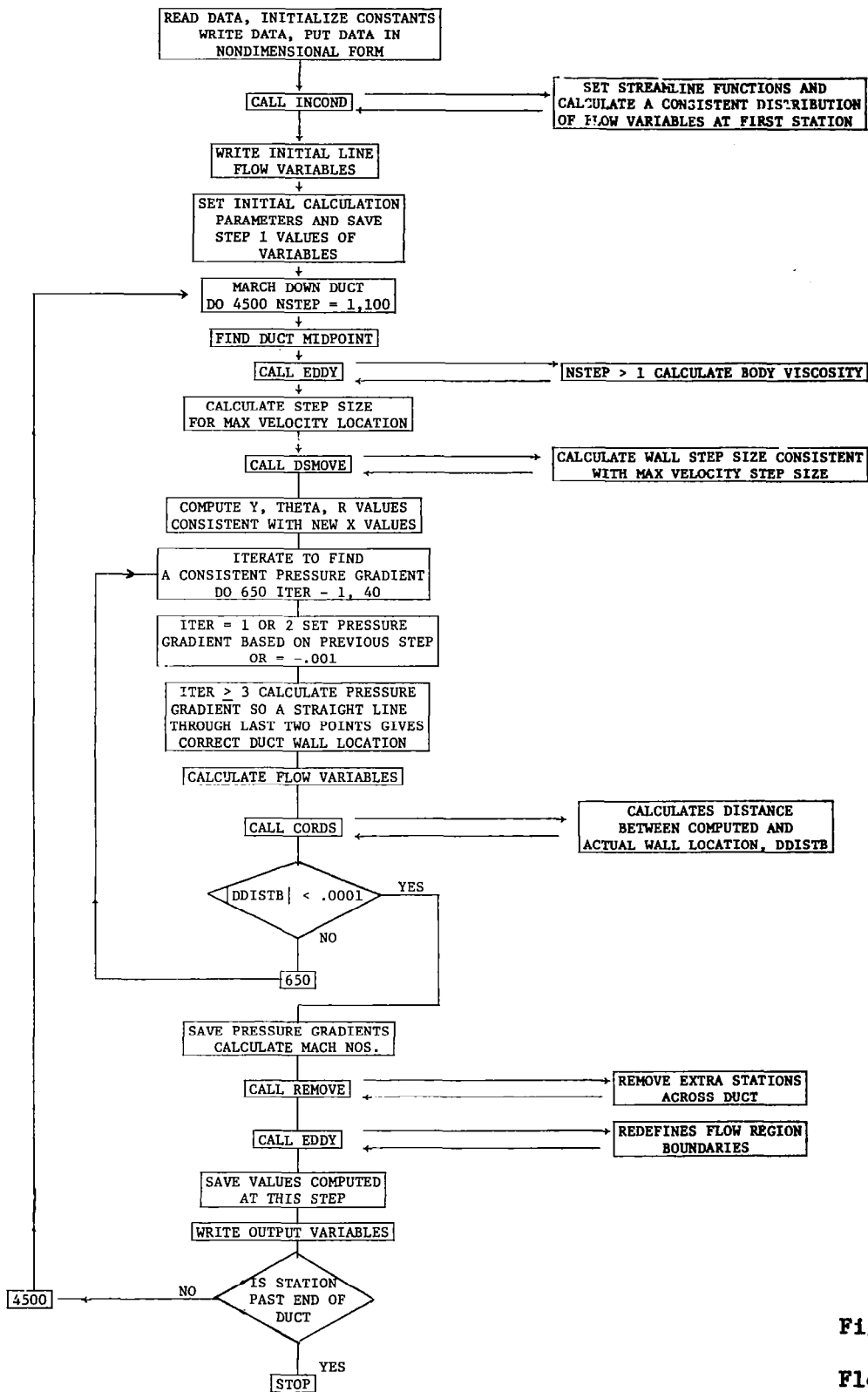


Figure D-2  
Flow Chart

## SUBROUTINES

### INCOND

This subroutine contains the computation which defines the flow conditions at the initial station (Fig. D-3). The steps used in the process are:

1. Points on the boundaries  $(x_1, y_1)$ ,  $(x_2, y_2)$  are defined such that two circular arcs are normal to the wall at  $(x_1, y_1)$  and  $(x_2, y_2)$  and pass through the nozzle  $(x_{noz}, y_{noz})$ .
2. A flow split is determined.
3. The subroutine OMGSET is called to set the streamline locations.
4. The subroutine TMPSET is called to determine the temperature distribution.
5. The remaining flow variables are calculated.
6. The location of the nozzle streamline checked to see if it is within tolerance of its specified location. If it is then the computation is returned, if not a new flow split is determined and the calculation is repeated.

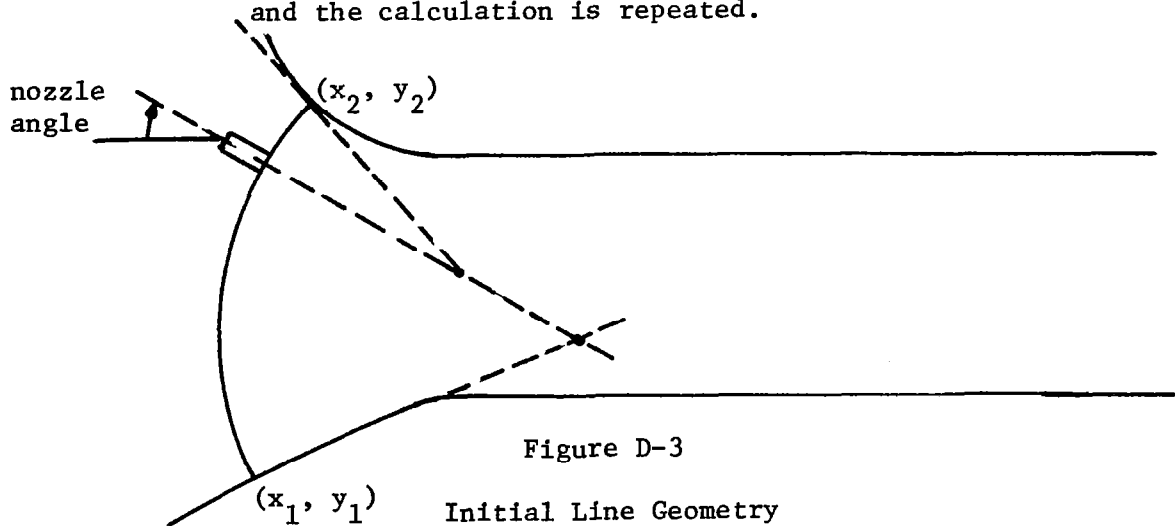


Figure D-3

Initial Line Geometry

## OMGSET

This subroutine sets the distribution of streamlines. The flow is divided into regions of width DOMG1, DOMG2, DOMG3 (Fig. D-4 and D-5). Given an initial spacing at each flow edge (DMWALL and DMJET) and a specific number of points (KWALL, KJET) the subroutine fills in the spaces close to the edges.

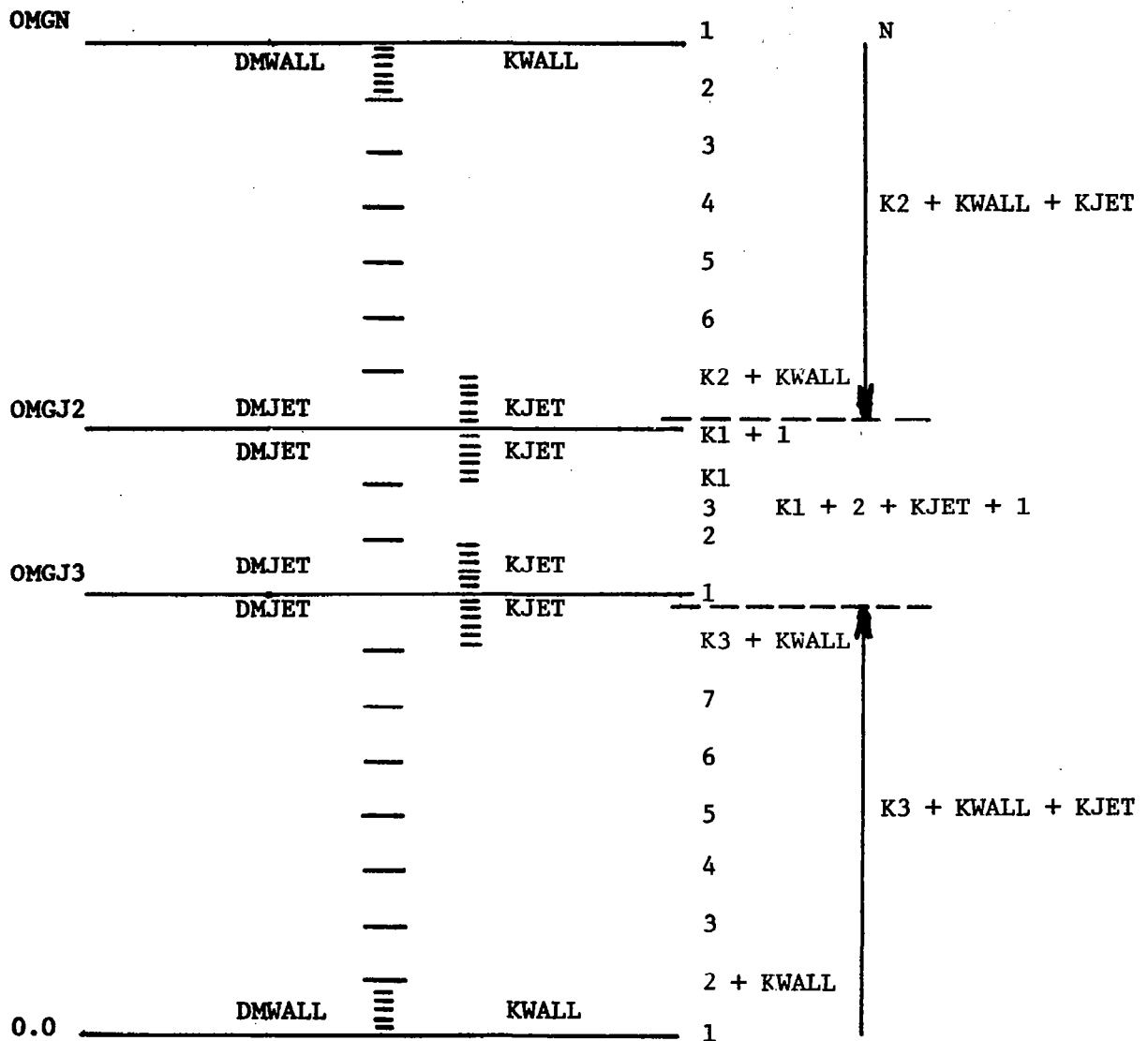


Figure D-4

Streamline Locations

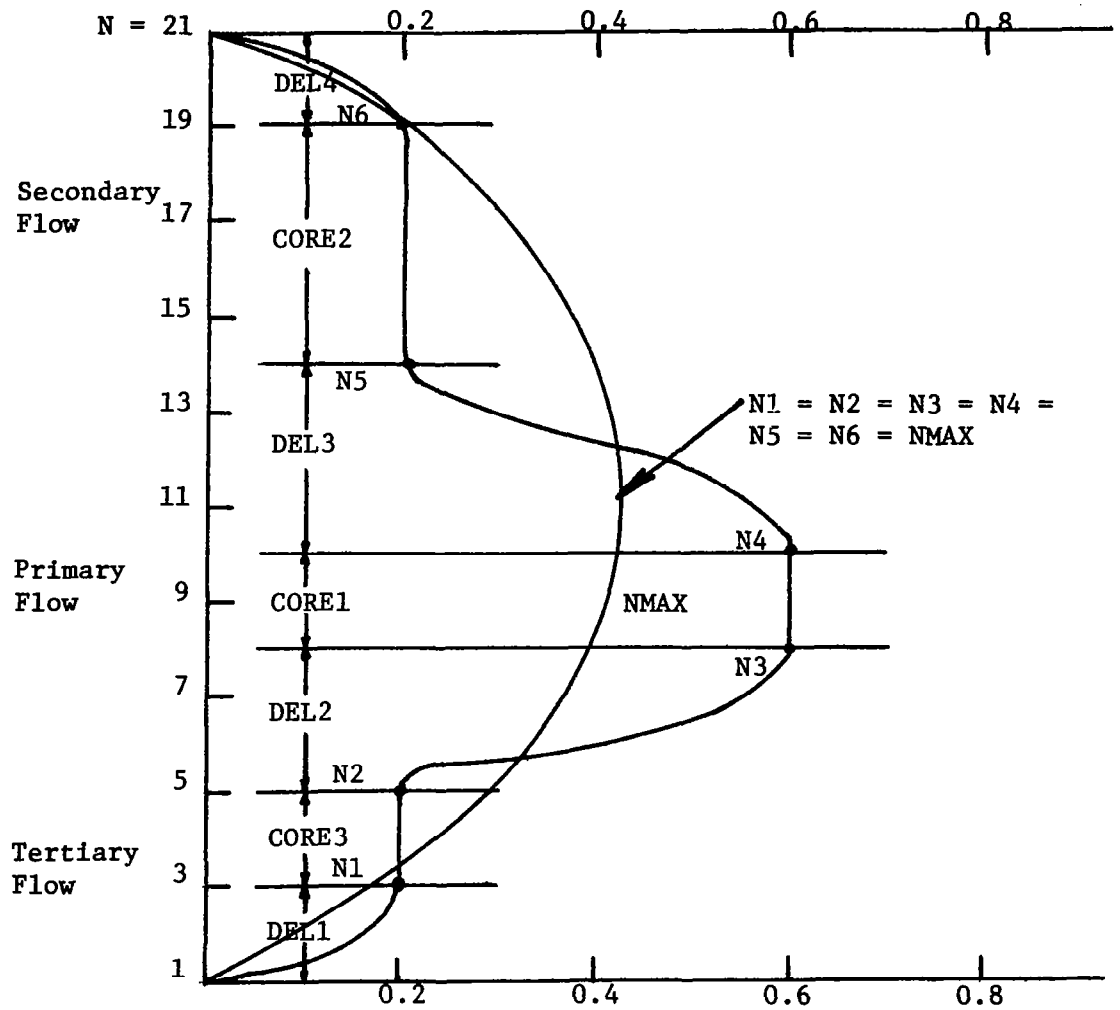


Figure D-5

Flow Structure

### TMPSET

An initial value of static temperature at the nozzle,  $TS(NM1D)$ , is selected. Values of the temperature are then computed at the outer wall con-

sistent with the wall curvature. The distance between the wall and nozzle consistent with TS(NMID) is returned to the calling program.

#### WALLS

Wall curvatures are calculated from the wall geometry data. This data is then used to interpolate values of wall location slope and curvature for specific axial locations. The computer calculates curvatures assuming the data is segmented with at least three points per segment. When five or more points are available a least squares parabola using the data points and the two points to either side is used to smooth the data. Subsequent interpolations are parabolic.

#### LSQ

This subroutine calculates the least squares parabolas for walls.

#### EDDY

This subroutine contains the eddy viscosity calculation.

#### DSMOVE

This subroutine contains the calculation of the distance moved along the walls in each step.

#### SDWE

This subroutine contains the calculation of temperatures and velocities at a computing station.

#### CALC

This is a simultaneous linear equation solution subroutine.

#### CORDS

The distance between streamlines and the x, y coordinates of the

streamlines are computed in this subroutine. The value DDIST, the difference between the y coordinates of the two arcs at the middle duct point, is returned.

#### LOOK

The algorithm identifies an edge of a core region and then checks to see if a core region that existed in the previous step has disappeared. If this has occurred that core width is set to 0.0. Core regions may disappear in any order.

#### REMOVE

This subroutine is used to remove grid points as the computation proceeds in order to reduce computing time. This is accomplished by every tenth step scanning the velocity array to see if the velocity gradients in the shear regions are less than a tolerance.

#### ARCDIS

This subroutine defines wall coordinates such that the arc is normal to the wall and nozzle.

#### FUNCTION SUBROUTINES

##### CENTRE

This is used in conjunction with ARCDIS. Given a nozzle and wall location it projects the two tangents, finds their intersection and returns the difference in length.

##### CURVDS

Used in DSMOVE to calculate distances along an arc.

##### SERIES

Used in OMGSET to determine streamline spacing in the boundary and shear layers.



## INPUT DATA

### Card 1 - Format 11,18A4

NUNIT - units indicator

NUNIT = 0 English units

NUNIT  $\neq$  0 S.I. units

TITLE(I) - title, printed at start of first output page

### Card 2 - Format 6F10.0

P01 - primary stagnation pressure - psia or pascal

T01 - primary stagnation temperature - Rankine or Kelvin

P02 - secondary stagnation pressure - psia or pascal

T02 - secondary stagnation temperature - Rankine or Kelvin

MASS1 - primary mass flow rate - lbm/sec-in or kg/sec-m

MASS2 - secondary + tertiary flow - lbm/sec-in or kg/sec-m

### Card 3 - Format 5I5

K1 - no. of grid points in primary flow (even)

K2 - no. of grid points in secondary flow (even)

K3 - no. of grid points in tertiary flow (even)

NPCYCL - print cycle, (e.g. NPCYCL = 10 causes print every ten steps)

NQUICK - 0 = full print; 1 = partial printout

### Card 4 - Format 4F10.0

XNOZ - x coordinate of nozzle center - in. or meters

YNOZ - y coordinate of nozzle center - in. or meters

TNOZ - nozzle angle - degrees

N - nozzle slot width - in. or meters

### Card 5 - Format 11I5

MST - number of segments used to describe lower wall geometry,

$$10 \geq \text{MST} \geq 1$$

NDI(I, 1) - MST values of the data point number of the segment end point. NDI(MST,1) = NPAIR1, the total number of points needed to describe the wall.

$$\text{NDI}(I + 1, 1) - \text{NDI}(I, 1) \geq 2$$

#### Cards 6 - Format 2F10.0 - NPAIR1 Cards

XW(I, 1) - x lower wall coordinate

YW(I, 1) - y lower wall coordinate

#### Card 7 - Format 11I5

MSS - number of segments used to describe upper wall geometry

$$10 \geq \text{MSS} \geq 1$$

NDI(I,2) - MSS values of the data point number of the segment end point, NDI(MSS, 1) = NPAIR2 the total number of points needed to describe the wall.

$$\text{NDI}(I + 1, 2) - \text{NDI}(I, 2) \geq 2$$

#### Cards 8 - Format 2F10.0 - NPAIR2 Cards

XW(I, 2) - x upper wall coordinate

YW(I, 2) - y upper wall coordinate

#### Card 9 - Format I5,10F5.0

NFULL - number of values of x for which full output is required

XPRØF(I) - NFULL values of x for which full output is required

The input data K1, K2, and K3 (Card 3) denote the number of grid points in the primary, secondary and tertiary flow. These are arbitrary numbers chosen so as to give the desired spacing of output data in the three regions. A total K1 + K2 + K3 of between 20 and 30 should give good results.

The data numbers MST, MSS (Cards 5 and 7) denote the number of segments needed to describe the wall geometry. For walls with continuous

curvature values, a value of 1 is sufficient. Values of MST and MSS greater than 1 allow the user to describe walls with discontinuous slopes and curvatures. The data  $NDI(I,2)$  and  $NDI(I,1)$  are the data point numbers at the boundary segment ends.

A sample of input data is shown in Table D-1.

Table D-1

## Sample of Input Data

1	TEST CASE	5	METRIC			
253934.	322.	102249.	301.	1.696	7.393	
8	10	10	5	0		
-.07366	.04023	-35.		.0030861		
6	12	16	26	31	35	43
-.1524	-.10396					
-.1511	-.1024					
-.1499	-.10033					
-.1397	-.08694					
-.1270	-.07450					
-.1143	-.06535					
-.1016	-.05837					
-.0889	-.05339					
-.0762	-.05004					
-.0635	-.04712					
-.0597	-.04678					
-.0587	-.04671					
-.0508	-.04517					
-.0254	-.04023					
0.0000	-.03523					
.01905	-.03160					
.02159	-.03115					
.02032	-.03137					
.02549	-.03048					
.03810	-.02870					
.05080	-.02736					
.06350	-.02647					
.07620	-.02596					
.08636	-.02564					
.08763	-.02560					
.08890	-.02558					
.10160	-.02536					
.12700	-.02492					
.15240	-.02448					
.17780	-.02404					
.19050	-.02383					
.20320	-.02383					
.22860	-.02383					
.25400	-.02383					
.26670	-.02383					
.27940	-.02449					
.30480	-.02582					
.35560	-.02848					
.40640	-.03115					
.45720	-.03381					
.50800	-.03649					
.55880	-.03913					
.58420	-.04038					

Table D-1 (Concluded)

5	16	26	31	35	43
-.0845		.156410			
-.0838		.143230			
-.0826		.134200			
-.0794		.120680			
-.0762		.111240			
-.0699		.097190			
-.0635		.086450			
-.0572		.077670			
-.0508		.070240			
-.0381		.058260			
-.0254		.049090			
-.0127		.042030			
0.0000		.036730			
.01270		.032960			
.01588		.032240			
.01905		.03160			
.02032		.03137			
.02159		.03115			
.02549		.03048			
.03810		.02870			
.05080		.02736			
.06350		.02647			
.07620		.02596			
.08636		.02564			
.08763		.02560			
.08890		.02558			
.10160		.02536			
.12700		.02492			
.15240		.02448			
.17780		.02404			
.19050		.02383			
.20320		.02383			
.22860		.02383			
.25400		.02383			
.26670		.02383			
.27940		.02449			
.30480		.02582			
.35560		.02848			
.40640		.03115			
.45720		.03381			
.50800		.03649			
.55880		.03913			
.58420		.04038			
1	.038				

## OUTPUT DATA

The program printed output consists of:

1. Title

2. Input flow and geometry (Cards 2, 3, and 4)

3. Wall Geometry

X, Y - wall coordinates (smoothed)

CURV - negative of wall curvature

4. Values of X at which full output is required (values of XPROF(I) from Card 9)

5. Flow split

SPLIT - ratio of tertiary flow to the sum of secondary and tertiary flow

6. Initial conditions along the computing station through the nozzle

ISENTROPIC NOZZLE THRUST PER UNIT WIDTH - lb/in or N/m

AMASS1, AMASS2, AMASS3 - primary, secondary and tertiary flow rates - lb/s/in or kg/s/m

J - station number

X, Y - coordinates of computing station streamline intersection - in or m

DN - distance from lower wall - in or m

THETA - streamline angle - degrees

K - negative of streamline curvature - 1/in or 1/m

OMG - streamline function

PO, PS - total and static pressure - psi or pascal

TS - static temperature - degrees Rankine or Kelvin

RHO - density - lb/f\*\*3 or kg/m\*\*3

U - speed - f/s or m/s

M - Mach number

7. Flow description downstream - partial output

NSTEP - step number

XI(XO) - coordinate along lower (upper) wall - in or m

PI(PO) - static pressure at lower (upper) wall - in of  
H<sub>2</sub>O or pascal

USTARI(USTARØ) - friction velocity at lower (upper) wall - f/s  
or m/s

KI(KO) - negative of wall curvature at XI(XO) - 1/in or  
1/m

RNI(RNØ) - Richardson number at lower (upper) wall

UMAX - maximum velocity, f/s or m/s

NUMBER OF ITERATIONS

DELTA X - step size at lower wall - in or m

DELTASI(DELTASO) - distance increment of lower (upper) wall - in  
or m

SI(SO) - cumulative distance along lower (upper) wall -  
in or m

SHRIN(SHROR) - shear stress at lower (upper) wall - psi or  
pascal

TOTAL AXIAL MOMENTUM PER UNIT WIDTH - lbf/in or N/m

THRUST RATIO -

8. Flow description downstream - full output

Same as partial output plus

N1 ... N6 - streamline numbers at edges of boundary and  
shear layers

CORE1, CORE2, CORE3 - widths of primary, secondary, and tertiary  
regions - in or m

DEL1, DEL2, DEL3, DEL4 - widths of boundary and shear layers - in or m

J - streamline number

1/R - negative of curvature - 1/in or 1/m

THETA - angle of streamline - degrees

X, Y - computing mode location - in or m

YREL - relative y coordinate

UREL - relative velocity (UMAX is normalizing quantity)  
POREL - relative total pressure  
PS - static pressure - psia or pascal  
E - Eddy viscosity  
TOTEMP - total temperature - degrees Rankine or Kelvin

9. In addition several warnings may be printed.

NO CONVERGENCE, DPDSA = ... DPDSB = ... DDISTB ... STEP 4

Convergence was not achieved in establishing the pressure gradient. The criteria for convergence is that  $|DDISTB|$  the distance between the computed and actual wall location be  $<.0001$ . DPDSA and DPDSB are pressure gradient increments.

NEGATIVE SHEAR STRESS AT THIS STATION. THIS INDICATES POSSIBLE SEPARATION. SUBSEQUENT RESULTS SHOULD BE USED WITH CAUTION.

The message is self-explanatory. When this occurs the calculation may be unstable and may stop for a variety of reasons. When this occurs the full output at that station is printed.



Figure D-6

## COMPUTER PROGRAM LISTING

```

C      PROGRAM NAS(INPUT,OUTPUT,TAPES=INPUT,TAPE6=OUTPUT)          $NA 0
C                                                                    $NA 10
C      THIS PROGRAM CALCULATES VENTED COANDA JET IN A CURVED DUCT   $NA 20
C      BY A FINITE-DIFFERENCE METHOD                                $NA 30
C                                                                    $NA 40
C      THE PROGRAM CALCULATES THE SPLIT BETWEEN SECONDARY          $NA 50
C      AND TERTIARY MASS FLOW                                       $NA 60
C                                                                    $NA 70
C      READ VARIABLES                                              $NA 80
C      NUNIT...UNITS INDICATOR--NUNIT=0..ENGLISH UNITS             $NA 90
C      NUNIT GT 0 ..SI UNITS                                       $NA 100
C      P01...PRIMARY STAGNATION PRESSURE--PSIA OR PASCAL           $NA 110
C      T01...PRIMARY STAGNATION TEMPERATURE--RANKINE OR KELVIN     $NA 120
C      P02...SECONDARY STAGNATION PRESSURE--PSIA OR PASCAL         $NA 130
C      T02...SECONDARY STAGNATION TEMPERATURE--RANKINE OR KELVIN   $NA 140
C      MASS1...PRIMARY MASS FLOW RATE--LBM/SEC-IN. OR KG/SEC-M     $NA 150
C      MASS2...TOTAL OF SECONDARY PLUS TERTIARY FLOW RATE--LBM/SEC-IN. $NA 160
C      OR KG/SEC-M                                                 $NA 170
C      K1...NO. OF GRID PTS. IN PRIMARY FLOW (EVEN NUMBER)        $NA 180
C      K2...NO. OF GRID PTS. IN SECONDARY FLOW (EVEN NUMBER)       $NA 190
C      K3...NO. OF GRID PTS. IN TERTIARY FLOW (EVEN NUMBER)        $NA 200
C                                                                    $NA 210
C      NPCYCL... PRINT CYCLE                                       $NA 220
C      NQUICK... 0 - FULL PRINT, 1 - PARTIAL PRINT OUT             $NA 230
C      XNOZ... X COORD OF CENTRE OF NOZZLE ...IN. OR METERS        $NA 240
C      YNOZ... Y COORD OF CENTRE OF NOZZLE .. IN. OR METERS       $NA 250
C      TNOZ... NOZZLE ANGLE DEGREES                                $NA 260
C      D... NOZZLE SLOT WIDTH INCHES OR METERS                     $NA 270
C                                                                    $NA 280
C      NPAIR1... NUMBER OF DUCT WALL COORDS, LOWER WALL            $NA 290
C      NPAIR2... NUMBER OF DUCT WALL COORDS, UPPER WALL            $NA 300
C      XW... ARRAY OF DUCT WALL X COORDS .. IN. OR METERS         $NA 310
C      YW... ARRAY OF DUCT WALL Y COORDS .. IN. OR METERS         $NA 320
C                                                                    $NA 330
C      NFULL... NUMBER OF PROFILE PRINTOUT LOCATIONS               $NA 340
C      XPROF... ARRAY OF LOCATIONS WHERE PROFILES REQUIRED IN. OR M. $NA 350
C                                                                    $NA 360
C      IT IS ASSUMED THAT P02=P03, T02=T03                          $NA 370
C                                                                    $NA 380
C      DATA VARIABLES                                             $NA 390
C      RG...GAS CONSTANT--FT.LBF/LBM.R OR N-M/KG-K                $NA 400
C      G...SPECIFIC HEAT CONSTANT                                  $NA 410
C      NUREF...REFERENCE KIN. VISCOSITY--FT**2/SEC OR M**2/SEC    $NA 420
C      TREF...REFERENCE TEMPERATURE--RANKINE OR KELVIN             $NA 430
C      RHOREF...REFERENCE DENSITY--LBM/FT**3 OR KG/M**3           $NA 440
C      GC...CONSTANT--LBM.FT/LBF.SEC**2 OR KG-M/N-SEC**2         $NA 450
C      PR...PRANDTL NUMBER                                         $NA 460
C      PRT...TURBULENT PRANDTL NUMBER                              $NA 470
C      TW...WALL TEMPERATURE--RANKINE OR KELVIN                    $NA 480
C                                                                    $NA 490
C      DIMENSIONLESS VARIABLES                                     $NA 500
C      OMG(I) STREAMLINE FUNCTIONS                                 $NA 510
C      S1,S2,S3,S4,S5(I) OMG DIFFERENCES                          $NA 520
C      DS(I) STREAMLINE STEP SIZES                                 $NA 530
C      R(I,2) CURVATURES                                           $NA 540
C      THETA(I,2) ANGLES                                           $NA 550
C      PS(I,2) PRESSURES                                           $NA 560
C      U(I,2) VELOCITIES                                           $NA 570
C      TS(I,2) TEMPERATURES                                         $NA 580
C      RHO(I,2) DENSITIES                                           $NA 590
C      VIS(I) VISCOSITIES                                           $NA 600
C      DN(I) DISTANCE TO INNER WALL                                $NA 610
C      X(I) X COORDINATE                                           $NA 620

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C	Y(I)	Y COORDINATE OF STEAMLINES	\$NA 630
C	M(I)	MACH NUMRERS	\$NA 640
C	P0(I)	STAGNATION PRESSURES	\$NA 650
C	E(I)	EDDY VISCOSITIES	\$NA 660
C			\$NA 670
C		PRESSURES NORMALIZED BY...P01	\$NA 680
C		TEMPERATURES NORMALIZED BY...T01	\$NA 690
C		VELOCITIES NORMALIZED BY...UREFP	\$NA 700
C		LENGTHS NORMALIZED BY ...0	\$NA 710
C		DYNAMIC VISCOSITES NORMALIZED BY...RH001*UREFP*D	\$NA 720
C		KINEMATIC EDDY VISCOSITY NORMALIZED BY...UREFP*D	\$NA 730
C		STREAM FUNCTION NORMALIZED BY...SQRT(RH001*UREFP*D)	\$NA 740
C		MASS FLOW RATE NORMALIZED BY...RH001*UREFP*D	\$NA 750
C		R(I) IS THE INVERSE OF THE DIMENSIONLESS RADIUS OF CURVATURE	\$NA 760
C			\$NA 770
	REAL NUREF,KS,LM,MASS1,MASS2,M(190)		\$NA 780
	DIMENSION XW(99,2),YW(99,2),S1(190),S2(190),S3(190), S4(190),		\$NA 790
1	S5(190),E(190),E1(190),DS(190),X(190),Y(190),VIS(190),		\$NA 800
2	P0(190),H(190),THETA(190,2),R(190,2),U(190,2),		\$NA 810
3	TS(190,2),RHO(190,2),DIP(190),NX(6),XPROF(10)		\$NA 820
	DIMENSION OMG(190),DN(190),PS(190,2)		\$NA 830
C	DOUBLE OMG(190),DN(190),PS(190,2)		\$NA 840
	DIMENSION PSI(180)		\$NA 850
C	DOUBLE PSI(190)		\$NA 860
	DIMENSION NOI(10,2),CURV(99,2),TITLE(18)		\$NA 870
	COMMON/KURV/MST,MSS,NDI,CURV		\$NA 880
	COMMON/SOLV/ S1,S2,S3,E,E1,H,S4,S5		\$NA 890
	COMMON/INEDY/PSI,RHO,VIS		\$NA 900
	COMMON/INCD/ U,PS,TS,M,P0,THETA,R,X,Y,DN		\$NA 910
	COMMON /WALL/ NPAIR1, NPAIR2, XW, YW		\$NA 920
	COMMON /CONST/ P01,P02,P03, T01,T02,T03, X1, XC, YC		\$NA 930
	DATA RG,G,GC,TRFF,RHOREF,TW/53.2,1.4,32.2,530.,0.075,530.0/		\$NA 940
	DATA NUREF,PR,PRT/0.00014,0.7,0.9/		\$NA 950
	DATA PRT1,PRT1/1.0,1.0/		\$NA 960
	DATA A1,A2,A3/144.,198.7,.03611/		\$NA 970
C	-----		\$NA 980
C			\$NA 990
C	INPUT DATA SECTION		\$NA1000
C			\$NA1010
C	-----		\$NA1020
	READ(5,800) NUNIT,(TITLE(I),I=1,18)		\$NA1030
	READ(5,801)P01,T01,P02,T02,MASS1,MASS2		\$NA1040
	READ(5,802) K1,K2,K3,NPCYCL,NQUICK		\$NA1050
	READ(5,801) XNOZ,YNOZ,TNOZ,D		\$NA1060
	READ(5,802) MST,(NDI(I,1),I=1,MST)		\$NA1070
	NPAIR1=NDI(MST,1)		\$NA1080
	READ(5,803) (XW(I,1),YW(I,1),I=1,NPAIR1)		\$NA1090
	READ(5,802) MSS,(NDI(I,2),I=1,MSS)		\$NA1100
	NPAIR2=NDI(MSS,2)		\$NA1110
	READ(5,803) (XW(I,2),YW(I,2),I=1,NPAIR2)		\$NA1120
	READ(5,804) NFULL,(XPROF(J),J=1,NFULL)		\$NA1130
800	FORMAT(I1,18A4)		\$NA1140
801	FORMAT(6F10.0)		\$NA1150
802	FORMAT(11I5)		\$NA1160
803	FORMAT(2F10.0)		\$NA1170
804	FORMAT(15,10F5.0)		\$NA1180
	XPROF(NFULL+1)= XW(NPAIR1,1)		\$NA1190
	WRITE(6,900)(TITLE(I),I=1,18)		\$NA1200
	IF(NUNIT.GT.0) GO TO 9		\$NA1210
	WRITE(6,901) P01,T01,P02,T02,MASS1,MASS2,D		\$NA1220
	WRITE(6,902) XNOZ,YNOZ,TNOZ		\$NA1230
	GO TO 16		\$NA1240
9	WRITE(6,914) P01,T01,P02,T02,MASS1,MASS2,D		\$NA1250
	WRITE(6,915) XNOZ,YNOZ,TNOZ		\$NA1260
	RG=284.		\$NA1270
	GC=1..		\$NA1280

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TREF=295.
RHOREF=1.201
TW=295.
NUREF=1.3E-5
A1=1.
A2=110.
A3=1.0
16 WRITE(6,903) K1,K2,K3,NPCYCL
WRITE(6,904)
CALL WALLS(XW,YW,1000.,0.,T1,R1,NPAIR1,1)
CALL WALLS(XW,YW,1000.,0.,T2,R2,NPAIR2,2)
NPAIR= NPAIR1
IF (NPAIR1.GT. NPAIR2) NPAIR= NPAIR2
WRITE(6,905) (J,XW(J,1),YW(J,1),CURV(J,1),XW(J,2),YW(J,2),
1 CURV(J,2),J=1,NPAIR)
WRITE(6,906) (XPROF(I),I=1,NFULL)
900 FORMAT(1H1,39X,43(1H*)/40X,1H*,41X,1H*/40X,1H*,2X,
1 40HCOANDA EFFECTS IN A PLANE CURVED DUCT */
2 40X,1H*,41X,1H*/40X,43(1H*)//18A4//)
901 FORMAT(25X,29HPRIMARY STAGNATION PRESSURE =,F6.2,5H PSIA/
1 25X,26HPRIMARY STAGNATION TEMP. =,F7.2,16H DEGREES RANKINE
2/ 25X,31HSECONDARY STAGNATION PRESSURE =,F6.2,5H PSIA/
3 25X,28HSECONDARY STAGNATION TEMP. =,F7.2,16H DEGREES RANKI
4NE// 25X,24HPRIMARY MASS FLOW RATE =,F8.5,12H LBM/SEC-IN./
5 25X,44HTOTAL OF SECONDARY PLUS TERTIARY FLOW RATE =,F8.5,
612H LBM/SEC-IN.//25X,19HNOZZLE SLOT WIDTH =,F6.3,4H IN./)
902 FORMAT(25X,21HNOZZLE X COORDINATE =,F10.4,4H IN./
1 25X,21HNOZZLE Y COORDINATE =,F10.4,4H IN./
2 25X,14HNOZZLE ANGLE =,F10.3,8H DEGREES/)
903 FORMAT(25X,39HNUMBER OF GRID POINTS IN PRIMARY FLOW =,I4/
1 25X,41HNUMBER OF GRID POINTS IN SECONDARY FLOW =,I4/
2 25X,40HNUMBER OF GRID POINTS IN TERTIARY FLOW =,I4/
3 25X,12HPRINT CYCLE ,I4)
904 FORMAT(1H1,10X,25HMIXING SECTION DIMENSIONS////
1 17X,10HLOWER WALL,35X,10HUPPER WALL/
2 6X,1HJ,10X,1HX,9X1HY,11X,4HCURV,19X,1HX,9X,1HY,11X,4HCURV
3/)
905 FORMAT(17,3X,2F10.4,F15.5,10X,2F10.4,F15.5)
906 FORMAT(/7X,12HPROFILES AT,10F10.5/)
914 FORMAT(25X,29HPRIMARY STAGNATION PRESSURE =,E12.5,7H PASCAL/
1 25X,26HPRIMARY STAGNATION TEMP. =,F7.2,15H DEGREES KELVIN/
2 25X,31HSECONDARY STAGNATION PRESSURE =,E12.5,7H PASCAL/
3 25X,28HSECONDARY STAGNATION TEMP. =,F7.2,15H DEGREES KELVI
4N// 25X,24HPRIMARY MASS FLOW RATE =,F8.5, 9H KG/SEC-M/
5 25X,44HTOTAL OF SECONDARY PLUS TERTIARY FLOW RATE =,F8.5,
69H KG/SEC-M //25X,19HNOZZLE SLOT WIDTH =,F9.6,7H METERS/)
915 FORMAT(25X,21HNOZZLE X COORDINATE =,F10.6,7H METERS/
1 25X,21HNOZZLE Y COORDINATE =,F10.6,7H METERS/
2 25X,14HNOZZLE ANGLE =,F10.3,8H DEGREES/)
XNOZ= XNOZ/D
YNOZ= YNOZ/D
TNOZ= TNOZ / 180.0 * 3.1416
DO 601 J=1,NPAIR1
XW(J,1)= XW(J,1) / D
YW(J,1)= YW(J,1) / D
CURV(J,1)=CURV(J,1)*D
601 CONTINUE
DO 602 J=1,NPAIR2
XW(J,2)= XW(J,2) / D
YW(J,2)= YW(J,2) / D
CURV(J,2)=CURV(J,2)*D
602 CONTINUE
RHOO1=A1*P01/RG/T01
UREFP=SQRT( GC*RG*T01)
C
C K.IFT IS NUMBR OF POINTS CLOSE TO JET NOZZLE WALLS

```

e  
e

KWALL IS NUMBER OF POINTS CLOSE TO DUCT WALLS

KJET= 6  
KWALL= 20  
N= K3 + K1 + K2 + 2 \* KWALL + 4 \* KJET + 1  
NN=N-1  
NX1=1  
NX2=K3+KWALL+KJET+1  
NX3=NX2  
NX4=N-K2-KWALL-KJET  
NX5=NX4  
NX6=N  
NX(1)=NX1  
NX(2)=NX2  
NX(3)=NX3  
NX(4)=NX4  
NX(5)=NX5  
NX(6)=NX6  
L1= NX(3)  
L2= NX(4)  
NMAX= (L1 + L2) / 2  
NF= 1  
P03=P02  
T03=T02  
MASS1=A1\*MASS1/RH001/UREFP/D  
MASS2=A1\*MASS2/RH001/UREFP/D

\$NA1950  
\$NA1960  
\$NA1970  
\$NA1980  
\$NA1990  
\$NA2000  
\$NA2010  
\$NA2020  
\$NA2030  
\$NA2040  
\$NA2050  
\$NA2060  
\$NA2070  
\$NA2080  
\$NA2090  
\$NA2100  
\$NA2110  
\$NA2120  
\$NA2130  
\$NA2140  
\$NA2150  
\$NA2160  
\$NA2170  
\$NA2180  
\$NA2190  
\$NA2200

c  
c  
c

SUBROUTINE INCOND IS USED TO CALCULATE THE INITIAL CONDITIONS

	CALL	INCOND(N,NN,K1,K2,K3,KWALL,KJET,L1,L2,	\$NA2210
		MASS1,MASS2,AMASS1,AMASS2,AMASS3,	\$NA2220
1		RG,GC,G,XNOZ,YNOZ,TNOZ,XW,YW,NPAIR1,NPAIR2)	\$NA2230
2		AMA1=AMASS1*UREFP*D*RHO01/A1	\$NA2240
		AMA2=AMASS2*UREFP*D*RHO01/A1	\$NA2250
		AMA3=AMASS3*UREFP*D*RHO01/A1	\$NA2260
		WRITE(6,921)	\$NA2270
		G2=G/(G-1.)	\$NA2280
		VNOZI=UREFP*SQR(2.*G2*(1.-(P02/P01)**(1./G2)))	\$NA2290
		TNOZI=AMA1*VNOZI	\$NA2300
		WRITE(6,927) TNOZI	\$NA2301
		WRITE(6,922) AMA1,AMA2,AMA3	\$NA2302
		WRITE(6,923)	\$NA2303
		DO 28 J=1,N	\$NA2304
			\$NA2310
			\$NA2320
			\$NA2330
		XS=X(J)*D	\$NA2340
		DNS=DN(J)*D	\$NA2350
		YS=Y(J)*D	\$NA2360
		TX=THETA(J,1)*57.2958	\$NA2370
		XR=R(J,1)/D	\$NA2380
28		WRITE(6,924) J,XS,YS,DNS,TX,XR,PSI(J)	\$NA2390
		WRITE(6,925)	\$NA2400
		DO 29 J=1,N	\$NA2410
		P0D=P0(J)*P01	\$NA2420
		PSD=PS(J,1)*P01	\$NA2430
		T0=TS(J,1)*T01	\$NA2440
		RH0D=RHO(J,1)*RHO01	\$NA2450
		UV=U(J,1)*UREFP	\$NA2460
29		WRITE(6, 26) J,P0D,PSD,T0,RH0D,UV,M(J)	\$NA2470
921		FORMAT(1H1,///,25X,18HINITIAL CONDITIONS//)	\$NA2480
922		FORMAT(7X,8HAMASS1 =F10.5,5X,8HAMASS2 =F10.5,5X,8HAMASS3 =F10.5/)	\$NA2490
923		FORMAT(7X,1HJ,9X,1HX,9X,1HY,8X,2HDN,6X,5HTHETA,8X,1HK,7X,3HOMG/)	\$NA2500
924		FORMAT(5X,13,2X,4F10.3,F10.6,F10.4)	\$NA2510
925		FORMAT(///7X,1HJ,8X,2HP0,8X,2HPS,8X,2HTS,7X,3HRHO,9X,1HU,9X,1HM/)	\$NA2520
926		FORMAT(5X,13,2X,2E10.3,F10.1,F10.5,F10.1,F10.3)	\$NA2530
927		FORMAT(7X,41HISENTROPIC NOZZLE THRUST PER UNIT WIDTH =,E10.3)	\$NA2531
		DO 30 J=1,N	\$NA2540
		THETA(J,2)=THETA(J,1)	\$NA2550
		PS(J,2)=PS(J,1)	\$NA2560
		TS(J,2)=TS(J,1)	\$NA2570
		RHO(J,2)=RHO(J,1)	\$NA2580
		U(J,2)=U(J,1)	\$NA2590
		R(J,2)=R(J,1)	\$NA2600

30	E(J) = 0.	\$NA2610
	DEL1=0.0	\$NA2620
	DEL2=0.0	\$NA2630
	DEL3=0.	\$NA2640
	DEL4=0.	\$NA2650
	CORE3=DN(NX2)-DN(NX1)	\$NA2660
	CORE1=DN(NX4)-DN(NX3)	\$NA2670
	CORE2=DN(NX6)-DN(NX5)	\$NA2680
	DPDSA= -0.001	\$NA2690
	USTARI = 0.0	\$NA2700
	USTARO = 0.0	\$NA2710
	UGO=0.	\$NA2720
	DGI=0.0	\$NA2730
	NGRID=0	\$NA2740
	NPRINT= 0	\$NA2750
	DSA= .02 * DN(N)	\$NA2760
	CC= .07	\$NA2770
	NCORE=1	\$NA2780
	NSEP=0	\$NA2790
	RDECAY=4.	\$NA2800
	ICOUNT=0	\$NA2810
C		\$NA2820
C	-----	\$NA2830
C	BEGINNING OF MARCHING CALCULATION	\$NA2840
C		\$NA2850
C	-----	\$NA2860
C		\$NA2870
	DO 4500 NSTEP = 1,100	\$NA2880
	ICHS = 0	\$NA2890
C		\$NA2900
C	FIND MIDDLE OF DUCT	\$NA2910
C		\$NA2920
	DO 527 J=1,N	\$NA2930
	IF (DN(J) .LT. DN(N)*0.5) NDUCT= J	\$NA2940
527	CONTINUE	\$NA2950
	IF (NSTEP .EQ. 1) GO TO 77	\$NA2960
		\$NA2970
C		\$NA2980
C	SUBROUTINE EDDY CALCULATES THE EDDY VISCOSITY FOR SUB. SOLV	\$NA2990
C		\$NA3000
	CALL EDDY(N,NN,NX,U,PS,CC,DS,E,E1,RHO,VIS,R,DN,	\$NA3010
1	S4, SS, DEL1,DEL2,DEL3,DEL4,CORE1,CORE2,CORE3)	\$NA3020
77	CONTINUE	\$NA3030
	E1(1) = 0.	\$NA3040
	DO 40 J=2,NN	\$NA3050
	E1(J) = RHO(J,1)*U(J,1)	\$NA3060
40	CONTINUE	\$NA3070
	E1(N)=0.	\$NA3080
C		\$NA3090
C	MOVE TO NEXT POINT ON WALL.....	\$NA3100
C		\$NA3110
	DS(NMAX)= DSA * (1.04) ** (NSTEP-1)	\$NA3120
	IF (DS(NMAX) .LT. .02 * DN(N) ) DS(NMAX)= .02 * DN(N)	\$NA3130
C		\$NA3140
C	CALCULATE DS(I) VALUES	\$NA3150
C	(MIDDLE DS VALUE CALCULATED FIRST	\$NA3160
C	THEN DS VALUES CALCULATED OUT TO BOTH WALLS	\$NA3170
C		\$NA3180
	NPRR=NMAX+1	\$NA3190
	DO 50 J=NPRR,N	\$NA3200
	JM= J-1	\$NA3210
	C1=RHO(J,2)+RHO(JM,2)	\$NA3220
	C2=U(J,2)+U(JM,2)	\$NA3230
	C3=R(J,2)+R(JM,2)	\$NA3240
	C4= PSI(J) - PSI(J-1)	\$NA3250
	C5 = C4*C3/(C1*C2)	\$NA3260

	DS(J) = (1.+C5)/(1.-C5)*DS(JM)	\$NA3270
50	CONTINUE	\$NA3280
	NPRR=NMAX-1	\$NA3290
	DO 52 MM=1,NPRR	\$NA3300
	J= NMAX - MM	\$NA3310
	JM= J+1	\$NA3320
	C1=RHO(J,2)+RHO(JM,2)	\$NA3330
	C2=U(J,2)+U(JM,2)	\$NA3340
	C3= R(J,2) + R(JM,2)	\$NA3350
	C4= PSI(J) - PSI(JM)	\$NA3360
	C5 = C4*C3/(C1*C2)	\$NA3370
	DS(J) = (1.+C5)/(1.-C5)*DS(JM)	\$NA3380
52	CONTINUE	\$NA3390
C		\$NA3400
C	MOVE TO NEW POINT X(1) AND X(N) ALONG WALL SURFACES	\$NA3410
C	AT A DISTANCE SPECIFIED BY DS VALUES	\$NA3420
C		\$NA3430
	DX1= X(1)	\$NA3440
	CALL DSMOVE(XW,YW,X(1),Y(1),DS(1),THETA(1,1), R(1,1), NPAIR1,1)	\$NA3450
	CALL DSMOVE(XW,YW,X(N),Y(N), DS(N), THETA(N,1), R(N,1), NPAIR2,2)	\$NA3460
	IF(X(1).GT. XW(NPAIR1 ,1)) STOP	\$NA3470
	IF(X(N).GT. XW(NPAIR2 ,2)) STOP	\$NA3480
	DX1= X(1) - DX1	\$NA3490
C		\$NA3500
C	COMPUTE Y, THETA, AND R VALUES CORRESPONDING TO NEW X VALUES	\$NA3510
C		\$NA3520
	CALL WALLS(XW,YW, X(1), Y(1), THETA1, R1, NPAIR1, 1)	\$NA3530
	CALL WALLS(XW,YW, X(N), Y(N), THETA2, R2, NPAIR2, 2)	\$NA3540
	Y2= Y(N)	\$NA3550
C		\$NA3560
C	COMPUTE CURVATURE R(J,2)	\$NA3570
C		\$NA3580
	DO 400 J=1,N	\$NA3590
	DW1= DN(J)	\$NA3600
	DW2= DN(N) - DN(J)	\$NA3610
	R(J,2)=R1*(1.-DW1/DN(N))*EXP(-RDECAY*DW1*ABS(R1))+R2*DW1/DN(N)*	\$NA3620
1	EXP(-RDECAY*DW2*ABS(R2))	\$NA3630
400	CONTINUE	\$NA3640
C		\$NA3650
C	COMPUTE THETA(J,2)	\$NA3660
C		\$NA3670
	DO 51 J=2,NN	\$NA3680
	THETA(J,2)= THETA1 + (THETA2-THETA1) * DN(J)/DN(N)	\$NA3690
51	CONTINUE	\$NA3700
	THETA(1,2)= THETA1	\$NA3710
	THETA(N,2)= THETA2	\$NA3720
C		\$NA3730
C	-----	\$NA3740
C		\$NA3750
C	THIS SECTION ATTEMPTS TO SATISFY CONTINUITY	\$NA3760
C	LOOKS FOR A PS(1,2) SUCH THAT Y2 - Y(N) = 0.0	\$NA3770
C	DPDSA, DPDSB ARE PRESSURE GRADIENTS	\$NA3780
C		\$NA3790
C	-----	\$NA3800
C		\$NA3810
	IF(ABS(DPDSA) .LT. 1.E-08) DPDSA=-.001	\$NA3820
	DPDSB= DPDSA * 0.9	\$NA3830
	DDISTB= 1.0	\$NA3840
	DO 650 ITER= 1,40	\$NA3850
	IF (ITER .EQ. 1) PS(1,2)= PS(1,1) + DPDSA	\$NA3860
C		\$NA3870
C	AT THIRD ITERATION GUESS NEW DPDSB	\$NA3880
C	USING EXTRAPOLATION THROUGH DPDSA AND PREVIOUS DPDSB	\$NA3890
	IF (ITER .GE. 3)	\$NA3900
1	DPDSB= (DDISTB*DPDSA - DDISTA*DPDSB) / (DDISTB-DDISTA)	\$NA3910
	IF (ITER .GE. 2) PS(1,2)= PS(1,1) + DPDSB	\$NA3920

	DO 60 J=2,N	\$NA3930
	JM= J-1	\$NA3940
	C2=U(J,2)+U(JM,2)	\$NA3950
	C3=R(J,2)+R(JM,2)	\$NA3960
	C4= PSI(J) - PSI(J-1)	\$NA3970
	PS(J,2)=PS(JM,2)+ C2*C3*C4/4.0	\$NA3980
60	CONTINUE	\$NA3990
C		\$NA4000
C	SUBROUTINE SOLV IS USED TO SOLVE U,T ON M=2 LINE	\$NA4010
C		\$NA4020
	CALL SOLV(DS,N,NN,PR1,PRT1,G)	\$NA4030
	U(1,2) = 0.	\$NA4040
	DO 70 J = 2,NN	\$NA4050
	U(J,2) = H(J-1)	\$NA4060
70	CONTINUE	\$NA4070
	U(N,2) = 0.	\$NA4080
	CALL SOLV(DS,N,NN,PR,PRT,G)	\$NA4090
	DO 80 J = 2,NN	\$NA4100
	TS(J,2) = H(J-1)	\$NA4110
80	CONTINUE	\$NA4120
	TS(1,2) = TS(2,2)	\$NA4130
	TS(N,2) = TS(N-1,2)	\$NA4140
	DO 90 I = 1,N	\$NA4150
	RHO(I,2) = PS(I,2)/TS(I,2)	\$NA4160
	RHO(I,1) =RHO(I,2)	\$NA4170
	VIS(I)=(TREF+A2)/(T01*TS(I,2)+A2)	\$NA4180
	VIS(I)=VIS(I)*(T01*TS(I,2)/TREF)**0.5	\$NA4190
	VIS(I)=VIS(I)*RHOREF*NUREF/(P01*SQRT(GC/RG/T01))	\$NA4200
90	CONTINUE	\$NA4210
C		\$NA4220
C	GIVEN PS(1,2), CALCULATE UPPER WALL COORDINATE Y(N)	\$NA4230
C		\$NA4240
	CALL CORDS(N,X,Y,THETA,U,RHO,PSI,DN,NDUCT,DDIST)	\$NA4250
C		\$NA4260
C	IF Y2 .EQ. Y(N) THEN THIS PS(1,2) SATISFIES CONTINUITY	\$NA4270
C		\$NA4280
	IF (ITER .EQ. 1) DDISTA= DDIST	\$NA4290
	IF (ITER .GE. 2) DDISTB= DDIST	\$NA4300
C		\$NA4310
C	IF SUFFICIENTLY SMALL INTERVAL, THEN EXIT	\$NA4320
C		\$NA4330
	IF (ABS(DDISTB) .LE. 0.1 ** 4) GO TO 660	\$NA4340
650	CONTINUE	\$NA4350
	WRITE(6,916) DPDSA,DPDSB,DDISTB,NSTEP	\$NA4360
916	FORMAT(/7X,23HNO CONVERGENCE, DPDSA =,F10.5,8H DPDSB =,F10.5,	\$NA4370
1	9H DDISTB =,F10.5,6H STEP=I4)	\$NA4380
660	CONTINUE	\$NA4390
C		\$NA4400
C	SAVE PRESSURE GRADIENT THIS STEP	\$NA4410
C		\$NA4420
	DPDSA= DPDSB	\$NA4430
	G1=G-1.	\$NA4440
	G2=G/G1	\$NA4450
C		\$NA4460
C	COMPUTE MACH NUMBERS M(I)	\$NA4470
C		\$NA4480
	DO 220 I=1,N	\$NA4490
	M(I)=U(I,2)/SQRT(G*TS(I,2))	\$NA4500
	G4 = 1. + G1/2.*(M(I)*M(I))	\$NA4510
	P0(I) = PS(I,2)*G4**G2	\$NA4520
220	CONTINUE	\$NA4530
C		\$NA4540
C	-----	\$NA4550
C	REMOVE EXTRA GRID POINTS, FIND EDGES OF FLOWS	\$NA4560
C	-----	\$NA4570



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CALL REMOVE(N,NN,NGRID,NX,NMAX,PS1,S1,S2,S3,S4,S5,DS,VIS,DN,      $NA4590
1 X,Y,M,P0,H,THETA,R,PS,U,TS,RHO)                                $NA4600
CALL LOOK(N,NN,K3,NCORE,DEL1,DEL2,DEL3,DEL4,CORE1,CORE2,CORE3,    $NA4610
1 NX,P0,DN,U,DIP,NMAX)                                           $NA4620
SHRIN=U(2,2)*VIS(2)/(DN(2)-DN(1)) - .5*(PS(1,2)-PS(1,1))/DS(1)* $NA4630
1 (DN(2)-DN(1))                                                    $NA4640
SHROR=U(NN,2)*VIS(NN)/(DN(N)-DN(NN)) -.5*(PS(N,2)-PS(N,1))/    $NA4650
1 DS(N)*(DN(N)-DN(NN))                                             $NA4660
IF((SHRIN.LE.0.).OR.(SHROR.LE.0.)) NSEP=1                          $NA4670
IF(NSEP.EQ.1) NPRINT=NPCYCL-1                                     $NA4680
IF(NSEP.EQ.1) ICOUNT=1                                            $NA4690
C-----                                                         $NA4700
C-----                                                         $NA4710
C SAVE VALUES COMPUTED THIS STEP                                $NA4720
C-----                                                         $NA4730
C-----                                                         $NA4740
C DO 222 I=1,N                                                    $NA4750
C PS(I,1) = PS(I,2)                                                $NA4760
C U(I,1) = U(I,2)                                                  $NA4770
C TS(I,1) = TS(I,2)                                                $NA4780
C THETA(I,1)=THETA(I,2)                                           $NA4790
222 R(I,1) = R(I,2)                                                $NA4800
C-----                                                         $NA4810
C-----                                                         $NA4820
C OUTPUT SECTION                                                  $NA4830
C-----                                                         $NA4840
C-----                                                         $NA4850
DS1Q= DS(1) * D                                                    $NA4860
DS2Q= DS(N) * D                                                    $NA4870
DSUM1= DSUM1 +DS1Q                                                 $NA4880
DSUM2= DS2Q + DSUM2                                               $NA4890
NPRINT= NPRINT+1                                                  $NA4900
IF ( (X(1)*D) .GE. XPROF(NF) ) NPRINT= NPCYCL                    $NA4910
IF(X(1)*D.GE.XPROF(NF)) ICOUNT=1                                  $NA4920
IF ( (X(1)*D) .GE. XPROF(NF) ) NF= NF + 1                         $NA4930
IF (NSTEP .EQ. 1) NPRINT= NPCYCL                                  $NA4940
IF (NPRINT .LT. NPCYCL) GO TO 45                                  $NA4950
NPRINT= 0                                                         $NA4960
DXQ= DX1 * D                                                       $NA4970
XI= X(1) * D                                                       $NA4980
XQ= X(N) * D                                                       $NA4990
PH201=(PS(1,2)*P01-P02)/A3                                         $NA5000
PH200=(PS(N,2)*P01-P02)/A3                                         $NA5010
IF (SHRIN .GT. 0.0)                                                $NA5020
1 USTARI=SQRT(SHRIN/RHO(2,2))*UREFP                                $NA5030
IF (SHROR .GT. 0.0)                                                $NA5040
1 USTARO=SQRT(SHROR/RHO(NN,2))*UREFP                                $NA5050
RI= R(1,2) / D                                                     $NA5060
RO= R(N,2) / D                                                     $NA5070
RNI= 2.0 * (R(2,1) / RHO(2,1))                                     $NA5080
RNI=RNI/(S5(2)*(U( 3,1)-U(2,1))+S4(2)*(U(2,1)-U(1,1)))          $NA5090
RNO= 2.0 * (R(NN,1) / RHO(NN,1))                                   $NA5100
RNO=RNO/(S5(NN)*(U( N,1)-U(NN,1))+S4(NN)*(U(NN,1)-U(NN-1,1)))   $NA5110
UMAX= U(NMAX,2) * UREFP                                           $NA5120
XTS1= TS(1,2) * T01 + 0.5/G2 * T01 * (U(1,2))**2              $NA5130
XTS0= TS(N,2) * T01 + 0.5/G2 * T01 * (U(N,2))**2              $NA5140
WRITE(6,908) NSTEP,XI,PH201,USTARI,RI,RNI,XQ,PH200,USTARO,RO,RNO, $NA5150
1 UMAX,ITER                                                         $NA5160
908 FORMAT(1H1,5X,5HNSTEP,6X,2HXI,8X,2HP1,4X,6HUSTARI,8X,2HKI,7X, $NA5170
1 3HRNI,8X,2HXO,8X,2HPO,4X,6HUSTARO,8X,2HKO,7X,3HRNO,6X,4HUMAX/ $NA5180
2 5X,13,2(F10.4,E11.3,2F10.4,F10.5),F10.1//7X,                  $NA5190
3 22HNUMBER OF ITERATIONS =,I3/)                                  $NA5200
WRITE(6,907)DXQ,DS1Q,DS2Q,DSUM1,DSUM2                            $NA5210
907 FORMAT( 7X, 9HDELTA X =,F10.5,12H DELTA SI =,F10.5,          $NA5220
1 12H DELTA SO =,F10.5,6H SI =F10.5,6H SO =,F10.5/)             $NA5230
SHRIN=SHRIN*P01                                                    $NA5240

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	SHROR=SHROR*P01	\$NA5250
	WRITE(6,912) SHRIN,SHROR	\$NA5260
	TMOMX=0.	\$NA5261
	NNY=N-1	\$NA5262
	DO 95 J=1,NNY	\$NA5263
95	TMOMX=TMOMX+(RHO(J,1)*U(J,1)**2*COS(THETA(J,1))+RHO(J+1,1)*	\$NA5264
1	U(J+1,1)**2*COS(THETA(J+1,1)))/2.*(Y(J+1)-Y(J))	\$NA5265
	TMOMX=TMOMX*P01*D	\$NA5266
	CT=TMOMX/TNOZI	\$NA5267
	WRITE(6,931) TMOMX,CT	\$NA5268
912	FORMAT( 7X,7HSHRIN =E12.5,5X,7HSHROR =E12.5/)	\$NA5270
931	FORMAT(/7X,37HTOTAL AXIAL MOMENTUM PER UNIT WIDTH =,E10.3//7X,	\$NA5271
1	14HTHRUST RATIO =,E10.3)	\$NA5272
	IF (NQUICK.GT. 0) GO TO 45	\$NA5280
	IF(ICOUNT.EQ.0) GO TO 45	\$NA5290
	ICOUNT=0	\$NA5300
	CORE1Q= CORE1 * D	\$NA5310
	CORE2Q= CORE2 * D	\$NA5320
	CORE3Q= CORE3 * D	\$NA5330
	DEL1Q= DEL1 * D	\$NA5340
	DEL2Q= DEL2 * D	\$NA5350
	DEL3Q= DEL3 * D	\$NA5360
	DEL4Q= DEL4 * D	\$NA5370
	WRITE(6,909) (NX(I),I=1,6),CORE1Q,CORE2Q,CORE3Q,DEL1Q,DEL2Q,	\$NA5380
1	DEL3Q,DEL4Q	\$NA5390
909	FORMAT(7X,2HN1,5X,2HN2,5X,2HN3,5X,2HN4,5X,2HN5,5X,2HN6/5X,I4,	\$NA5400
1	5(3X,I4)// 7X,5HCORE1,5X,5HCORE2,5X,5HCORE3,/4X,3F10.5//	\$NA5410
2	7X,4HDEL1,6X,4HDEL2,6X,4HDEL3,6X,4HDEL4/4X,4F10.5//)	\$NA5420
	WRITE(6,910)	\$NA5430
910	FORMAT(7X,14J,7X,3H1/R,5X,5HTHETA,9X,1HX,9X,1HY,6X,4HYREL,6X,	\$NA5440
1	4HUREL,5X,5HPOREL,8X,2HPS,9X,1HE,5X,6HTOTEMP/)	\$NA5450
	DO 100 J=1,N	\$NA5460
	THEIO=THETA(J,2)*180./3.1416	\$NA5470
	XR=R(J,2)/D	\$NA5480
	YS=Y(J)*D	\$NA5490
	YREL= Y(J) / Y(N)	\$NA5500
	XS= X(J)*D	\$NA5510
	UV=U(J,2) / U(NMAX,2)	\$NA5520
	P0D=P0(J)	\$NA5530
	PSD=PS(J,2)*P01	\$NA5540
	T0= TS(J,2) * T01 + 0.5/G2 * T01 * (U(J,2))**2	\$NA5550
	T0MAX= TS(NMAX,2) * T01 + 0.5/G2 * T01 * (U(NMAX,2))**2	\$NA5560
	XTS= (T0 - XTS0) / (T0MAX - XTS0)	\$NA5570
	XT=E(J)*UREFP*D/SQRT(A1)	\$NA5580
	DND=DN(J)*D	\$NA5590
	WRITE(6,911)J,XR,THEIO,XS,YS,YREL,UV,P0D,PSD,XT,T0	\$NA5600
911	FORMAT(5X,I3,2X,F10.5,3F10.4,3F10.5,E11.3,F10.5,F10.2)	\$NA5610
100	CONTINUE	\$NA5620
45	CONTINUE	\$NA5630
	IF(NSEP.EQ.1) GO TO 300	\$NA5640
4499	CONTINUE	\$NA5650
4500	CONTINUE	\$NA5660
	STOP	\$NA5670
300	WRITE(6,930)	\$NA5680
930	FORMAT(7X,24(1H*)/7X,38HNEGATIVE SHEAR STRESS AT THIS STATION./	\$NA5690
1	7X,36H THIS INDICATES POSSIBLE SEPARATION./	\$NA5700
2	7X,48H SUBSEQUENT RESULTS SHOULD BE USED WITH CAUTION.//	\$NA5710
3	7X,24(1H*))	\$NA5720
	NSEP=0	\$NA5730
	GO TO 4499	\$NA5740
	END	\$NA5750

	SUBROUTINE INCOND(N,NN,K1,K2,K3,KWALL,KJEL,L1,L2,	\$IN	0
1	MASS1,MASS2,AMASS1,AMASS2,AMASS3,	\$IN	10
2	RG,GC,G,XNOZ,YNOZ,TNOZ,XW,YW,NPAIR1,NPAIR2)	\$IN	20
C		\$IN	30
C	*****	\$IN	40
C	SUBROUTINE INCOND	\$IN	50
C	CALCULATES INITIAL CONDITIONS	\$IN	60
C	INPUT VARIABLES	\$IN	70
C	NPAIR1: NUMBER OF POINTS ON LOWER WALL	\$IN	80
C	NPAIR2: NUMBER OF POINTS ON UPPER WALL	\$IN	90
C	XW(I,1), YW(I,1): DATA POINTS SPECIFYING LOWER WALL	\$IN	100
C	XW(I,2), YW(I,2): DATA POINTS SPECIFYING UPPER WALL	\$IN	110
C	X1: STARTING POINT ON LOWER WALL	\$IN	120
C	XC,YC: CENTRE OF INITIAL RADIUS OF CURVATURE	\$IN	130
C	*****	\$IN	140
C		\$IN	150
	REAL MASS1,MASS2,NUD,M(190)	\$IN	160
	DIMENSION XW(99,2),YW(99,2),S1(190),S2(190),S3(190), S4(190),	\$IN	170
1	S5(190),E(190),E1(190),DS(190),X(190),Y(190),	\$IN	180
2	P0(190),H(190),THETA(190,2),R(190,2),U(190,2),	\$IN	190
3	TS(190,2),RHO(190,2),VIS(190)	\$IN	200
	DIMENSION OMG(190),DN(190),PS(190,2)	\$IN	210
C	DOUBLE OMG(190),DN(190),PS(190,2)	\$IN	220
	DIMENSION PSI(190)	\$IN	230
C	DOUBLE PSI(190)	\$IN	240
C	DOUBLE XT,P3,RY,PP,DY,PSI,DPSI,ZY,DSQRT,DPS,Z	\$IN	250
	COMMON/SOLV/ S1,S2,S3,E,E1,H,S4,S5	\$IN	260
	COMMON/INEDY/PSI,RHO,VIS	\$IN	270
	COMMON/INCD/ U,PS,TS,M,P0,THETA,R,X,Y,DN	\$IN	280
	COMMON /CONST/ P01,P02,P03, T01,T02,T03, X1, XC, YC	\$IN	290
	DIST(XX,YY,XXC,YYC)=SQRT((ABS(XX-XXC)**2)+(ABS(YY-YYC)**2))	\$IN	300
C		\$IN	310
C	LOOK FOR ARC TANGENT TO BOTH WALLS WHICH PASSES THROUGH NOZZLE	\$IN	320
C	X1 AND X2 ARE STARTING POINTS ON LOWER AND UPPER DUCT WALL	\$IN	330
C		\$IN	340
C	XNOZ, YNOZ IS MIDDLE OF NOZZLE	\$IN	350
C	FIND COORDS OF EACH EDGE	\$IN	360
C	NOZZLE WIDTH IS 1.0	\$IN	370
C		\$IN	380
	XNOZ1= XNOZ + ( (SIN(TNOZ)) * 0.5)	\$IN	390
	XNOZ2= XNOZ - ( (SIN(TNOZ)) * 0.5)	\$IN	400
	YNOZ1= YNOZ - (ABS(COS(TNOZ)) * 0.5)	\$IN	410
	YNOZ2= YNOZ + (ABS(COS(TNOZ)) * 0.5)	\$IN	420
C		\$IN	430
C	COMPUTE ARC LENGTH ACROSS DUCT (DN2)	\$IN	440
C	ARC LENGTHS TO NOZZLE (DNI, DNO)	\$IN	450
C		\$IN	460
	CALL ARCDIS(XNOZ1,YNOZ1,TNOZ,X1,RC1,XW,YW,NPAIR1,1)	\$IN	470
	CALL WALLS(XW,YW,X1,Y1,T1,R1,NPAIR1,1)	\$IN	480
	IF (RC1 .LT. 100.0) DNI= RC1 * ABS(TNOZ-T1)	\$IN	490
	IF (RC1 .GE. 100.0) DNI= DIST(XNOZ1,YNOZ1,X1,Y1)	\$IN	500
	CALL ARCDIS(XNOZ2,YNOZ2,TNOZ,X2,RC2,XW,YW,NPAIR2,2)	\$IN	510
	CALL WALLS(XW,YW,X2,Y2,T2,R2,NPAIR2,2)	\$IN	520
	IF (RC2 .LT. 100.0) DNO= RC2 * ABS(TNOZ-T2)	\$IN	530
	IF (RC2 .GE. 100.0) DNO= DIST(XNOZ2,YNOZ2,X2,Y2)	\$IN	540
	DN2= DNI + DNO +1.0	\$IN	550
	X(1)=X1	\$IN	560
	Y(1)=Y1	\$IN	570
	X(N)= X2	\$IN	580
	Y(N)= Y2	\$IN	590
	THETA(1,1) = T1	\$IN	600
	THETA(N,1) = T2	\$IN	610
		\$IN	620

C	SOLVE FOR MASS SPLIT (NOZZLE EXIT CORRESPONDS TO DN(L1) )	\$IN 630
C	DUCT COMPUTED MUST CORRESPOND TO ACTUAL DUCT WIDTH	\$IN 640
C	THUS DN(1) SHOULD EQUAL 0.0	\$IN 650
C		\$IN 660
C	FIND LOCATION OF NOZZLE	\$IN 670
C	(DIFFERENT CALCULATIONS FOR SPLMAX DEPENDING ON	\$IN 680
C	NOZZLE LOCATION)	\$IN 690
C		\$IN 700
	IF (DNI .LT. DNO) NZONE= 0	\$IN 710
	IF (DNI .GE. DNO) NZONE= 1	\$IN 720
	IF (NZONE .EQ. 0) NMID= L2	\$IN 730
	IF (NZONE .EQ. 1) NMID= L1	\$IN 740
	G1= 1.0 / (G-1.0)	\$IN 750
	G2= G/(G-1.0)	\$IN 760
	RDECAY=4.	\$IN 770
	A=DNI+1.+DNO	\$IN 780
	IF (NZONE.EQ.0) RTEST=R1*(1.-(DNI+1.)/A)*EXP(-RDECAY*(DNI+1.))*	\$IN 790
1	ABS(R1))+R2*(DNI+1.)/A*EXP(-RDECAY*DNO*ABS(DNO))	\$IN 800
	IF (NZONE.EQ.1) RTEST=R1*(1.-DNI/A)*EXP(-RDECAY*DNI*ABS(R1))+	\$IN 810
1	R2*DNI/A*EXP(-RDECAY*(DNO+1.)*ABS(R2))	\$IN 820
	TMAX= (T02/T01 - .0001*((G-1.0)/(G*2.0)) )	\$IN 830
1	* (P03/P01) ** ((G-1.0)/G) * (T01/T02)	\$IN 840
	M(NMID)= SQRT(2.0 * G1 * (1.0/ TMAX - 1.0) )	\$IN 850
	PSNMID= TMAX ** G2	\$IN 860
	PHO(NMID,1)= P02 / P01 * T01 / T02	\$IN 870
	U2MIN=ABS(.0001-2.*G*PSNMID*(M(NMID))**2*RTEST/RHO(NMID,1))	\$IN 880
	U2MIN= SQRT(U2MIN)	\$IN 890
	SPLMAX= 1.0 - ( P02 / P01 * T01 / T02 * U2MIN * DNO / MASS2)	\$IN 900
	SPLMAX= ABS(SPLMAX)	\$IN 910
	IF (SPLMAX .GT. 0.99) SPLMAX= 0.99	\$IN 920
C		\$IN 930
C	THIS IS ENTRY POINT FOR LOWERING SPLMAX IF TMAX IS EXCEEDED	\$IN 940
C		\$IN 950
555	CONTINUE	\$IN 960
	SPLMIN= 1.0 - SPLMAX	\$IN 970
	SPLTST= DNI / (DNI + DNO)	\$IN 980
	XA= SPLTST	\$IN 990
	IF ((NZONE .EQ. 0).AND.(XA .LT. SPLMIN*1.1)) XA= SPLMIN*1.11	\$IN1000
	IF ((NZONE .EQ. 1).AND.(XA .GT. SPLMAX* 1.0)) XA= SPLMAX* 1.0	\$IN1010
	XB= XA * 0.9	\$IN1020
	DIFFB= 1.0	\$IN1030
	NMID= (L1 + L2) / 2	\$IN1040
	DO 40 ITER= 1.50	\$IN1050
	IF (ITER .EQ. 1) SPLIT=XA	\$IN1060
	IF (ITER .GE. 3) XB= (DIFFB*XA - DIFFA*XB) / (DIFFB-DIFFA)	\$IN1070
	IF (ITER .GE. 2) SPLIT= XB	\$IN1080
	AMASS1=MASS1	\$IN1090
	AMASS3= SPLIT * MASS2	\$IN1100
	AMASS2= (1.0-SPLIT) * MASS2	\$IN1110
C		\$IN1120
C	STREAM FUNCTIONS	\$IN1130
C		\$IN1140
	OMGN= AMASS1 + AMASS2 + AMASS3	\$IN1150
	OMGJ3= AMASS3	\$IN1160
	OMGJ2= AMASS3 + AMASS1	\$IN1170
C		\$IN1180
C	DMJET IS THE SMALL DELTA OMG AT JET WALLS	\$IN1190
C	DMWALL IS THE SMALL DELTA OMG STARTING IN FROM DUCT WALLS	\$IN1200
C		\$IN1210
	DMWALL= .00001 * OMGN	\$IN1220
	DMJET= DMWALL * 100.0	\$IN1230
	DOMG3= OMGJ3 / FLOAT(K3)	\$IN1240
	DOMG2= (OMGN - OMGJ2) / FLOAT(K2)	\$IN1250

	D0MG1= (OMGJ2 - OMGJ3) / FLOAT(K1)	\$IN1260
	PSI(1)= 0.0	\$IN1270
	PSI(L1)= OMGJ3	\$IN1280
	PSI(L2)= OMGJ2	\$IN1290
	PSI(N)= OMGN	\$IN1300
	CALL OMGSET(PSI, 1,L1,KWALL,KJET,DMWALL,DMJET,D0MG3,RST3,RFIN3)	\$IN1310
	CALL OMGSET(PSI,L1,L2,KJET,KJET,DMJET,DMJET,D0MG1,RST1,RFIN1)	\$IN1320
	CALL OMGSET(PSI,L2,N,KJET,KWALL,DMJET,DMWALL,D0MG2,RST2,RFIN2)	\$IN1330
	CALL TMSFT(N,NN,RG,GC,G,DN2,R1,R2,L1,L2,DNI,DNO,	\$IN1340
1	PSI, R, M, TS, PS, P0, RHO, U, DN,MAXERR)	\$IN1350
	IF ((MAXERR .EQ. 1).AND.(ITER .EQ. 1)) GO TO 589	\$IN1360
	IF ((MAXERR .EQ. 1).AND.(ITER .NE. 1)) GO TO 599	\$IN1370
	IF (ITER .EQ. 1) DIFFA= DN(1)	\$IN1380
	IF (ITER .GE. 2) DIFFB= DN(1)	\$IN1390
C		\$IN1400
C	CHECK FOR TERMINATION	\$IN1410
C		\$IN1420
	IF (ABS(DIFFB) .LE. 0.001) GO TO 42	\$IN1430
40	CONTINUE	\$IN1440
42	WRITE(6,902)SPLIT,ITER	\$IN1450
902	FORMAT(/7X,7HSPLIT =,F10.5/	\$IN1460
1	7X,36HNUMBER OF ITERATIONS TO FIND SPLIT =,I4//)	\$IN1470
C		\$IN1480
C	COMPUTE OMG DIFFERENCE ARRAYS S1... S5	\$IN1490
C		\$IN1500
	DO 20 J= 2,NN	\$IN1510
	JP = J+1	\$IN1520
	JM = J-1	\$IN1530
	S1(J)= PSI(JP) - PSI(JM)	\$IN1540
	S2(J)= PSI(JP) - PSI(J)	\$IN1550
	S3(J)= PSI(J) - PSI(JM)	\$IN1560
	S4(J) = S2(J)/S3(J)/S1(J)	\$IN1570
	S5(J) = S3(J)/S2(J)/S1(J)	\$IN1580
20	CONTINUE	\$IN1590
C		\$IN1600
C	COMPUTE ANGLES THETA (I)	\$IN1610
C		\$IN1620
	DO 410 J=2,NN	\$IN1630
	IF ( (J .GE. L1).AND.(J .LE. L2)) THETA(J,1)= TNOZ	\$IN1640
	IF (J.LT.L1) THETA(J,1)=THETA(J-1,1)-(DN(J)-DN(J-1))/RC1	\$IN1650
	IF (J.GT.L2) THETA(J,1)=THETA(J-1,1)-(DN(J)-DN(J-1))/RC2	\$IN1660
410	CONTINUE	\$IN1670
	DO 2 J=2,NN	\$IN1680
	JM=J-1	\$IN1690
	D=.5*(THETA(J,1)+THETA(JM,1))	\$IN1700
	X(J)=X(JM) -(DN(J)-DN(JM))*SIN(D)	\$IN1710
	Y(J) = Y(JM) +(DN(J)-DN(JM))*COS(D)	\$IN1720
2	CONTINUE	\$IN1730
	RETURN	\$IN1740
C		\$IN1750
C	THIS SPLMAX GIVES A TS EXCEEDING TMAX	\$IN1760
C	SO TRY LOWERING SPLMAX 1 PERCENT	\$IN1770
C		\$IN1780
589	CONTINUE	\$IN1790
	SPLMAX= SPLMAX * 0.99	\$IN1800
	GO TO 555	\$IN1810
C		\$IN1820
C	NO SOLUTION FOR SPLIT	\$IN1830
C		\$IN1840
599	CONTINUE	\$IN1850
	WRITE(6,911)	\$IN1860
911	FORMAT(/7X,21HNO SOLUTION FOR SPLIT)	\$IN1870
	STOP	\$IN1880
	END	\$IN1890

	SUBROUTINE TMPSET(N,NN,RG,GC,G,DN2,R1,R2,L1,L2,DNI,DNO,	\$TM 0
	1 PSI, R, M, TS, PS, P0, RHO, U, DN, MAXERR)	\$TM 10
C	-----	\$TM 20
C	SETS AN INITIAL TS IN MIDDLE OF NOZZLE,	\$TM 30
C	DETERMINES CORRESPONDING M(I), TS(I), DN(I), ETC VALUES	\$TM 40
C	CHECKS DIFFERENCE BETWEEN DN(N) AND DN2	\$TM 50
C	-----	\$TM 60
		\$TM 70
	REAL M(190)	\$TM 80
	DIMENSION P0(190),THETA(190,2),R(190,2),U(190,2),TS(190,2),	\$TM 90
	1 RHO(190,2)	\$TM 100
	DIMENSION OMG(190),DN(190),PS(190,2)	\$TM 110
C	DOUBLE OMG(190),DN(190),PS(190,2)	\$TM 120
	DIMENSION PSI(180)	\$TM 130
C	DOUBLE PSI(190)	\$TM 140
	COMMON /CONST/ P01,P02,P03, T01,T02,T03, X1, XC, YC	\$TM 150
C		\$TM 160
C	THE STARTING ITERATION POINT IS AT PRESENT DETERMINED BY THE	\$TM 170
C	LOCATION OF THE NOZZLE IN THE DUCT	\$TM 180
C	IF NOZZLE IS CLOSE TO UPPER WALL, THE TMAX CALCULATION IS	\$TM 190
C	VALID ONLY FOR TS(L1,1), AND SO NMID SHOULD EQUAL L1	\$TM 200
C	IF NOZZLE IN MIDDLE OF DUCT NMID SHOULD = (L1 + L2) / 2	\$TM 210
C	IF NOZZLE CLOSE TO LOWER WALL NMID SHOULD = L2	\$TM 220
C	AND CHANGE COMPUTATION OF R(J,NMID), AND DN(NMID)	\$TM 230
C		\$TM 240
	IF (DNI .LT. DNO) NZONE= 0	\$TM 250
	IF (DNI .GE. DNO) NZONE= 1	\$TM 260
	MAXERR= 0	\$TM 270
	G1=1./(G-1.)	\$TM 280
	G2=G/(G-1.)	\$TM 290
	RDECAY=4.	\$TM 300
	IF (NZONE .EQ. 0) NMID= L2	\$TM 310
	IF (NZONE .EQ. 0) R(NMID,1)= R1 * EXP(-RDECAY*R1*(DNI+1.0))	\$TM 320
	+ R2 * EXP( RDECAY*R2*DNO)	\$TM 330
1	IF (NZONE .EQ. 0) DN(NMID)= DNI + 1.0.	\$TM 340
	IF (NZONE .EQ. 1) NMID= L1	\$TM 350
	IF (NZONE .EQ. 1) R(NMID,1)= R1 * EXP(-RDECAY*R1*DNI)	\$TM 360
1	+ R2 * EXP( RDECAY*R2*(DNO+1.0))	\$TM 370
	IF (NZONE .EQ. 1) DN(NMID)= DNI	\$TM 380
C		\$TM 390
C	TMAX IS THE TEMPERATURE WHICH PREVENTS VELOCITY FROM	\$TM 400
C	FALLING BELOW .01 (IE GOING NEGATIVE)	\$TM 410
C		\$TM 420
	TMAX= (T02/T01 - .0001*((G-1.0)/(G*2.0)) )	\$TM 430
1	* (P03/P01) ** ((G-1.0)/G) * (T01/T02)	\$TM 440
595	CONTINUE	\$TM 445
	TSA=, TMAX	\$TM 450
	TSB= TSA - (5.0/T01)	\$TM 460
	DIFFB= 1.0	\$TM 470
C		\$TM 480
C	DO LOOP WHICH SOLVES FOR TSB SUCH THAT DN2-DN(N)= 0.0	\$TM 490
C	GUESS NEW TSB USING FIXED TSA AND PREVIOUS TSB VALUE	\$TM 500
C		\$TM 510
	DO 40 J=1,100	\$TM 520
	IF (J .EQ. 1) TS(NMID,1) = TSA	\$TM 530
	IF (J .GE. 3) TSB= (DIFFB*TSA-DIFFA*TSB) / (DIFFB-DIFFA)	\$TM 540
	IF (J .GE. 2) TS(NMID,1)= TSB	\$TM 550
	IF (TS(NMID,1) .GT. TMAX) GO TO 599	\$TM 560

		\$TM 570
C	SOLVE FOR TEMPRATURES,MACH NUMBERS,ETC. AT NODE POINTS	\$TM 580
C	WORKING FROM NOZZLE MIDDLE (NMID) OUT TO BOTH WALLS	\$TM 590
C		\$TM 600
C	NOZZLE CENTRE LINE VALUES	\$TM 610
C		\$TM 620
	M(NMID)= SQRT(2.0 * G1 * (1.0/TS(NMID,1) - 1.0) )	\$TM 630
	PS(NMID,1)= TS(NMID,1) ** G2	\$TM 640
	RHO(NMID,1)= PS(NMID,1) / TS(NMID,1)	\$TM 650
	U(NMID,1)= M(NMID) * SQRT(G * TS(NMID,1))	\$TM 660
	P0(NMID)= PS(NMID,1) * (1.0 + 0.5 / G1 * M(NMID)**2) **G2	\$TM 670
C		\$TM 680
C	REGION FROM NOZZLE MIDDLE TO OUTER WALL	\$TM 690
C	SET M, TS, PS	\$TM 700
C	(L2 IS JUMP FROM NOZZLE TO SECONDARY STREAM)	\$TM 710
C		\$TM 720
	PK=1.	\$TM 730
	QK=1.	\$TM 740
	RK=1.	\$TM 750
	NPRR=NMID+1	\$TM 760
	DO 42 I=NPRR,N	\$TM 770
	IM=I-1	\$TM 780
C		\$TM 790
C	JUMP POINT TO SECONDARY STREAM	\$TM 800
C		\$TM 810
	IF (I .EQ. L2+1) PK= P02/P01	\$TM 820
	IF (I .EQ. L2+1) QK=P02/P01*(T01/T02)**G2	\$TM 830
	IF (I .EQ. L2+1) RK=T02/T01	\$TM 840
	SX=RK*(IM,1)*(PSI(I)-PSI(IM))	\$TM 850
	M(I)= M(IM) - SX/(G**0.5 * QK * TS( IM,1)**(G2+0.5) )	\$TM 860
	B1= ((PK/PS( IM,1)) ** (1./G2)-1.)*2.*G1	\$TM 870
	IF((I.EQ.L2+1).AND.(B1.LE.0.0)) TMAX=TMAX*0.995	\$TM 872
	IF((I.EQ.L2+1).AND.(B1.LE.0.0)) GO TO 595	\$TM 873
	IF (I .EQ. L2+1) M(I)= SQRT(B1)	\$TM 880
C		\$TM 890
C	SET TS, PS, RHO, U, P0, DN, R ACROSS STREAM FROM M(I)	\$TM 900
C		\$TM 910
	TS(I,1)=RK/(1.+0.5/G1*(M(I))**2.)	\$TM 920
	PS(I,1)=QK*(TS(I,1))**G2	\$TM 930
	RHO(I,1)=PS(I,1)/TS(I,1)	\$TM 940
	U(I,1)=M(I)*SQRT(G*TS(I,1))	\$TM 950
	P0(I)=PS(I,1)*(1.+0.5/G1*(M(I))**2.))**G2	\$TM 960
	Z= 2.0 / (RHO(IM,1)*U(IM,1)+RHO(I,1)*U(I,1))	\$TM 970
	DN(I)= DN(IM) + Z * (PSI(I) - PSI(IM))	\$TM 980
	DW1= DN(I)	\$TM 990
	DW2= DN2 - DN(I)	\$TM1000
	IF (DW2 .LT. 0.0) DW2= 0.0	\$TM1010
42	R(I,1)=R1*(1.+DW1/DN2)*EXP(-RDECAY*DW1*ABS(R1))+R2*DW1/DN2*	\$TM1020
1	EXP(-RDECAY*DW2*ABS(R2))	\$TM1030
C		\$TM1040
C	SEARCH FOR TSA WHERE DN2-DN(N) .EQ. 0.0	\$TM1050
C		\$TM1060
	IF (J .EQ. 1) DIFFA= DN2-DN(N)	\$TM1070
	IF (J .GE. 2) DIFFB= DN2-DN(N)	\$TM1080
C		\$TM1090
C	IF SUFFICIENTLY SMALL INTERVAL, THEN EXIT	\$TM1100
C		\$TM1110
	IF (ABS(TSB-TSA) .LE. 0.0000001) GO TO 41	\$TM1120
	IF (ABS(DIFFB) .LE. 0.0001) GO TO 41	\$TM1130
40	CONTINUE	\$TM1140
	WRITE(6,900) DN2,DN(N),TSA,TSB	\$TM1150
900	FORMAT(/7X,22HNO CONVERGENCE DN2 =,F10.5,5X,7HDN(N) =,F10.5,5X	\$TM1160
1	5HTSA =,F10.5,5X,5HTSB =,F10.5/)	\$TM1170
41	CONTINUE	\$TM1180

C		STM1190
C	REGION FROM NOZZLE MIDDLE TO INNER DUCT WALL	STM1200
C		STM1210
	PK=1.	STM1220
	OK=1.	STM1230
	RK=1.	STM1240
	NPRR=NMID-1	STM1250
	DO 46 KK=1,NPRR	STM1260
	I=NMID-KK	STM1270
	IM=I+1	STM1280
C		STM1290
C	JUMP POINT TO TERTIARY STREAM	STM1300
C		STM1310
	IF (I .EQ. LI-1) PK=P03/P01	STM1320
	IF (I .EQ. LI-1) OK = P03/P01*(T01/T03)**G2	STM1330
	IF (I .EQ. LI-1) RK = T03/T01	STM1340
	SX=RK*R(IM,1)*(PSI(I)-PSI(IM))	STM1350
	M(I)= M(IM) - SX/(G**.5 * OK * TS( IM,1)**(G2+0.5) )	STM1360
	B1= ((PK/PS( IM,1)) *(1./G2)-1.)*2.*G1	STM1370
	IF ((I.EQ.LI-1).AND.(B1.LE.0.0)) TMAX=TMAX*0.995	STM1373
	IF ((I.EQ.LI-1).AND.(B1.LE.0.0)) GO TO 595	STM1374
	IF (I .EQ. LI-1) M(I)=SQRT(B1)	STM1380
C		STM1390
C	SET TS, PS, RHO, U, P0, DN, R ACROSS STREAM FROM M(I)	STM1400
C		STM1410
	TS(I,1)=RK/(1.+0.5/G1*(M(I))**2.)	STM1420
	PS(I,1)=OK*(TS(I,1))**G2	STM1430
	RHO(I,1)=PS(I,1)/TS(I,1)	STM1440
	U(I,1)=M(I)*SQRT(G*TS(I,1))	STM1450
	P0(I)=PS(I,1)*(1.+0.5/G1*(M(I))**2.)*G2	STM1460
	Z= 2.0 / (RHO(IM,1)*U(IM,1)+RHO(I,1)*U(I,1))	STM1470
	DN(I)= DN(IM) + Z * (PSI(I) - PSI(IM))	STM1480
	DW1= DN(I)	STM1490
	IF (DW1 .LT. 0.0) DW1= 0.0	STM1500
	DW2= DN2 - DN(I)	STM1510
46	R(I,1)=R1*(1.-DW1/DN2)*EXP(-RDECAY*DW1*ABS(R1))+R2*DW1/DN2*	STM1520
1	EXP(-RDECAY*DW2*ABS(R2))	STM1530
	RETURN	STM1540
599	CONTINUE	STM1550
	MAXERR= 1	STM1560
	RETURN	STM1570
	END	STM1580



	SUBROUTINE EDDY(N,NN,NX,U,PS,CC,DS,E,E1,RHO,VIS,R,Y,	\$ED 0
1	S4, S5, DEL1,DEL2,DEL3,DEL4,CORE1,CORE2,CORE3)	\$ED 10
	*****	\$ED 20
C	SUBROUTINE EDDY	\$ED 30
C		\$ED 40
C	CALCULATES EDDY VISCOSITY VALUES - ARRAY E(N)	\$ED 50
C	REQUIRED PARAMETERS:	\$ED 60
C	ARRAYS OMG, RHO, VIS, Y, PS, DS, S1...S5	\$ED 70
C	VARIABLES N, NN, N1...N6, CORE1, CORE2, CORE3	\$ED 80
C	METHOD	\$ED 90
C	CALCULATES REFERENCE VARIABLES (DP, DELI, DELO,	\$ED 100
C	UREFI, UREFO, CC)	\$ED 110
C	FOR EACH J FROM 1 TO N	\$ED 120
C	DETERMINES THE REGION BY CHECKING N1...N5	\$ED 130
C	CALCULATES EDDY VISCOSITY: E(J)	\$ED 140
C	*****	\$ED 150
C		\$ED 160
	RFAL LZ,LZR,LM	\$ED 170
	DIMENSION S1(190),S2(190),S3(190),S4(190),S5(190),E(190),E1(190),	\$ED 180
1	DS(190),VIS(190),THETA(190,2),R(190,2),U(190,2),	\$ED 190
2	TS(190,2),RHO(190,2),NX(6)	\$ED 200
	DIMENSION OMG(190), Y(190),PS(190,2)	\$ED 210
C	DOUBLE OMG(190), Y(190),PS(190,2)	\$ED 220
C		\$ED 230
C	CALCULATE REFERENCE VARIABLES	\$ED 240
C		\$ED 250
	DP=(PS(1,2)-PS(1,1))/DS(1)	\$ED 260
	DELI=Y(2)-Y(1)	\$ED 270
	DELO=Y(N)-Y(NN)	\$ED 280
	UREFI=SQRT((U(2,2)/DELI-DELI*DP/(2.*VIS(2))) / RHO(2,2) *	\$ED 290
1	VIS(2))	\$ED 300
	DP=(PS(N,2)-PS(N,1))/DS(N)	\$ED 310
	UREFO=SQRT((U(NN,2)/DELO-DELO*DP/(2.*VIS(NN)) )/RHO(NN,2)*	\$ED 320
1	VIS(NN))	\$ED 330
	IF (CORE1.EQ. 0.0) CC= 0.08	\$ED 340
	SCORE = CORE1 + CORE2 + CORE3	\$ED 350
C		\$ED 360
C	DO LOOP WHICH COMPUTES EDDY VALUES E(J)	\$ED 370
C	YS: DIMENSIONLESS DIST TO WALL	\$ED 380
C	LZR: CURVATURE EFFECT	\$ED 390
C	RN: RICHARDSON NUMBER	\$ED 400
C		\$ED 410
C	DIFFERING CALCULATIONS FOR EACH REGION (N1...N6)	\$ED 420
C		\$ED 430
	DO 10 J=2,NN	\$ED 440
	JM=J-1	\$ED 450
	JP=J+1	\$ED 460
	IF (J.LT. NX(1) ) GO TO 20	\$ED 470
	IF (J.GT. NX(6) ) GO TO 21	\$ED 480
	IF (J.LE. NX(2) ) GO TO 30	\$ED 490
	IF (J.GE. NX(5) ) GO TO 30	\$ED 500
	IF (J.LT. NX(3) ) GO TO 29	\$ED 510
	IF (J.GT. NX(4) ) GO TO 40	\$ED 520
	GO TO 30	\$ED 530
C		\$ED 540
C	INNER WALL	\$ED 550
C		\$ED 560
20	DRN=Y(J)-Y(1)	\$ED 570

	YS=DRN*UREFI*RHO(J,1)/VIS(J)	\$ED 580
	LZ=.41*(1.-EXP(-YS/26.))*DRN	\$ED 590
	AZ1=0.09*DEL1	\$ED 600
	DDZ= 1.0 - 2.0 * DRN / Y(N)	\$ED 610
	IF(SCORE.EQ.0.0) AZ1=(.14-.08*DDZ**2-.06*DDZ**4)*Y(N)*.5	\$ED 620
	LZ= LZ * 1.2	\$ED 630
	IF(LZ .GT. AZ1) LZ= AZ1	\$ED 640
	GO TO 25	\$ED 650
C		\$ED 660
C	CORE REGIONS	\$ED 670
C		\$ED 680
30	E(J)=0.	\$ED 690
	GO TO 10	\$ED 700
C		\$ED 710
C	DEL2 REGION	\$ED 720
C		\$ED 730
29	NX3=NX(3)	\$ED 740
	NX2=NX(2)	\$ED 750
	LZ=CC*DEL2*(1+.6*U(NX2,2)/U(NX3,2))	\$ED 760
	GO TO 25	\$ED 770
C		\$ED 780
C	DEL3 REGION	\$ED 790
C		\$ED 800
40	NX5=NX(5)	\$ED 810
	NX4=NX(4)	\$ED 820
	LZ=CC*DEL3*(1+.6*U(NX5,2)/U(NX4,2))	\$ED 830
	GO TO 25	\$ED 840
C		\$ED 850
C	OUTER WALL	\$ED 860
C		\$ED 870
21	CONTINUE	\$ED 880
	DRN= Y(N)-Y(J)	\$ED 890
	YS=DRN*UREFO*RHO(J,1)/VIS(J)	\$ED 900
	LZ=0.41*(1.-EXP(-YS/26.))*DRN	\$ED 910
	AZ4=0.09*DEL4	\$ED 920
	DDZ=1.0-2.0*DRN/Y(N)	\$ED 930
	IF (SCORE.EQ.0.0) AZ4=(0.14-0.08*DDZ**2-0.06*DDZ**4)*Y(N)*0.5	\$ED 940
	IF(LZ .GT. AZ4) LZ=AZ4	\$ED 950
	LZ= LZ * 1.2	\$ED 960
25	CONTINUE	\$ED 970
	RN= 2.0 * (R(J,1) / RHO(J,1))	\$ED 980
	PN=RN/(S5(J)*(U(JP,1)-U(J,1))+S4(J)*(U(J,1)-U(JM,1)))	\$ED 990
	IF (RN .LE. 0.0) LZR= LZ*(2.0 - EXP(3.0*RN))	\$ED1000
	IF (RN .GT. 0.0) LZR= LZ*EXP(-3.0*RN)	\$ED1010
	EI=Y(J)-Y(JM)	\$ED1020
	EJ=Y(JP)-Y(J)	\$ED1030
	EK=Y(JP)-Y(JM)	\$ED1040
	DUIY=ABS(EI*(U(JP,2)-U(J,2))/(EK*EJ)+EJ*(U(J,2)-U(JM,2))/(EK*EI))	\$ED1050
	E(J)=DUIY*LZR*LZR	\$ED1060
10	CONTINUE	\$ED1070
	E(1)= 0.0	\$ED1080
	E(N)= 0.0	\$ED1090
	RETURN	\$ED1100
	END	\$ED1110
	SUBROUTINE WALLS(XW,YW,XX,YY,T,CUR,N,MQ)	\$WA 0
C	*****	\$WA 10
C	SUBROUTINE WALLS SMOOTHS THE BOUNDARY DATA USING A LEAST	\$WA 20
C	SQUARES PROCEDURE. IT ALSO INTERPOLATES THE SMOOTHED	\$WA 30
C	DATA TO GET THE CURVATURE AND SLOPE AT ANY POINT	\$WA 40
C	*****	\$WA 50
	DIMENSION XW(99,2),YW(99,2),YP(99,2),CURV(99,2),NDI(10,2),	\$WA 60
1	YB(99,2)	\$WA 70

	COMMON/KURV/MST,MSS,NDI,CURV	SWA 80
	IF(MQ.EQ.1) M=MST	SWA 90
	IF(MQ.EQ.2) M=MSS	SWA 100
	IF(XX.LT.999.) GO TO 60	SWA 110
	N=N+M-1	SWA 120
	NX=N	SWA 130
	MP=M-1	SWA 140
	DO 4 J=1,MP	SWA 150
	JK=M-J	SWA 160
	JK1=JK+1	SWA 170
	NP=NDI(JK1,MQ)-NDI(JK,MQ)+1	SWA 180
	DO 3 I=1,NP	SWA 190
	K=NX+1-I	SWA 200
	L=K-JK	SWA 210
	XW(K,MQ)=XW(L,MQ)	SWA 220
3	YW(K,MQ)=YW(L,MQ)	SWA 230
	NDI(JK1,MQ)=NX	SWA 240
4	NX=NX-NP	SWA 250
	DO 5 I=1,N	SWA 260
5	YB(I,MQ)=YW(I,MQ)	SWA 270
	KM=1	SWA 280
	DO 198 MP=1,M	SWA 290
	NDA=NDI(MP,MQ)	SWA 300
	KL=NDA-KM+1	SWA 310
	IF(KL.LT.3) GO TO 700	SWA 320
	IF(KL.EQ.3) GO TO 30	SWA 330
	IF(KL.EQ.4) GO TO 40	SWA 340
	L=5	SWA 350
16	DO 20 KS=KM,NDA	SWA 360
	IF(L.EQ.4) GO TO 17	SWA 370
	CALL LSQ(KS,KM,NDA,L,A1,A2,A3,XW(1,MQ),YB(1,MQ))	SWA 380
17	YW(KS,MQ)=A1+A2*XW(KS,MQ)+A3*XW(KS,MQ)**2	SWA 390
	YP(KS,MQ)=A2+2.*A3*XW(KS,MQ)	SWA 400
	CURV(KS,MQ)=-2.*A3/(1.+YP(KS,MQ)**2)**1.5	SWA 410
	IF(ABS(CURV(KS,MQ)).LT..001) CURV(KS,MQ)=0.	SWA 420
20	CONTINUE	SWA 430
	GO TO 198	SWA 440
30	C1=XW(KM,MQ)-XW(KM+1,MQ)	SWA 450
	C2=XW(KM,MQ)-XW(NDA,MQ)	SWA 460
	C3=XW(KM+1,MQ)-XW(NDA,MQ)	SWA 470
	YPP=2.*(YW(KM,MQ)/C1/C2-YW(KM+1,MQ)/C1/C3+YW(NDA,MQ)/C2/C3)	SWA 480
	DO 31 K=KM,NDA	SWA 490
	YP(K,MQ)=(2.*XW(K,MQ)-XW(NDA,MQ)-XW(KM+1,MQ))*YW(KM,MQ)/C1/C2	SWA 500
1	- (2.*XW(K,MQ)-XW(KM,MQ)-XW(NDA,MQ))*YW(KM+1,MQ)/C1/C3	SWA 510
2	+ (2.*XW(K,MQ)-XW(KM,MQ)-XW(KM+1,MQ))*YW(NDA,MQ)/C2/C3	SWA 520
	CURV(K,MQ)=-2.*YPP/(1.+YP(K,MQ)**2)**1.5	SWA 530
	IF(ABS(CURV(K,MQ)).LT.0.001) CURV(K,MQ)=0.	SWA 540
31	CONTINUE	SWA 550
	GO TO 198	SWA 560
40	L=4	SWA 570
	CALL LSQ(KS,KM,NDA,L,A1,A2,A3,XW(1,MQ),YB(1,MQ))	SWA 580
	GO TO 16	SWA 590
198	KM=NDA+1	SWA 600
	RETURN	SWA 610
60	NDA=1	SWA 620
	DO 70 I=1,M	SWA 630
	J=NDI(I,MQ)	SWA 640
	IF((XX.GE.XW(NDA,MQ)).AND.(XX.LE.XW(J,MQ))) GO TO 71	SWA 650
	IF(I.EQ.M) GO TO 71	SWA 660
70	NDA=J+1	SWA 670
71	JJ=J-1	SWA 680
	DO 72 K=NDA,JJ	

	IF((XX.GE.XW(K,MQ)).AND.(XX.LE.XW(K+1,MQ))) GO TO 73	\$WA 690
72	CONTINUE	\$WA 700
73	IF(K.EQ.NDA) K=NDA+1	\$WA 710
	C1=XW(K+1,MQ)-XW(K,MQ)	\$WA 720
	C2=XW(K+1,MQ)-XW(K-1,MQ)	\$WA 730
	C3=XW(K,MQ)-XW(K-1,MQ)	\$WA 740
	S1=XX-XW(K+1,MQ)	\$WA 750
	S2=XX-XW(K,MQ)	\$WA 760
	S3=XX-XW(K-1,MQ)	\$WA 770
	YY=S3*S2*YW(K+1,MQ)/C1/C2-S1*S3*YW(K,MQ)/C1/C3+S1*S2*YW(K-1,MQ)/	\$WA 780
	1 C2/C3	\$WA 790
	CUR=S2*S3*CURV(K+1,MQ)/C1/C2-S1*S3*CURV(K,MQ)/C1/C3	\$WA 800
	1 +S1*S2*CURV(K-1,MQ)/C2/C3	\$WA 810
	T=S3*S2*YP(K+1,MQ)/C1/C2-S1*S3*YP(K,MQ)/C1/C3+S1*S2*YP(K-1,MQ)	\$WA 820
	1 /C2/C3	\$WA 830
	T=ATAN(T)	\$WA 840
	RETURN	\$WA 850
700	WRITE(6,900)	\$WA 860
900	FORMAT(/4X,34H LESS THAN THREE POINTS IN SEGMENT)	\$WA 870
	STOP	\$WA 880
	END	\$WA 890
	SUBROUTINE LSQ(KS,KM,NDA,L,A1,A2,A3,C,B)	\$WA 900
	DIMENSION B(99),C(99),S(2,4)	\$LS 0
	K=KS	\$LS 10
	IF((KS.EQ.NDA).OR.(KS.EQ.NDA-1)) K=NDA-2	\$LS 20
	IF((KS.EQ.KM).OR.(KS.EQ.KM+1)) K=KM+2	\$LS 30
	IF(L.EQ.4) K=KM+2	\$LS 40
	DO 5 I=1,2	\$LS 50
	DO 5 J=1,4	\$LS 60
	N=I-1	\$LS 70
	S(I,J)=0.	\$LS 80
	DO 4 M=1,L	\$LS 90
	KK=K+M-3	\$LS 100
	IF(C(KK).EQ.0.) C(KK)=1.E-06	\$LS 110
	IF(B(KK).EQ.0.) B(KK)=1.E-06	\$LS 120
4	S(I,J)= S(I,J)+B(KK)**N*C(KK)**(J-N)	\$LS 130
5	CONTINUE	\$LS 140
	A=L	\$LS 150
	D=A*(S(1,2)*S(1,4)-S(1,3)**2)-S(1,1)*(S(1,1)*S(1,4)-S(1,2)*S(1,3))	\$LS 160
	1 +S(1,2)*(S(1,1)*S(1,3)-S(1,2)**2)	\$LS 170
	A1=(S(2,1)*(S(1,2)*S(1,4)-S(1,3)**2)-S(1,3)*(S(1,1)*S(1,4)-S(1,2)*S(1,3))	\$LS 180
	1 S(1,2)*S(1,3))+S(2,3)*(S(1,1)*S(1,3)-S(1,2)**2))/D	\$LS 190
	A2=(S(2,1)*(S(1,2)*S(1,3)-S(1,1)*S(1,4))+S(2,2)*(A*S(1,4)-	\$LS 200
	1 S(1,2)**2)+S(2,3)*(S(1,1)*S(1,2)-A*S(1,3)))/D	\$LS 210
	A3=(S(2,1)*(S(1,1)*S(1,3)-S(1,2)**2)-S(2,2)*(A*S(1,3)-	\$LS 220
	1 S(1,1)*S(1,2))+S(2,3)*(A*S(1,2)-S(1,1)**2))/D	\$LS 230
	RETURN	\$LS 240
	END	\$LS 250
	SUBROUTINE SOLV(DS,N,NN,PR,PRT,G)	\$LS 260
	REAL M(190)	\$SO 0
	DIMENSION S1(190),S2(190),S3(190),S4(190),S5(190),E(190),E1(190),	\$SO 10
1	DS(190),VIS(190),X(190),Y(190),P0(190),H(190),	\$SO 20
2	THETA(190,2),R(190,2),U(190,2),TS(190,2),RHO(190,2),	\$SO 30
3	A(190),R(190),C(190),D(190)	\$SO 40
	DIMENSION OMG(190),DN(190),PS(190,2)	\$SO 50
C	DOUBLE, OMG(190),DN(190),PS(190,2)	\$SO 60
	DIMENSION PSI(180)	\$SO 70
C	DOUBLE PSI(190)	\$SO 80
	COMMON/INCD/ U,PS,TS,M,P0,THETA,R,X,Y,DN	\$SO 90
	COMMON/SOLV/ S1,S2,S3,E,E1,H,S4,S5	\$SO 100
	COMMON/INEDY/PSI,RHO,VIS	\$SO 110
	CL=(G-1.)/G	\$SO 120
	C0 = 0.	\$SO 130

	IF (PR.EQ..7) C0 = 1.	\$S0 140
	DO 10 J=2,NN	\$S0 150
	JP=J+1	\$S0 160
	JM= J-1	\$S0 170
	Y1 =E1(JP)*(VIS(JP)/PR + RH0(JP,1)*E(JP)/PRT)	\$S0 180
	Y2=E1(JM)*( VIS(JM)/PR + RH0(JM,1)*E(JM)/PRT)	\$S0 190
	Y3 =E1(J) *(VIS(J )/PR + RH0(J ,1)*E(J )/PRT)	\$S0 200
	Y4 = (Y1+Y3)/S2(J)	\$S0 210
	Y5 =(Y3+Y2)/S3(J)	\$S0 220
	Y6= Y4 / S1(J)	\$S0 230
	Y7= Y5 / S1(J)	\$S0 240
	Y8= Y6 * U(J,1)	\$S0 250
	Y9= Y7 * U(J,1)	\$S0 260
	A(JM) = U(J,1)/DS(J) + Y8 + Y9	\$S0 270
	B(JM) = -Y8	\$S0 280
	C(JM) = -Y9	\$S0 290
	IF (PR.NE.0.7) D(JM)=U(J,1)**2/DS(J)-(PS(J,2)-PS(J,1))/DS(J)*(1./	\$S0 300
	1 (RH0(J,1)))	\$S0 310
	IF (PR.EQ.0.7) D(JM)=TS(J,1)*U(J,1)/DS(J)+(CL*U(J,1)/(RH0(J,1)))	\$S0 320
	1 *(PS(J,2)-PS(J,1))/DS(J)+RH0(J,1)*CL*(VIS(J)+RH0(J,1)*E(J))*(U	\$S0 330
	2 (J,1)**2*(S5(J)*(U(JP,2)-U(J,2))+S4(J)*(U(J,2)-U(JM,2))))**2	\$S0 340
10	CONTINUE	\$S0 350
	A(1) =A(1) + C0*C(1)	\$S0 360
	A(N-2) = A(N-2) + C0*B(N-2)	\$S0 370
	CALL CALC(A,B,C,D,H,N)	\$S0 380
413	RETURN	\$S0 390
	END	\$S0 400
	SUBROUTINE CALC(A,B,C,D,H,J)	\$S0 410
	1 DIMENSION A(190),B(190),C(190),D(190),W(190),G(190),H(190),Q(190)	\$CA 0
	N2 = J-2	\$CA 10
	N1= J-2	\$CA 20
	W(1) = A(1)	\$CA 30
	G(1) = D(1)/W(1)	\$CA 40
	DO 1 K = 2,N2	\$CA 50
	K1 = K-1	\$CA 60
	Q(K1) = R(K1)/W(K1)	\$CA 70
	W(K) = A(K)-C(K)*Q(K1)	\$CA 80
1	G(K) = ( D(K)-C(K)*G(K1))/W(K)	\$CA 90
	H(N2) = G(N2)	\$CA 100
	N3 =J-3	\$CA 110
	DO 2 K = 1,N3	\$CA 120
	KK = N2-K	\$CA 130
2	H(KK) = G(KK) -Q(KK)*H(KK+1)	\$CA 140
	RETURN	\$CA 150
	END	\$CA 160
	SUBROUTINE DSMOVE(XW,YW,XX,YY,DSX,THETA,RR,NPAIR,J)	\$CA 170
C		\$DS 0
C	*****	\$DS 10
C		\$DS 20
C	SUBROUTINE DSMOVE	\$DS 30
C	FINDS NEW X VALUE AT A GIVEN DISTANCE DSX ALONG WALL SURFACE	\$DS 40
C	METHOD:	\$DS 50
C	GUESSES VALUES OF DX UNTIL (DSX - DS COMPUTED) = 0.0	\$DS 60
C	*****	\$DS 70
C		\$DS 80
C		\$DS 90
C	DIMENSION XW(99,2),YW(99,2)	\$DS 100
C		\$DS 110
C	APPROXIMATE POINT X(N), Y(N)	\$DS 120
C	AT DISTANCE DS(N) ALONG UPPER WALL	\$DS 130
C	(TWO METHODS OF APPROXIMATING, DEPENDING ON R1)	\$DS 140

		\$DS 150
	DXSTAR= DSX * ABS(COS(THETA))	\$DS 160
	XSTAR= DXSTAR + XX	\$DS 170
	CALL WALLS(XW,YW, XSTAR, YSTAR, TSTAR, RSTAR,NPAIR, J)	\$DS 180
	DDSTAR= DSX - CURVDS(DXSTAR, RR,RSTAR, THETA,TSTAR,YY,YSTAR)	\$DS 190
	DXB= DSX * ABS(COS((THETA + TSTAR)*0.5))	\$DS 200
	XB= DXB + XX	\$DS 210
	CALL WALLS(XW,YW,XB, YB, TB, RB, NPAIR, J)	\$DS 220
	DDIFFB= DSX - CURVDS(DXB, RR, RB, THETA, TB,YY, YB)	\$DS 230
	DO 510 I=1,50	\$DS 240
	IF (ABS(DDIFFB-DDSTAR) .LE. 0.1 **8) GO TO 512	\$DS 250
	DXB=(DDIFFB*DXSTAR - DDSTAR*DXB) / (DDIFFB-DDSTAR)	\$DS 260
	IF (DXB .LT. 0.0) DXB= DXSTAR * 0.1	\$DS 270
	XB= DXB + XX	\$DS 280
	CALL WALLS(XW, YW, XB, YB, TB, RB, NPAIR, J)	\$DS 290
	DDIFFB= DSX - CURVDS(DXB, RR, RB, THETA, TB,YY, YB)	\$DS 300
C		\$DS 310
C	CHECK FOR TERMINATION	\$DS 320
C		\$DS 330
	IF (ABS(XB-XSTAR).LE. 0.1 **6) GO TO 512	\$DS 340
	IF (ABS(DDIFFB) .LE. 0.1 ** 7) GO TO 512	\$DS 350
510	CONTINUE	\$DS 360
	WRITE(6,900) XR, DDIFFB	\$DS 370
900	FORMAT(/7X,21HNO CONVERGENCE DSMOVE,5X,4HXB =F10.4,5X,8HDDIFFB =,	\$DS 380
1	F10.5/)	\$DS 390
512	CONTINUE	\$DS 400
	XX= XB	\$DS 410
	RETURN	\$DS 420
	END	\$DS 430
	REAL FUNCTION CURVDS(DX,RA,RB,TA,TR,YA,YB)	\$CU 0
	IF (ABS(TA-TB) .LE. 0.0001) GO TO 500	\$CU 10
	DR= ABS(RA+RB)	\$CU 20
	IF (DR .LE. 0.02 ) GO TO 500	\$CU 30
	CURVDS= (2.0/DR) * ABS(TB-TA)	\$CU 40
	RETURN	\$CU 50
500	CURVDS= SQRT(DX ** 2 + (ABS(YB-YA)**2) )	\$CU 60
	RETURN	\$CU 70
	END	\$CU 80
	SUBROUTINE LOOK(N,NN,K3,NCORE,DEL1,DEL2,DEL3,DEL4,	\$LO 0
1	CORE1,CORE2,CORE3, NX, P0, DN, U, DIP, NMAX)	\$LO 10
	*****	\$LO 20
C		\$LO 30
C	SUBROUTINE LOOK	\$LO 40
C		\$LO 50
C	CALCULATES	\$LO 60
C	EDGES OF POTENTIAL FLOWS - N1... N6	\$LO 70
C	WIDTH OF FLOWS - DEL1, DEL2, DEL3, DEL4	\$LO 80
C	WIDTH OF CORES - CORE1, CORE2, CORE3	\$LO 90
C	REQUIRED PARAMETERS:	\$LO 100
C	ARRAYS DN, P0	\$LO 110
C	VARIABLES N, NN, NSTEP, K3, P01, P02, P03	\$LO 120
C	METHOD	\$LO 130
C	CALCULATES TOLERANCES TOL1, TOL2, TOL3, TOL4	\$LO 140
C	CALCULATES ARRAY DIP ~ RATE OF CHANGE OF P0	\$LO 150
C		\$LO 160
C	FOR EACH J FROM 2 TO NN	\$LO 170
C	SCAN ARRAY DIP(J)	\$LO 180
C	FOR REGION WHERE DIP GT TOLERANCE	\$LO 190
C	SET FLOW REGION EDGE NX	\$LO 200
C	SET FLOW REGION WIDTH DELX	\$LO 210
C	FOR REGION WHERE DIP(J) LT TOLERANCE	\$LO 220
C	SET FLOW REGION EDGE NX	\$LO 230
C	COMPUTE CORE REGION AREA COREX	\$LO 240
C	*****	

C	DIMENSION X(190),Y(190),U(190,2),DIP(190),NX(6),P0(190),H(190)	\$LO 250
	DIMENSION DN(190)	\$LO 260
C	DOUBLE DN(190)	\$LO 270
	COMMON/CONST/ P01,P02,P03, T01,T02,T03, X1, XC, YC	\$LO 280
C		\$LO 290
C	SET TOLERANCES	\$LO 300
C		\$LO 310
	TOL1=0.01*P03/DN(N) / P01	\$LO 320
	TOL2=0.01*(P01-P03)/DN(N) / P01	\$LO 330
	TOL3= 0.01*(P01-P02)/DN(N) / P01	\$LO 340
	TOL4= 0.01*P02/DN(N) / P01	\$LO 350
	IF (K3.EQ.0) CORE3=0.	\$LO 360
		\$LO 370
C		\$LO 380
C	THIS LOOP FINDS A VALUE FOR NMAX	\$LO 390
C		\$LO 400
	DO 11 J=2,N	\$LO 410
	DIP(J)=(P0(J)-P0(J-1))/(DN(J)-DN(J-1))	\$LO 420
	DIP(J)= ABS(DIP(J))	\$LO 430
11	CONTINUE	\$LO 440
	NMAX=1	\$LO 450
	DO 12 J=2,N	\$LO 460
	IF (U(J,2) .GT. U(NMAX,2) ) NMAX= J	\$LO 470
12	CONTINUE	\$LO 480
C		\$LO 490
C	THIS LOOP COMPUTES EDGE VALUES FOR THE PRIMARY,,SECONDARY	\$LO 500
C	AND TERTIARY POTENTIAL FLOWS	\$LO 510
C	TERTIARY FLOW REGION	\$LO 520
C		\$LO 530
	DO 50 I=2,N	\$LO 540
	IF (DIP(I) .LT. TOL1) GO TO 52	\$LO 550
50	CONTINUE	\$LO 560
52	CONTINUE	\$LO 570
	NX(1)= I-1	\$LO 580
C		\$LO 590
C	SECONDARY FLOW REGION	\$LO 600
C		\$LO 610
	DO 60 M=2,N	\$LO 620
	I= N + (2-M)	\$LO 630
	IF (DIP(I) .LT. TOL4) GO TO 62	\$LO 640
60	CONTINUE	\$LO 650
62	CONTINUE	\$LO 660
	NX(6)= I	\$LO 670
	IF (NX(1) .GT. NX(2) ) CORE3= 0.0	\$LO 680
	IF (NX(6) .LT. NX(5) ) CORE2= 0.0	\$LO 690
	IF (CORE3 .EQ. 0.0) GO TO 102	\$LO 700
C		\$LO 710
C	TERTIARY CORE REGION	\$LO 720
C		\$LO 730
	NPRR=NX(1)+1	\$LO 740
	DO 70 I=NPRR,N	\$LO 750
	IF (DIP(I) .GT. TOL2) GO TO 72	\$LO 760
70	CONTINUE	\$LO 770
72	CONTINUE	\$LO 780
	NX(2)= I-1	\$LO 790
	NPRR=NX(2)+1	\$LO 800
	DO 75 I=NPRR,N	\$LO 810
	IF (DIP(I) .LT. TOL2) GO TO 77	\$LO 820
75	CONTINUE	\$LO 830
77	CONTINUE	\$LO 840
	NX(3)= I-1	\$LO 850
102	CONTINUE	\$LO 860
	IF (CORE2 .EQ. 0.0) GO TO 101	\$LO 870

C		\$L0 880
C	SECONDARY CORE REGION	\$L0 890
C		\$L0 900
	NPRR=NX(6)	\$L0 910
	DO 80 M=2,NPRR	\$L0 920
	I= NX(6) + (2-M)	\$L0 930
	IF (DIP(I) .GT. TOL3) GO TO 82	\$L0 940
80	CONTINUE	\$L0 950
82	CONTINUE	\$L0 960
	NX(5)= I	\$L0 970
	NPRR=NX(5)	\$L0 980
	DO 85 M=2,NPRR	\$L0 990
	I= NX(5) + (2-M)	\$L01000
	IF (DIP(I) .LT. TOL3) GO TO 88	\$L01010
85	CONTINUE	\$L01020
88	CONTINUE	\$L01030
	NX(4)= I	\$L01040
101	CONTINUE	\$L01050
C		\$L01060
C	CHECK IF PRIMARY CORE EXISTS	\$L01070
C		\$L01080
	IF (CORE1 .EQ. 0.0) GO TO 110	\$L01090
	IF (CORE3 .NE. 0.0) GO TO 120	\$L01100
	DO 130 I=2,N	\$L01110
	IF (DIP(I) .LT. TOL2) GO TO 132	\$L01120
130	CONTINUE	\$L01130
132	CONTINUE	\$L01140
	NX(3)= I-1	\$L01150
120	CONTINUE	\$L01160
	IF (CORE2 .NE. 0.0) GO TO 140	\$L01170
	DO 135 M=2,N	\$L01180
	I= N + (2-M)	\$L01190
	IF (DIP(I) .LT. TOL3) GO TO 137	\$L01200
135	CONTINUE	\$L01210
137	CONTINUE	\$L01220
	NX(4)= I	\$L01230
140	CONTINUE	\$L01240
	IF (NX(3) .GE. NX(4) ) CORE1=0.0	\$L01250
110	CONTINUE	\$L01260
C		\$L01270
C	CHECK VALUES OF , CORE1, CORE2, CORE3	\$L01280
C	NMAX IS MAXIMUM U( ) VALUE	\$L01290
C		\$L01300
	IF (CORE1 .NE. 0.0 ) GO TO 670	\$L01310
C		\$L01320
C	CORE1 = 0.0	\$L01330
C		\$L01340
	NX(3)=NMAX	\$L01350
	NX(4)=NMAX	\$L01360
670	IF (CORE3 .NE. 0.0) GO TO 680	\$L01370
C		\$L01380
C	CORE3 = 0.0	\$L01390
C		\$L01400
	NX(1)=NX(3)	\$L01410
	NX(2)=NX(3)	\$L01420
680	IF (CORE2 .NE. 0.0) GO TO 90	\$L01430
C		\$L01440
C	CORE2 = 0.0	\$L01450
C		\$L01460
	NX(5)=NX(4)	\$L01470
	NX(6)=NX(4)	\$L01480
90	CONTINUE	\$L01490
	NX1=NX(1)	



		\$L01500
	NX2=NX(2)	\$L01510
	NX3=NX(3)	\$L01520
	NX4=NX(4)	\$L01530
	NX5=NX(5)	\$L01540
	NX6=NX(6)	\$L01550
	CORE1=DN(NX4)-DN(NX3)	\$L01560
	CORE2=DN(NX6)-DN(NX5)	\$L01570
	CORE3=DN(NX2)-DN(NX1)	\$L01580
	DEL1=DN(NX1)	\$L01590
	DEL3=DN(NX5)-DN(NX4)	\$L01600
	DEL2=DN(NX3)-DN(NX2)	\$L01610
	DEL4=DN(N)-DN(NX6)	\$L01620
	IF ((CORE1.EQ.0.).AND.(CORE2.EQ.0.).AND.(CORE3.EQ.0.)) NCORE= 0	\$L01630
	RETURN	\$L01640
	END	\$L01650
	SUBROUTINE CORDS(N,X,Y,THETA,U,RHO,PSI,DN,NDUCT,DDIST)	\$CO 0
C		\$CO 10
C	*****	\$CO 20
C	SUBROUTINE CORDS	\$CO 30
C	CALCULATES	\$CO 40
C	DISTANCE BETWEEN STREAM LINES (ARRAY DN)	\$CO 50
C	DRAWS ARCS FROM TOP AND BOTTOM WALL TO MIDDLE OF DUCT	\$CO 60
C	COMPARES Y COORD OF BOTH ARCS AT MIDDLE POINT	\$CO 70
C	*****	\$CO 80
C		\$CO 90
	REAL M(190)	\$CO 100
	DIMENSION X(190),Y(190),P0(190),H(190),THETA(190,2),R(190,2),	\$CO 110
1	U(190,2),TS(190,2),RHO(190,2)	\$CO 120
	DIMENSION OMG(190),DN(190)	\$CO 130
C	DOUBLE OMG(190),DN(190)	\$CO 140
	DIMENSION PSI(180)	\$CO 150
C	DOUBLE PSI(190)	\$CO 160
	DN(1) = 0.	\$CO 170
	DO 2 J=2,N	\$CO 180
	JM=J-1	\$CO 190
	7=2./(RHO(JM,2)*U(JM,2)+RHO(J,2)*U(J,2))	\$CO 200
	DN(J)= DN(JM) + 7 * (PSI(J) - PSI(JM))	\$CO 210
2	CONTINUE	\$CO 220
	DO 50 J=2,NDUCT	\$CO 230
	JM=J-1	\$CO 240
	X(J)= X(JM)-(DN(J)-DN(JM))*SIN(THETA(J,2))	\$CO 250
	Y(J)= Y(JM)+(DN(J)-DN(JM))*COS(THETA(J,2))	\$CO 260
50	CONTINUE	\$CO 270
	NPAR=NDUCT+1	\$CO 280
	NPARR=N-1	\$CO 290
	DO 52 MM=NPARR,NPARR	\$CO 300
	J=N-MM+NDUCT	\$CO 310
	JM=J+1	\$CO 320
	X(J)= X(JM)-(DN(J)-DN(JM))*SIN(THETA(J,2))	\$CO 330
	Y(J)= Y(JM)+(DN(J)-DN(JM))*COS(THETA(J,2))	\$CO 340
52	CONTINUE	\$CO 350
	YCOMP= Y(NDUCT+1) + (DN(NDUCT) - DN(NDUCT+1)) * COS(THETA(J,2))	\$CO 360
	DDIST= (Y(NDUCT) - YCOMP) / DN(N)	\$CO 370
	RETURN	\$CO 380
	END	\$CO 390
	SUBROUTINE REMOVE(N,NN,NGRID,NX,NMAX,OMG,S1,S2,S3,S4,S5,DS,VIS,	\$RE 0
1	DN,X,Y,M,P0,H,THETA,R,PS,U,TS,RHO)	\$RE 10
C		\$RE 20
C	*****	\$RE 30
C	SUBROUTINE REMOVEPTS	\$RE 40

C	REMOVES GRID POINTS AS PROGRAM ITERATES	\$RE 50
C	(TO REDUCE COMPUTING TIME)	\$RE 60
C	REQUIRED PARAMETERS:	\$RE 70
C	N, NGRID, ARRAY U(90)	\$RE 80
C	(ALL OTHER ARRAYS)	\$RE 90
C	METHOD:	\$RE 100
C	FOR EVERY TENTH NSTEP	\$RE 110
C	SCANS ARRAY U	\$RE 120
C	IF (U(I-1) - U(I) / U(I) ) LESS THAN TOLERANCE,	\$RE 130
C	THEN GRID POINT IS REMOVED	\$RE 140
C	CORE REGIONS ARE NOT SCANNED	\$RE 150
C	ARRAYS COPIED ARE (OMG,S1,S2,S3,S4,S5,DS,VIS,DN,X,Y,M,P0,H,	\$RE 160
C	THETA,R,PS,U,TS,RHO)	\$RE 170
C	*****	\$RE 180
C		\$RE 190
C		\$RE 200
C	REAL M(190)	\$RE 210
C	DIMENSION S1(190),S2(190),S3(190),S4(190),S5(190),E(190),E1(190),	\$RE 220
1	DS(190),VIS(190),X(190),Y(190),P0(190),H(190),	\$RE 230
2	THETA(190,2),R(190,2),U(190,2),TS(190,2),RHO(190,2),	\$RE 240
3	NXSAVE(6),NX(6)	\$RE 250
C	DIMENSION OMG(190),DN(190),PS(190,2)	\$RE 260
C	DOUBLE OMG(190),DN(190),PS(190,2)	\$RE 270
C	NGRID= NGRID+1	\$RE 280
C	IF (NGRID .LT. 10) RETURN	\$RE 290
C	NGRID= 0	\$RE 300
C	DO 300 I=1,6	\$RE 310
C	NXSAVE(I)= NX(I)	\$RE 320
300	CONTINUE	\$RE 330
C	J= 1	\$RE 340
C	DO 24 I=2,N	\$RE 350
C		\$RE 360
C	DO NOT REMOVE GRID POINTS FROM CORE REGIONS	\$RE 370
C		\$RE 380
C	IF ((I.GT. (NXSAVE(1)-3)) .AND. (I.LE.(NXSAVE(2)+3))) GO TO 51	\$RE 390
C	IF ((I.GT. (NXSAVE(3)-3)) .AND. (I.LE.(NXSAVE(4)+3))) GO TO 51	\$RE 400
C	IF ((I.GT. (NXSAVE(5)-3)) .AND. (I.LE.(NXSAVE(6)+3))) GO TO 51	\$RE 410
C	IF (I .EQ. N) GO TO 51	\$RE 420
C	UR= ABS((U(I,2) - U(J,2)) / U(I,2))	\$RE 430
C	DNN= ABS( (DN(I) - DN(J) ) / DN(N) )	\$RE 440
C	IF ( ( UR .LT. 0.02) .AND. (DNN .LT. 0.015) ) GO TO 50	\$RE 450
C		\$RE 460
C	GO TO STATEMENT 50 WILL REMOVE GRID POINT	\$RE 470
C		\$RE 480
51	J= J + 1	\$RE 490
50	OMG(J) = OMG(I)	\$RE 500
	DS(J) = DS(I)	\$RE 510
	VIS(J) = VIS(I)	\$RE 520
	DN(J) = DN(I)	\$RE 530
	X(J) = X(I)	\$RE 540
	Y(J) = Y(I)	\$RE 550
	M(J) = M(I)	\$RE 560
	P0(J) = P0(I)	\$RE 570
	H(J) = H(I)	\$RE 580
	THETA(J,2) = THETA(I,2)	\$RE 590
	R(J,2) = R(I,2)	\$RE 600
	PS(J,2) = PS(I,2)	\$RE 610
	U(J,2) = U(I,2)	\$RE 620
	TS(J,2) = TS(I,2)	\$RE 630
	RHO(J,2) = RHO(I,2)	\$RE 640
	IF (NX(1) .EQ. I) NX(1) = J	\$RE 650
	IF (NX(2) .EQ. I) NX(2) = J	\$RE 660
	IF (NX(3) .EQ. I) NX(3) = J	\$RE 670

350	GO TO 320	\$AR 450
	CONTINUE	\$AR 460
	XXA=XX	\$AR 470
	DRC= DRC	\$AR 480
	IF(FW*DRC .GT.0.0) DRCB= DRCB*0.5	\$AR 490
320	CONTINUE	\$AR 500
	FW= DRC	\$AR 510
C		\$AR 520
C	CHECK FOR TERMINATION	\$AR 530
C		\$AR 540
	IF(ABS(XXB-XXA) .LE. 1.E-6) GO TO 310	\$AR 550
	IF(ABS(DRC) .LE. 1.E-9 ) GO TO 310	\$AR 560
300	CONTINUE	\$AR 570
	WRITE(6,900)	\$AR 580
900	FORMAT(/7X,35HNO CONVERGENCE IN FINDING XC AND YC/)	\$AR 590
310	CONTINUE	\$AR 600
	RETURN	\$AR 610
	END	\$AR 620
	PEAL FUNCTION CENTRE(XNOZ,YNOZ,MNOZ,BNOZ,XX,RC,XW,YW,NPAIR,J)	\$CE 0
C		\$CE 10
C	*****	\$CE 20
C	FUNCTION CENTRE	\$CE 30
C	AIDS IN FINDING A SUITABLE TANGENT TO UPPER WALL	\$CE 40
C	GIVEN XNOZ, XX - PROJECTS TWO TANGENTS	\$CE 50
C	FINDS INTERSECTION XC, YC	\$CE 60
C	RETURNS DIFFERENCE BETWEEN LENGTH OF RADIUS	\$CE 70
C	*****	\$CE 80
C		\$CE 90
	REAL MNOZ, MM	\$CE 100
	DIMENSION XW(99,2),YW(99,2)	\$CE 110
	DIST(XX,YY,XXC,YYC)=SQRT((ABS(XX-XXC)**2)+(ABS(YY-YYC)**2))	\$CE 120
	CALL WALLS(XW,YW,XX,YY,TT,RR,NPAIR,J)	\$CE 130
	MM= TAN(TT)	\$CE 140
	RR= YY - (MM* XX)	\$CE 150
	IF (ABS(MNOZ-MM) .GT. 0.00001) XC= (BB-BNOZ) / (MNOZ-MM)	\$CE 160
	IF (ABS(MNOZ-MM) .LE. 0.00001) XC= (XNOZ+XX) * 0.5	\$CE 170
	YC= MNOZ*XC + BNOZ	\$CE 180
	RC= DIST(XNOZ,YNOZ,XC,YC)	\$CE 190
	CENTRE= DIST(XNOZ,YNOZ,XC,YC) - DIST(XX,YY,XC,YC)	\$CE 200
	RETURN	\$CE 210
	END	\$CE 220
	SUBROUTINE OMGSET(OMG,JSTART,JFIN,JTSER,JFNSER,DMST,DMFN,DMMID,	\$OM 0
1	RST,RFIN)	\$OM 10
C		\$OM 20
C	*****	\$OM 30
C	SUBROUTINE OMGSET	\$OM 40
C	SETS OMG DISTRIBUTION ARRAY OMG	\$OM 50
C	FOR EACH CORE REGION, OMGs ARE CONSTANT IN MIDDLE	\$OM 60
C	CLOSE TOGETHER AT EDGES	\$OM 70
C	OMG(JSTART) AND OMG(JFIN) MUST ALREADY BE SPECIFIED	\$OM 80
C	*****	\$OM 90
C		\$OM 100
	DIMENSION OMG(190)	\$OM 110
C	DOUBLE OMG(190)	\$OM 120
	JCONST= JSTART + JTSER + 1	\$OM 130
	OMG(JCONST)= OMG(JSTART) + DMMID	\$OM 140
	NPRR=JCONST+1	\$OM 150
	NPRRR=JFIN-JFNSER-1	\$OM 160
	DO 500 J=NPRR,NPRRR	\$OM 170
500	OMG(J)=OMG(J-1)+DMMID	\$OM 180
C	SOLVE FOR RST IN SERIES A+AR+AR**2+...	\$OM 190
	XA= 1.5	\$OM 200
	FXA= SERIES(DMMID,DMST,JTSER,XA)	\$OM 210

	IF (NX(4) .EQ. I) NX(4) = J	\$RE 680
	IF (NX(5) .EQ. I) NX(5) = J	\$RE 690
	IF (NX(6) .EQ. I) NX(6) = J	\$RE 700
	IF (NMAX .EQ. I) NMAX= J	\$RE 710
24	CONTINUE	\$RE 720
	N= J	\$RE 730
	NN = N - 1	\$RE 740
	DO 20 J=2,NN	\$RE 750
	JP= J+1	\$RE 760
	JM= J-1	\$RE 770
	S1(J) = OMG(JP) - OMG(JM)	\$RE 780
	S2(J)= OMG(JP) - OMG(J)	\$RE 790
	S3(J)= OMG(J) - OMG(JM)	\$RE 800
	S4(J)= S2(J) / S3(J) / S1(J)	\$RE 810
	S5(J)= S3(J) / S2(J) / S1(J)	\$RE 820
20	CONTINUE	\$RE 830
	RETURN	\$RE 840
	END	\$RE 850
	SUBROUTINE ARCDIS(XNOZ,YNOZ,TNOZ,XX,RC,XW,YW,NPAIR,J)	\$AR 0
C		\$AR 10
C	*****	\$AR 20
C	SUBROUTINE ARCDIS	\$AR 30
C	FINDS XX (ON WALL SPECIFIED) SUCH THAT	\$AR 40
C	ARC RADIUS RC IS TANGENT AT XX	\$AR 50
C	AND PASSES THROUGH NOZZLE AT ANGLE TNOZ	\$AR 60
C	*****	\$AR 70
C		\$AR 80
C	FIND EQUATION OF LINE TANGENT TO X1	\$AR 90
C		\$AR 100
	REAL MNOZ	\$AR 110
	DIMENSION XW(99,2),YW(99,2)	\$AR 120
	DIST(XX,YY,XXC,YYC)=SQRT((ABS(XX-XXC)**2)+(ABS(YY-YYC)**2))	\$AR 130
	MNOZ= TAN(TNOZ)	\$AR 140
	RNOZ= YNOZ - MNOZ * XNOZ	\$AR 150
C		\$AR 160
C	CHOOPPING TECHNIQUE TO FIND CENTRE(X2A) .LT. 0.0	\$AR 170
C	AND CENTRE(X2B) .GT. 0.0	\$AR 180
C		\$AR 190
	DO 200 K =1,NPAIR	\$AR 200
	XXB= XW(K,J)	\$AR 210
	DRCB= CENTRE(XNOZ,YNOZ,MNOZ,BNOZ, XXB,RC,XW,YW,NPAIR,J)	\$AR 220
	IF (DRCB .GT. 0.0) GO TO 210	\$AR 230
200	CONTINUE	\$AR 240
210	CONTINUE	\$AR 250
	XXA= XW(K-2,J)	\$AR 260
	DRCA= CENTRE(XNOZ,YNOZ,MNOZ,BNOZ,XXA,RC,XW,YW,NPAIR,J)	\$AR 270
	FW= DRCA	\$AR 280
C		\$AR 290
C	MODIFIED REGULA FALSI ALGORITHM	\$AR 300
C		\$AR 310
	DO 300 K=1,50	\$AR 320
C		\$AR 330
C	GUESS NEW X2	\$AR 340
C		\$AR 350
	XX= (DRCB*XXA - DRCA*XXB) / (DRCB-DRCA)	\$AR 360
	DRC= CENTRE(XNOZ,YNOZ,MNOZ,BNOZ,XX,RC,XW,YW,NPAIR,J)	\$AR 370
C		\$AR 380
C	CHANGE APPROPRIATE ENDPOINT TO MIDPOINT	\$AR 390
C		\$AR 400
	IF (DRCA*DRC .GT.0.0) GO TO 350	\$AR 410
	XXB= XX	\$AR 420
	DRCB= DRC	\$AR 430
	IF (FW*DRC.GT. 0.0) DRCA=DRCA*.5	\$AR 440

	XR= 1.3	\$OM 220
	FXB= SERIES(DMMID,DMST,JSTSER,XB)	\$OM 230
	DO 505 J=1,50	\$OM 240
	XB= (FXB*XA - FXA*XB)/(FXB-FXA)	\$OM 250
	IF (XB .GT. 2.0) XR= 2.0	\$OM 260
	FXB= SERIES(DMMID,DMST,JSTSER,XB)	\$OM 270
C		\$OM 280
C	CHECK FOR TERMINATION	\$OM 290
C		\$OM 300
	IF (ABS(FXB) .LE. 1.E-04) GO TO 510	\$OM 310
505	CONTINUE	\$OM 320
	WRITE(6,900) JSTART	\$OM 330
900	FORMAT(/7X,22HNO SOLUTION JSTART =,I5/)	\$OM 340
510	CONTINUE	\$OM 350
	RST= XB	\$OM 360
C		\$OM 370
C	SOLVE FOR RFIN IN SERIES	\$OM 380
C		\$OM 390
	XA= 1.5	\$OM 400
	FXA= SERIES(DMMID,DMFN,JFNSE,XA)	\$OM 410
	XB= 1.3	\$OM 420
	FXB= SERIES(DMMID,DMFN,JFNSE,XB)	\$OM 430
	DO 525 J=1,50	\$OM 440
	XB= (FXB*XA - FXA*XB) / (FXB-FXA)	\$OM 450
	IF (XB .GT. 2.0) XB= 2.0	\$OM 460
	FXB= SERIES(DMMID,DMFN,JFNSE,XB)	\$OM 470
C		\$OM 480
C	CHECK FOR TERMINATION	\$OM 490
C		\$OM 500
	IF (ABS(FXB) .LE. 1.E-04 ) GO TO 520	\$OM 510
525	CONTINUE	\$OM 520
	WRITE(6,901)JFIN	\$OM 530
901	FORMAT(/7X,20HNO SOLUTION JFIN =,I5/)	\$OM 540
520	CONTINUE	\$OM 550
	RFIN= XB	\$OM 560
	DO 550 J=1,JSTSER	\$OM 570
550	OMG(JSTART+J)= OMG(JSTART+J-1) + DMST * (RST**(J-1))	\$OM 580
	DO 560 J=1,JFNSE	\$OM 590
560	OMG(JFIN-J)= OMG(JFIN-J+1) - DMFN * (RFIN**(J-1))	\$OM 600
	RETURN	\$OM 610
	END	\$OM 620
	REAL FUNCTION SERIES(SUM,A,N,R)	\$SE 0
	IF (R .NE. 1.0) SERIES= (R**(N+1)-1.0) / (R-1.0) - (SUM/A)	\$SE 10
	IF (R .EQ. 1.0) SERIES= N+1 - (SUM/A)	\$SE 20
	RETURN	\$SE 30
	END	\$SE 40

## REFERENCES

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2. Gerald B. Gilbert and Philip G. Hill "Analysis and Testing of Two-Dimensional Slot Nozzle Ejectors with Variable Area Mixing Sections" NASA CR-2251, May 1973.
3. P. Bradshaw "Effects of Streamline Curvature on Turbulence Flow" AGARDograph No. 169 August 1973 Available from Report Distribution and Storage Unit, NASA, Langley Field, Virginia 23365.
4. H. Schlichting "Boundary Layer Theory" McGraw Hill Book Company, New York, 1968.
5. K.R. Hedges and P.G. Hill "A Finite Difference Method for Compressible Jet Mixing in Converging - Diverging Ducts" Queen's University Thermal Sciences Report No. 3/72 June 1, 1972 Department of Mechanical Engineering, Kingston, Ontario, Canada.
6. Hickman, K.E., Gilbert, G.B., and Carey, J.H., "Analytical and Experimental Investigation of High Entrainment Jet Pumps", NASA CR-1602, July 1970.
7. Kline, S.J., and McClintock, F.A., "Describing Uncertainties in Single Sample Experiments", Mechanical Engineering, 1953.

Table 1  
Mixing Section Dimensions  
(Inches)

J	LOWER WALL		CURV	UPPER WALL		CURV
	X	Y		X	Y	
1	-6.0000	-4.1028	.11979	-3.3270	6.0206	-.03137
2	-5.9500	-4.0266	.12608	-3.3000	5.7680	-.04041
3	-5.9000	-3.9523	.13277	-3.2500	5.3572	-.06887
4	-5.5000	-3.4324	.17553	-3.1250	4.7764	-.11339
5	-5.0000	-2.9342	.21607	-3.0000	4.3953	-.13194
6	-4.5000	-2.5697	.22500	-2.7500	3.8304	-.17358
7	-4.0000	-2.2988	.22416	-2.5000	3.4911	-.20405
8	-3.5000	-2.1047	.18402	-2.2500	3.0633	-.17665
9	-3.0000	-1.9528	.15850	-2.0000	2.7705	-.16650
10	-2.5000	-1.8637	.13904	-1.5000	2.2963	-.18696
11	-2.3500	-1.8403	.14043	-1.0000	1.9315	-.20381
12	-2.3125	-1.8349	.14075	-.5000	1.6544	-.20305
13	-2.3125	-1.8390	.00000	.0000	1.4449	-.21235
14	-2.0000	-1.7783	.00000	.5000	1.2971	-.21607
15	-1.0000	-1.5840	.00000	.6250	1.2689	-.22019
16	.0000	-1.3897	.00000	.7500	1.2444	-.22389
17	.7500	-1.2440	.00000	.7500	1.2443	-.08962
18	.7500	-1.2443	.08962	.8000	1.2350	-.08984
19	.8000	-1.2350	.08984	.8500	1.2260	-.09007
20	.8500	-1.2260	.09007	1.0000	1.2009	-.07313
21	1.0000	-1.2009	.07313	1.5000	1.1298	-.07108
22	1.5000	-1.1298	.07108	2.0000	1.0773	-.06664
23	2.0000	-1.0773	.06664	2.5000	1.0423	-.05432
24	2.5000	-1.0423	.05432	3.0000	1.0215	-.04063
25	3.0000	-1.0205	.04063	3.4000	1.0095	-.01971
26	3.4000	-1.0095	.01971	3.4500	1.0082	-.01971
27	3.4500	-1.0082	.01971	3.5000	1.0069	-.01971
28	3.5000	-1.0069	.01971	3.5000	1.0070	.00000
29	3.5000	-1.0070	.00000	4.0000	.9984	.00000
30	4.0000	-.9984	.00000	5.0000	.9811	.00000
31	5.0000	-.9811	.00000	6.0000	.9639	.00000
32	6.0000	-.9639	.00000	7.0000	.9466	.00000
33	7.0000	-.9466	.00000	7.5000	.9380	.00000
34	7.5000	-.9380	.00000	7.5000	.9380	.00000
35	7.5000	-.9380	.00000	8.0000	.9380	.00000
36	8.0000	-.9380	.00000	9.0000	.9380	.00000
37	9.0000	-.9380	.00000	10.0000	.9380	.00000
38	10.0000	-.9380	.00000	10.5000	.9380	.00000
39	10.5000	-.9380	.00000	10.5000	.9380	.00000
40	10.5000	-.9380	.00000	11.0000	.9642	.00000
41	11.0000	-.9642	.00000	12.0000	1.0166	.00000
42	12.0000	-1.0166	.00000	14.0000	1.1214	.00000
43	14.0000	-1.1214	.00000	16.0000	1.2262	.00000
44	16.0000	-1.2262	.00000	18.0000	1.3310	.00000
45	18.0000	-1.3310	.00000	20.0000	1.4358	.00000
46	20.0000	-1.4358	.00000	22.0000	1.5406	.00000
47	22.0000	-1.5406	.00000	23.0000	1.5930	.00000

Table 2  
Variation of Individual Integrated Traverse Mass  
Flows for Each Test Run

Run	Variation of Integrated Traverse Mass Flow Rate Around An Average Value	Number of Traverses
1	-4.1%, +7.3%	4
2	-3.1%, +5.1%	10
3	-2.1%, +2.2%	3
4	-3.7%, +5.9%	4
5	-2.7%, +3.7%	4
6	-1.6%, +3.0%	3
7	-3.5%, +4.6%	5
8	-3.4%, +3.5%	5
Average	-3.0%, +4.4%	



Table 3

## Traverse Locations and Data Summary

	Maximum Flow - Atmospheric Discharge				Reduced Flow - Back Pressure at Discharge			
Test Number	2	3	6	7	1	4	5	8
Nozzle Position (See Fig. 3)	1	2	3	4	1	2	3	4
Nozzle Angle	22.5°	45°	45°	67.5°	22.5°	45°	45°	67.5°
Nozzle	meter	0.020	0.020	0.034	0.018	0.020	0.020	0.018
Spacing	inch	0.80	0.80	1.32	0.70	0.80	1.32	0.70
Traverse Location †								
s = -0.070 meters, - 2.75 inches				P <sub>max</sub>				P <sub>max</sub>
s = -0.024 meters, - 0.95 inches		P <sub>max</sub>	P <sub>max</sub>	P, P <sub>max</sub>		P <sub>max</sub>	P <sub>max</sub>	P, P <sub>max</sub>
x = +0.013 meters, + 0.50 inches	P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>
x = +0.038 meters, + 1.50 inches					P <sub>max</sub>	P <sub>max</sub>	P <sub>max</sub>	
x = +0.064 meters, + 2.50 inches								
x = +0.114 meters, + 4.50 inches	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>
x = +0.165 meters, + 6.50 inches				U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>
x = +0.254 meters, +10.00 inches	U, P <sub>max</sub>							
x = +0.318 meters, +12.50 inches		U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>
x = +0.419 meters, +16.50 inches				U, P <sub>max</sub>				
x = +0.521 meters, +20.50 inches	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>	U, P <sub>max</sub>

Axial Wall Static Pressure Profiles (Figures 16 to 19)

P = Vertical Total Pressure Profile (Figures 21 and 22)

P<sub>max</sub> = Maximum Pressure Location (Figures 23 to 30)

U = Vertical Velocity Profile (Figures 31 to 38)

Table 4

## Summary of Experimental Test Conditions and Flow Rates

Nozzle Throat Area = .9688 in<sup>2</sup>

Mixing Section Throat Size = 1.875 inch

Nozzle Pressure = 36.70 psia (constant)

Run No	Nozzle Temp °R $T_N$	Nozzle Throat Coefficient $C_N$	Barometric Pressure in Hg $P_b$	Atmospheric Temp °R $T_a$	Measured Nozzle Flow Rate lb/(sec-in) $W_N$	Mixing Section Flow lb/(sec-in) $W_m$	Secondary Flow Rate lb/(sec-in) $W_s$	Flow Ratio $W_s/W_N$
1	556	.959	30.02	527	.0967	.4995	.4028	4.17
2	575	.966	29.96	540	.0952	.4332	.3380	3.55
3	564	.961	29.89	536	.0952	.4627	.3675	3.86
4	561	.957	29.91	536	.0955	.5238	.4283	4.48
5	576	.959	30.16	544	.0889	.5208	.4319	4.86
6	561	.959	30.00	545	.0958	.4220	.3262	3.41
7	576	.960	30.13	543	.0955	.4390	.3435	3.60
8	582	.954	30.04	548	.0932	.5242	.4310	4.62

Table 5

Comparison of Experimental and Analytical Mass Flow Rates

Run No.	Mixing Section Mass Flow Rate From Traverse Data lb/(sec-m) ①	Mixing Section Mass Flow Rate From Orifice Data lb/(sec-in) ②	Percent Difference in Measured Data ① - ② ①	Analytical Mass Flow for Best Static Pressure Match lb/(sec-m) ③	Comparison of Traverse to Analytical Mass Flow ① - ③ ①	Comparison of Orifice to Analytical Mass Flow ② - ③ ②
1	.4995	-	-	.4905	+ 1.8%	-
2	.4332	.3776	+ 12.8%	.4086	+ 5.7%	- 8.2%
3	.4627	.3954	+ 14.5%	.4175	+ 9.8%	- 5.6%
4	.5238	-	-	.4930	+ 5.9%	-
5	.5208	-	-	.5100	+ 2.1%	-
6	.4220	.3886	+ 7.9%	.4117	+ 2.4%	- 5.9%
7	.4390	.3846	+ 12.4%	.3960	+ 9.8%	- 3.0%
8	.5242	-	-	.4840	+ 7.7%	-

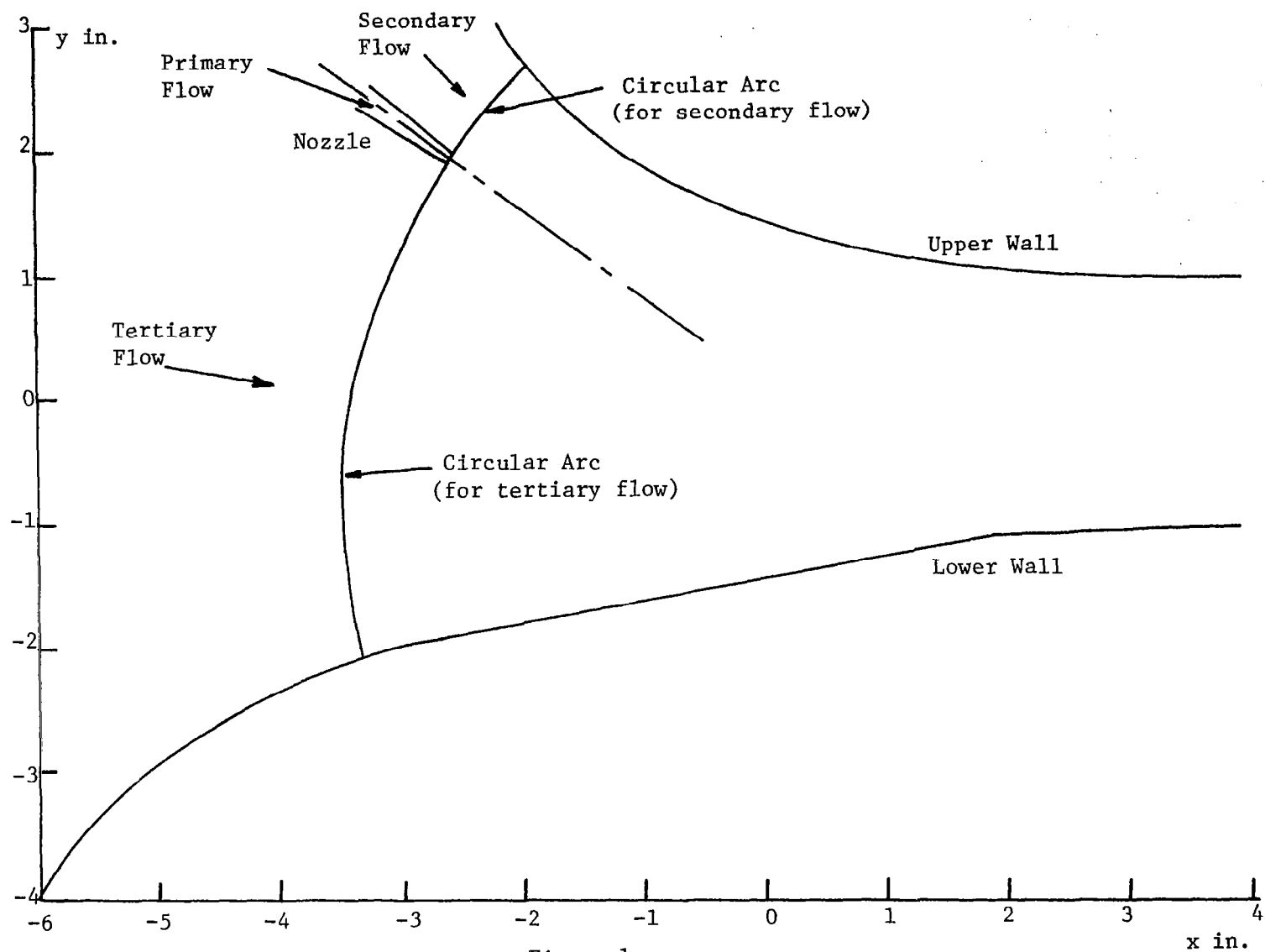
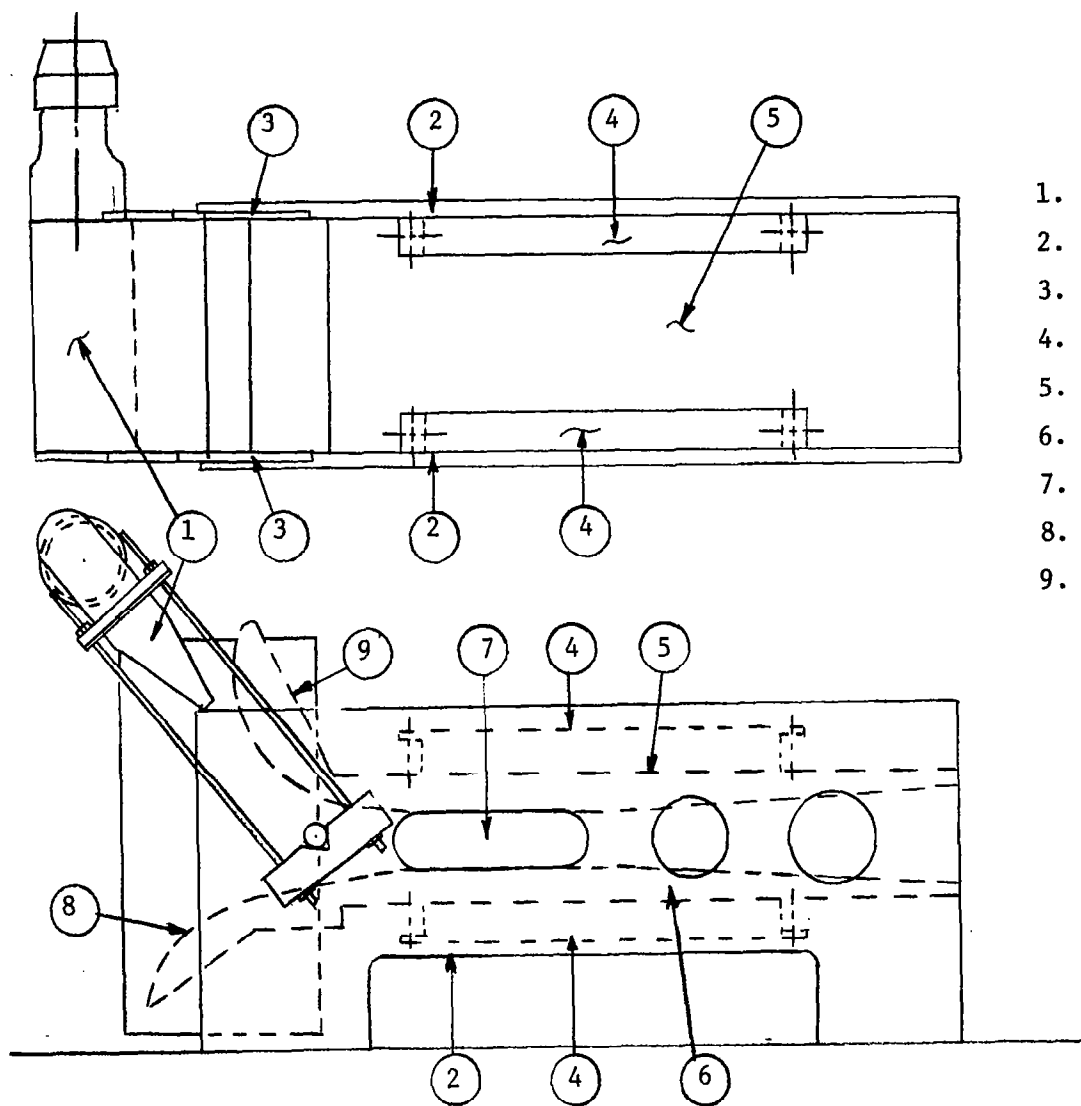


Figure 1

Initial Line for Flow Calculation



1. Primary Nozzle
2. Side Plates
3. Nozzle Positioning Plates
4. Suction Plenum
5. Upper Contoured Plate
6. Lower Contoured Plate
7. Plexiglass Windows
8. Inlet Bellmouth
9. Coanda Surface

Figure 2

Assembly Sketch of Two Dimensional Ejector Test Rig

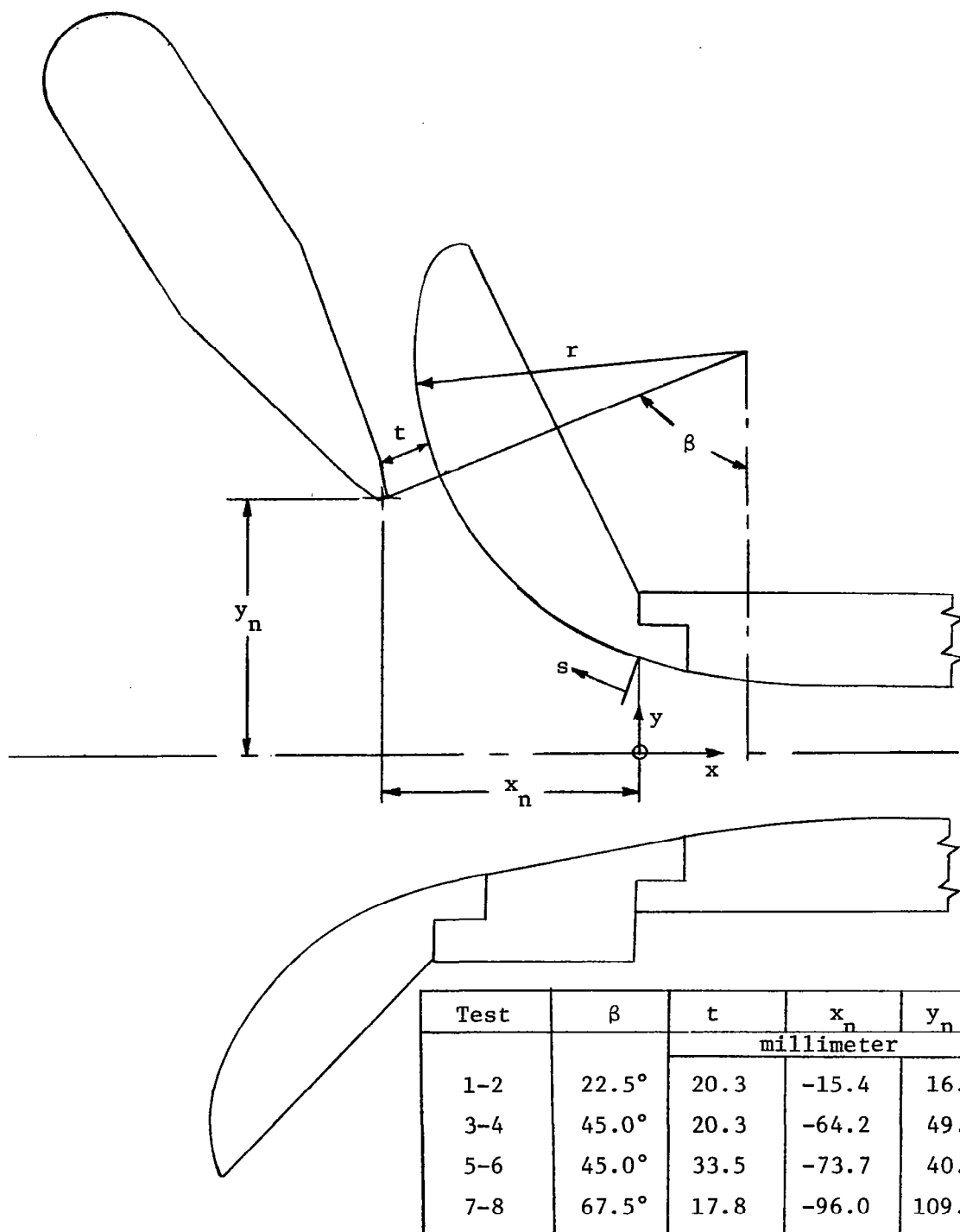


Figure 3

Drawing of Nozzle Coordinate System  
and Positioning Data for Eight Tests

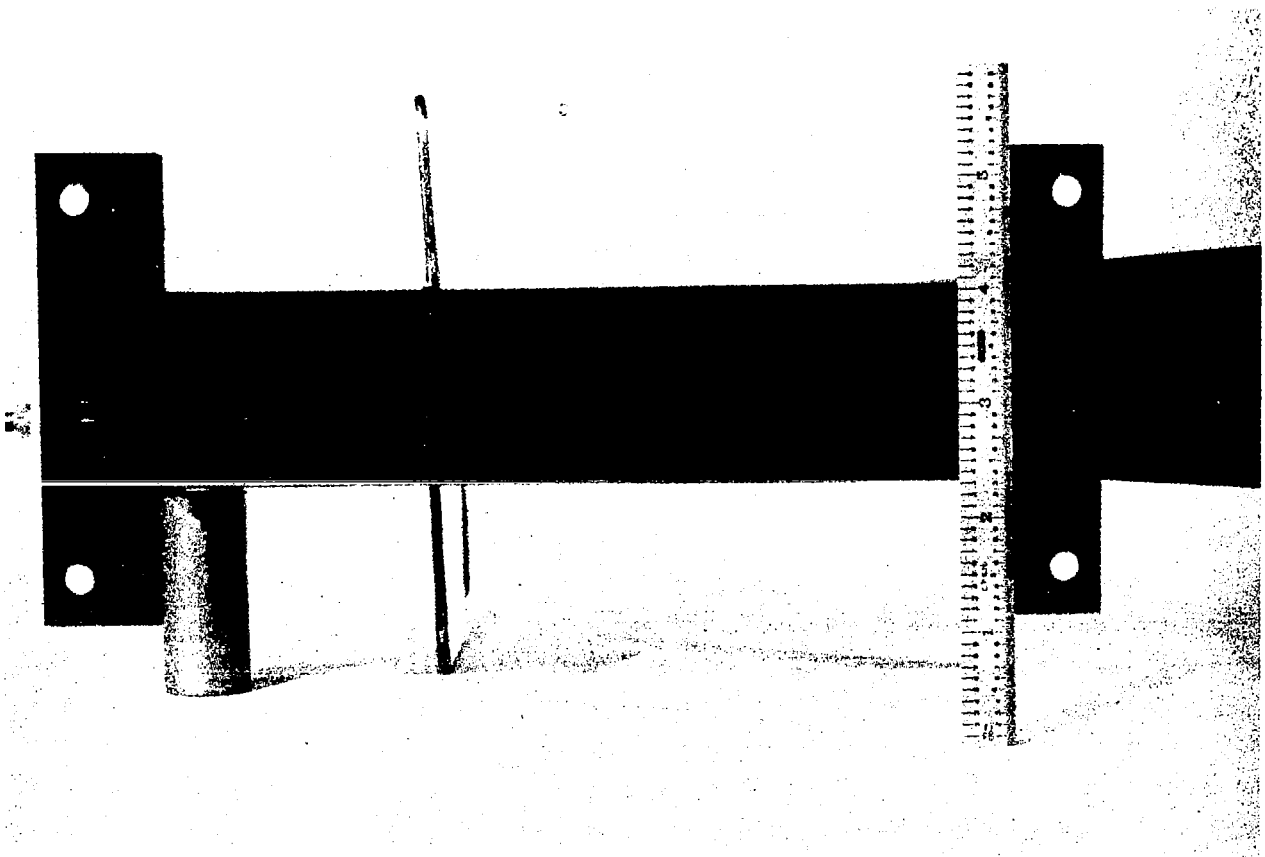


Figure 4  
Primary Nozzle

1. Nozzle
2. Mixing Section Side Plate
3. Top Contoured Plate
4. Bottom Contoured Plate
5. Table Top
6. Screen
7. Solid Side Plates
8. Nozzle Positioning Plates

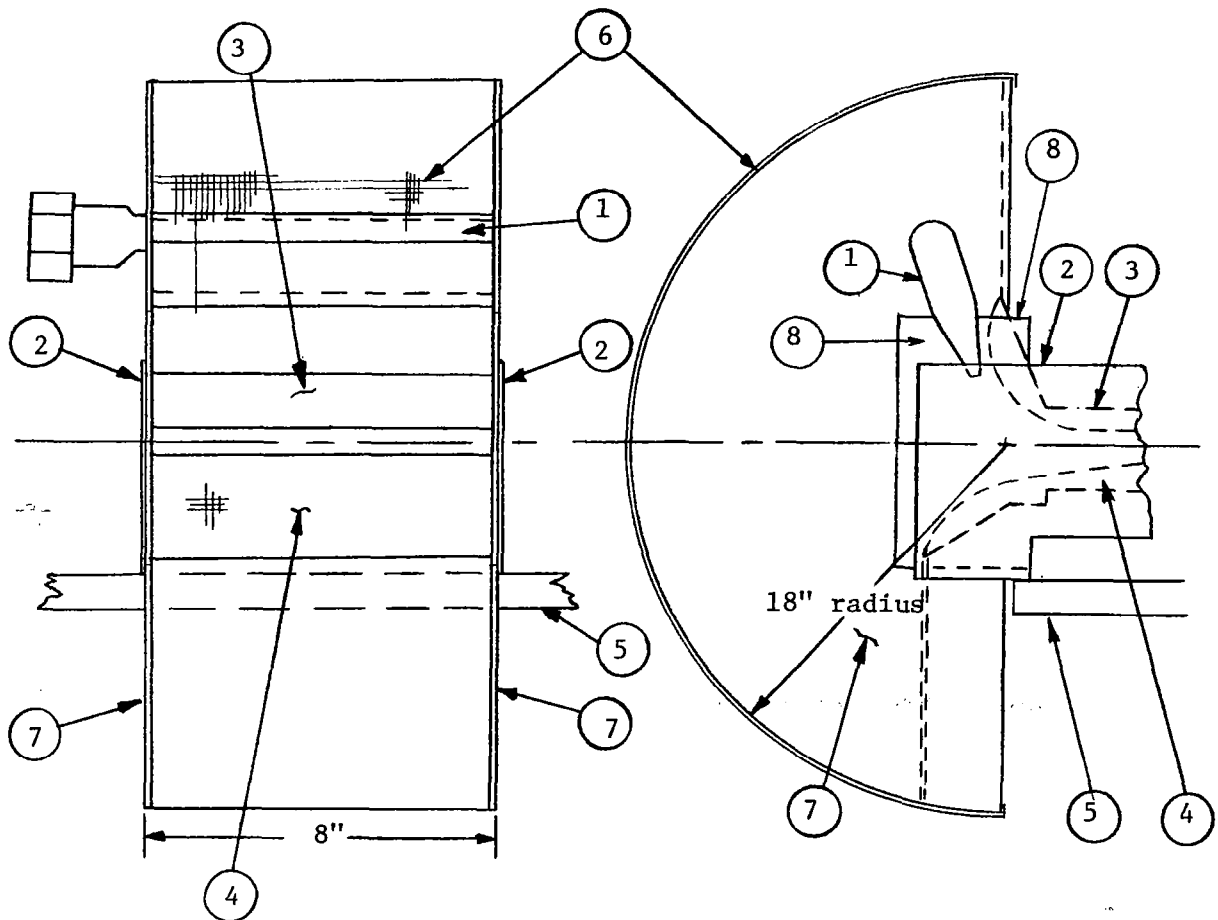


Figure 5

Extended Inlet on Ejector Test Rig



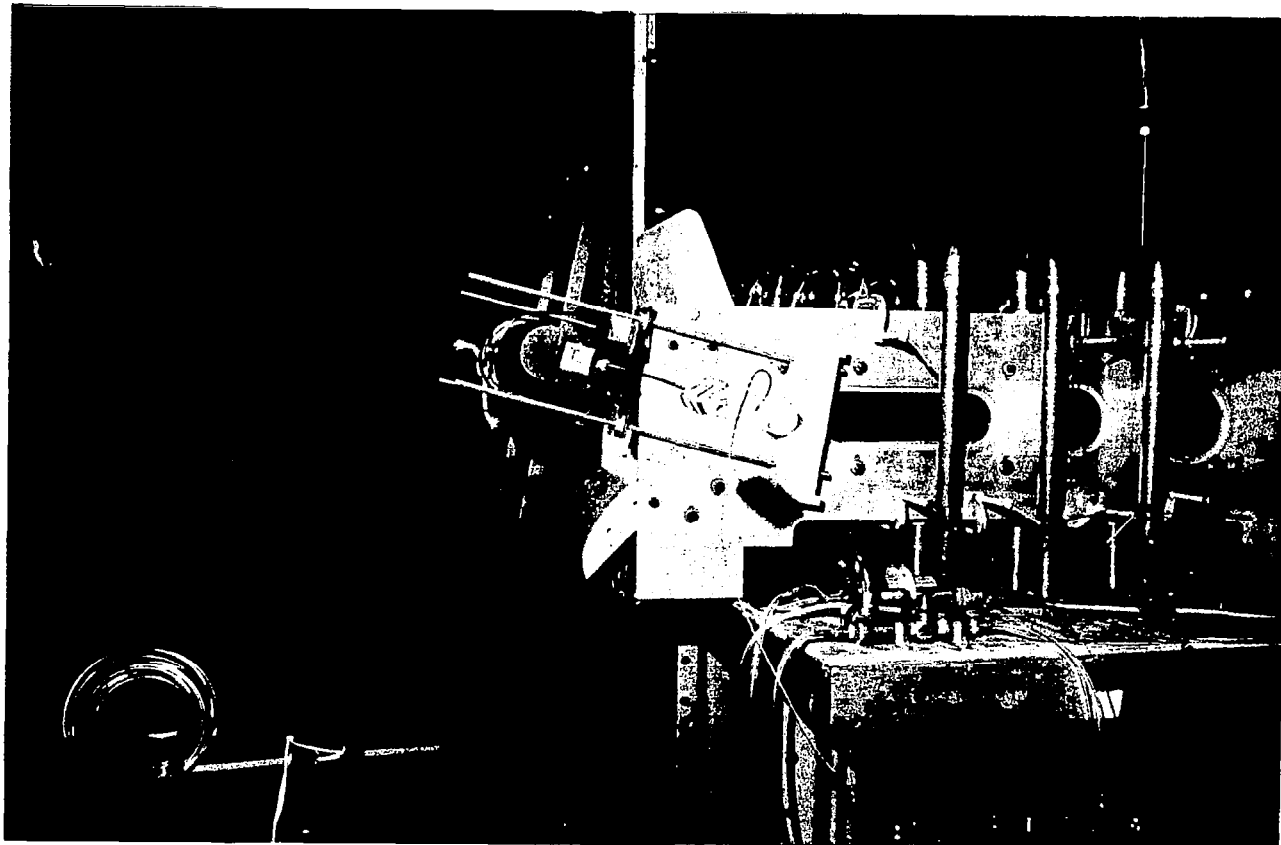


Figure 6

Side View of Test Section Showing  
Nozzle Positioned at  
22.5° and 0.020 meters (0.80 inches) spacing

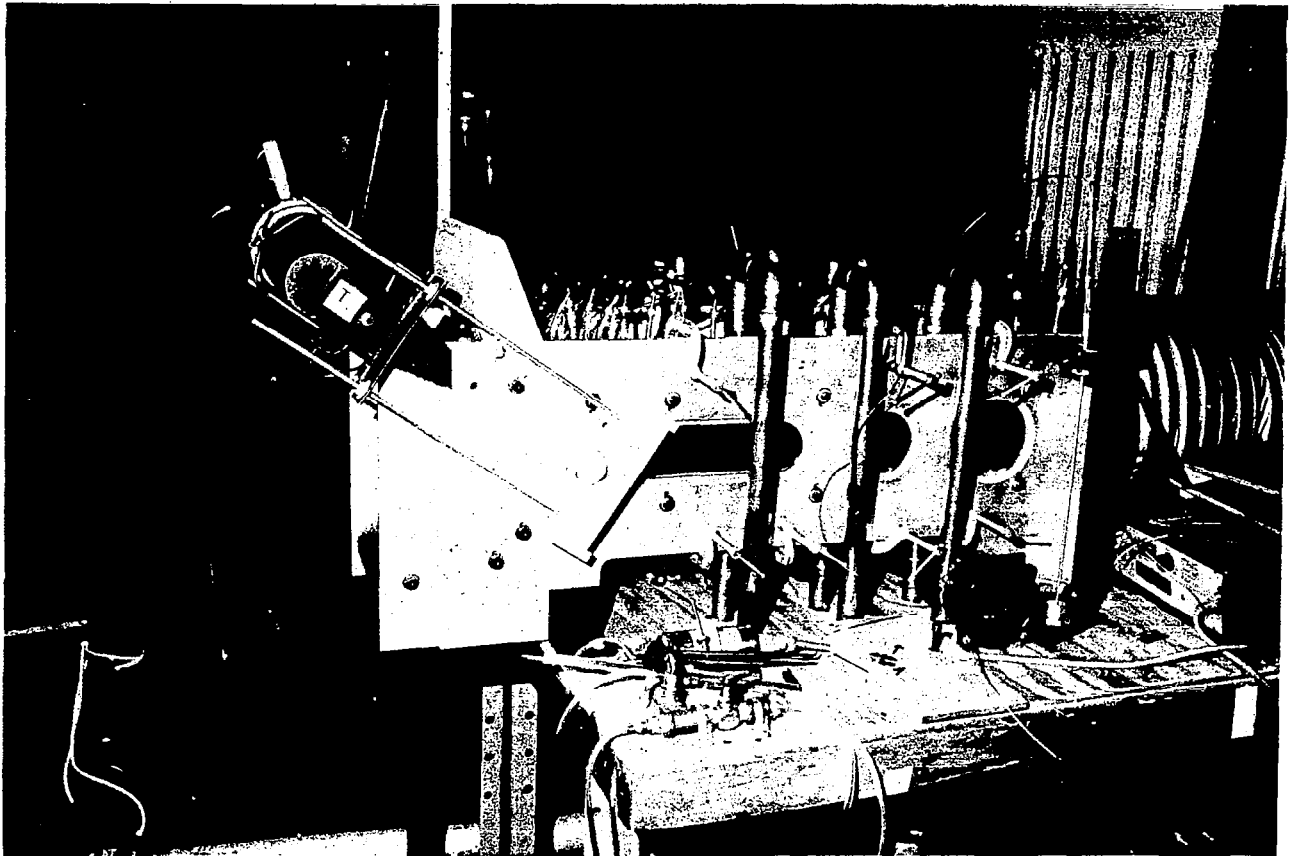


Figure 7

Side View of Test Section Showing  
Nozzle Positioned at  
45° and 0.034 meters (1.32 inches) spacing

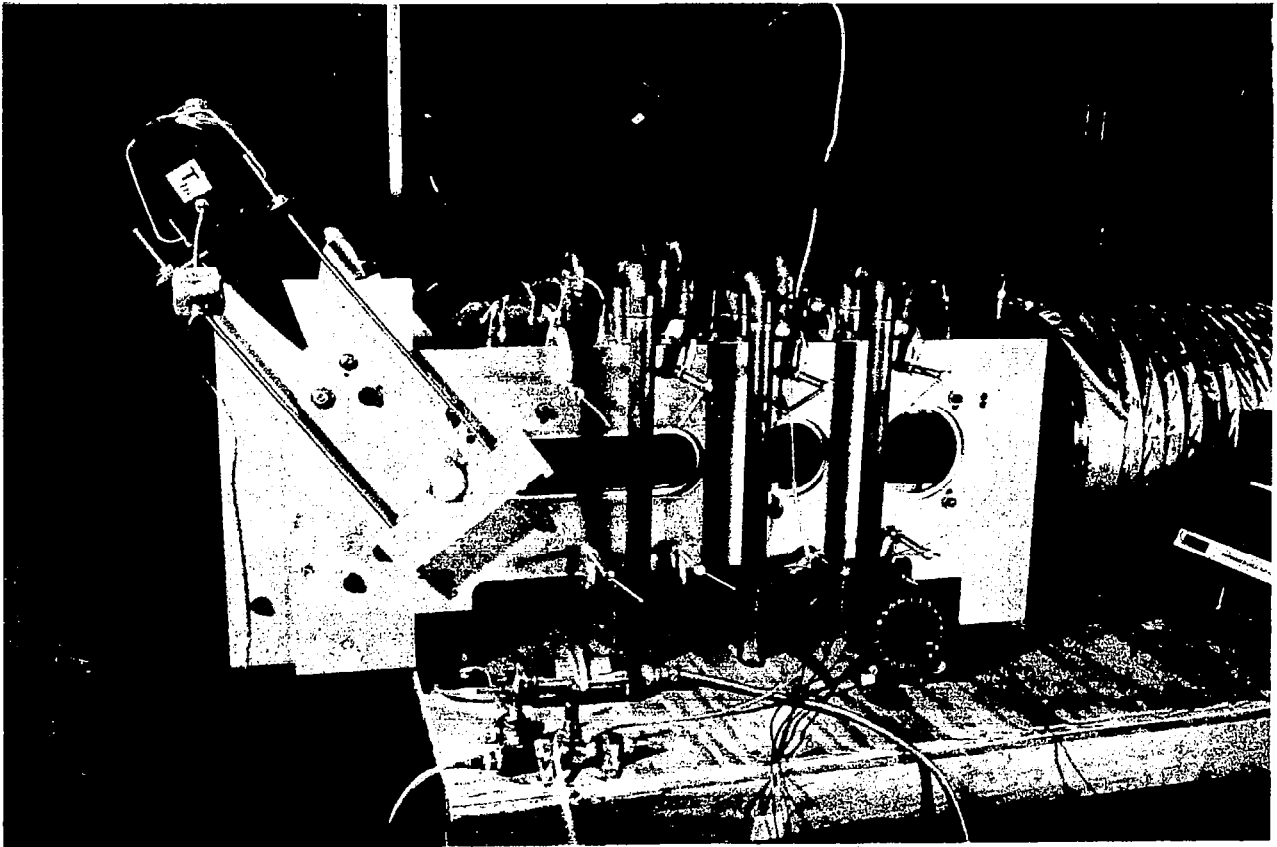


Figure 8

Side View of Test Section Showing  
Nozzle Positioned at  
 $67.5^\circ$  and 0.018 meters (0.7 inches) spacing

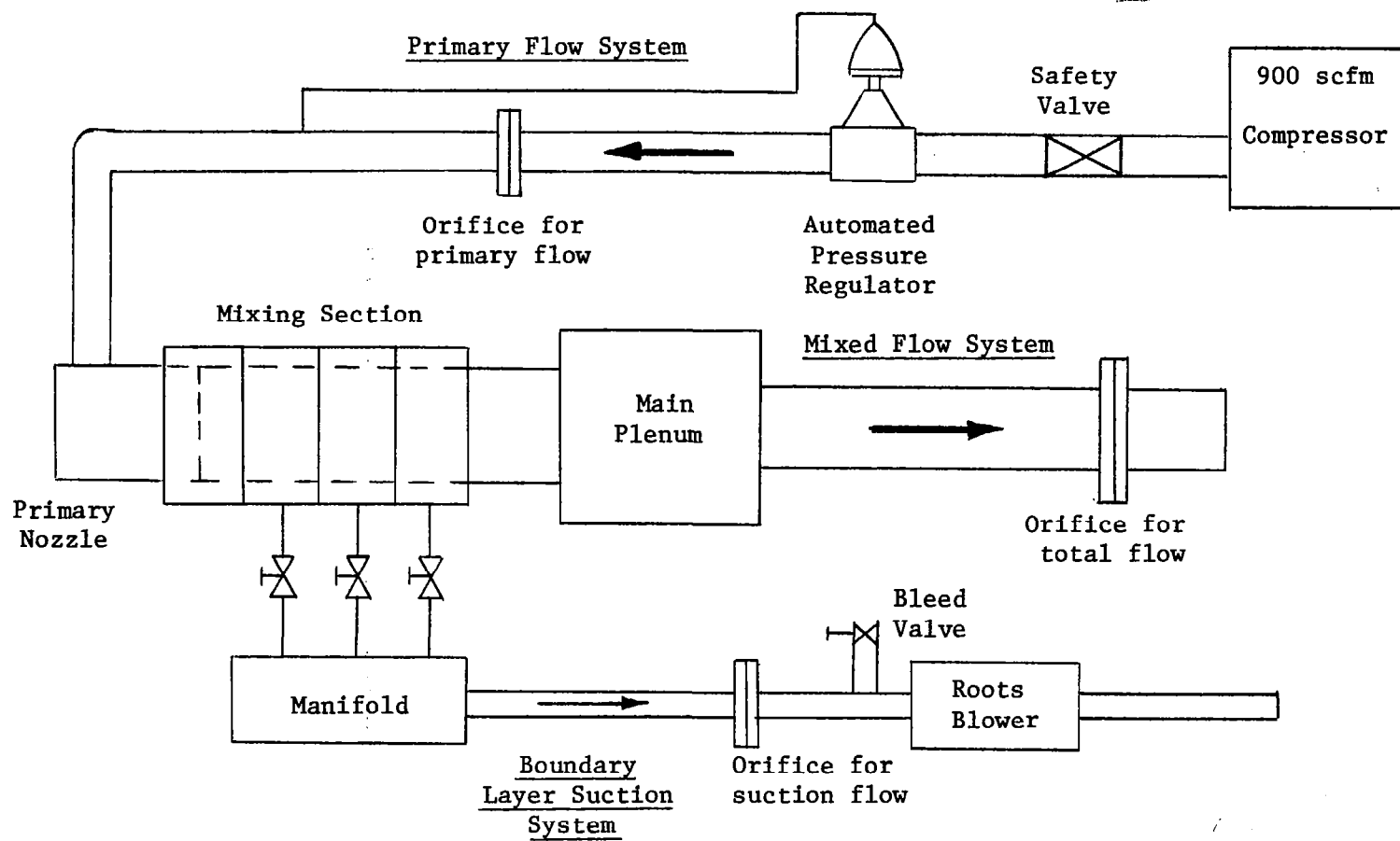


Figure 9

Schematic of Experimental Layout  
for Reduced Flow Conditions

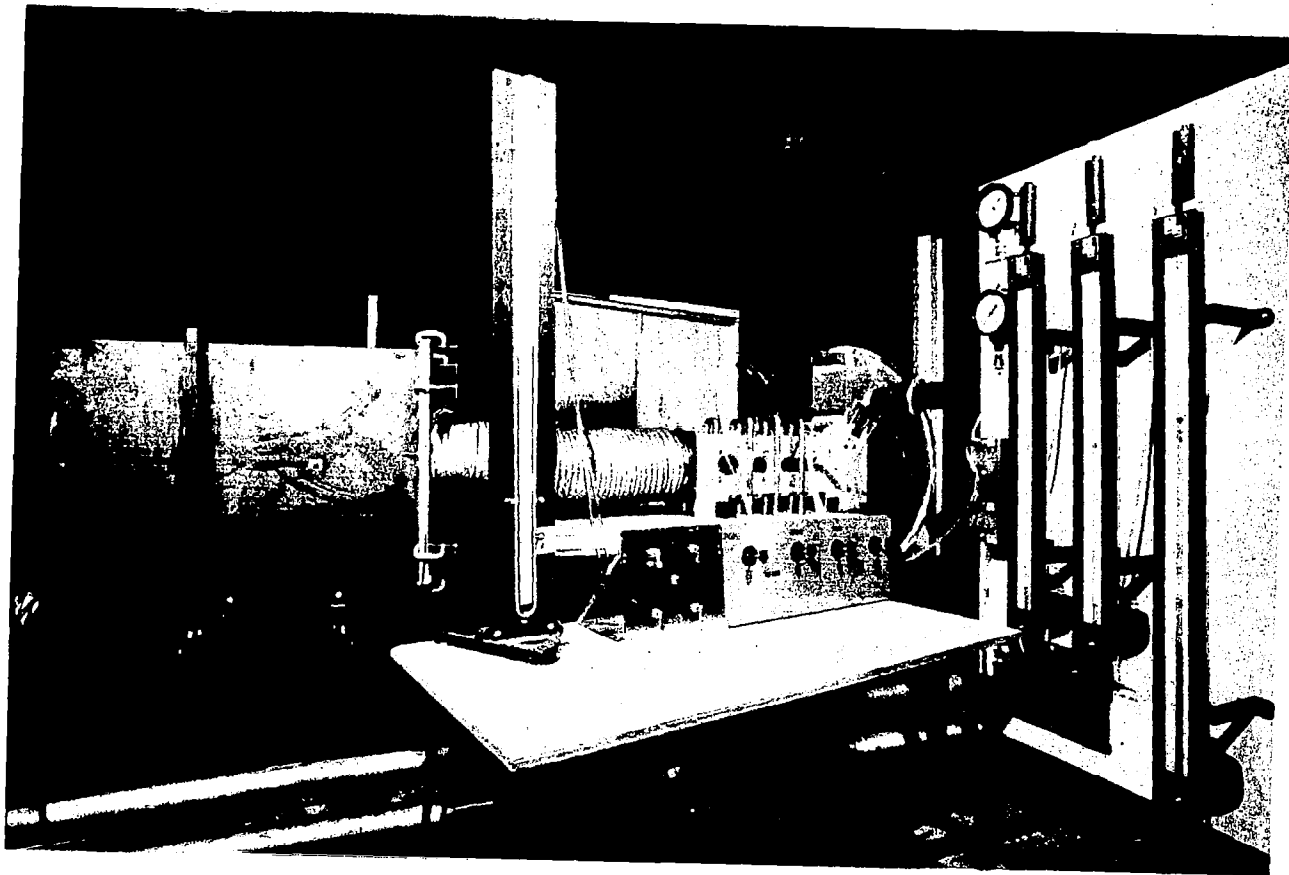


Figure 10  
View of Test Facility Showing Plenum,  
Mixing Section, Manometer Board, and  
Pressure Sampling Valves

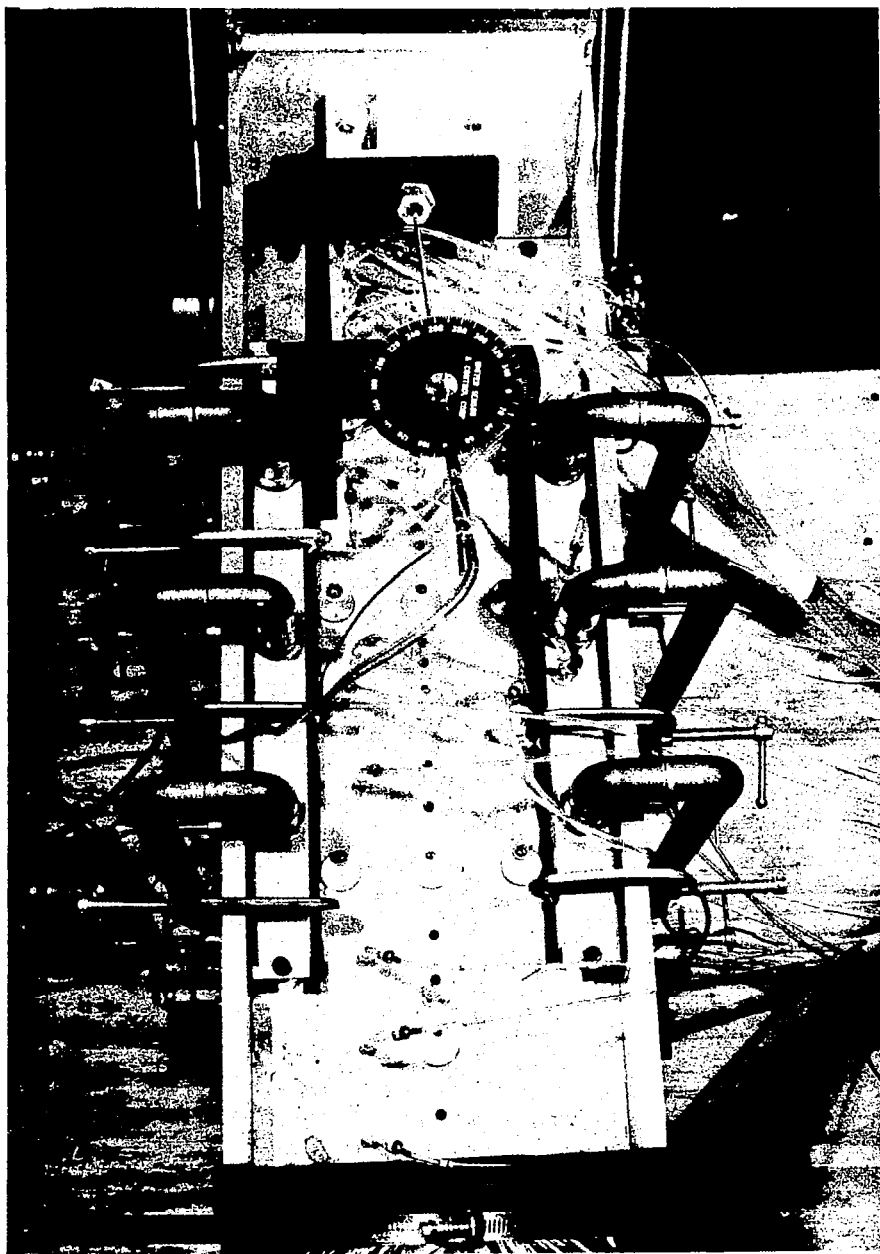


Figure 11

Top View of Mixing Section Showing  
Static and Traversing Tap Locations  
Traversing Probe is Positioned at  
"A" Location on Coanda Surface

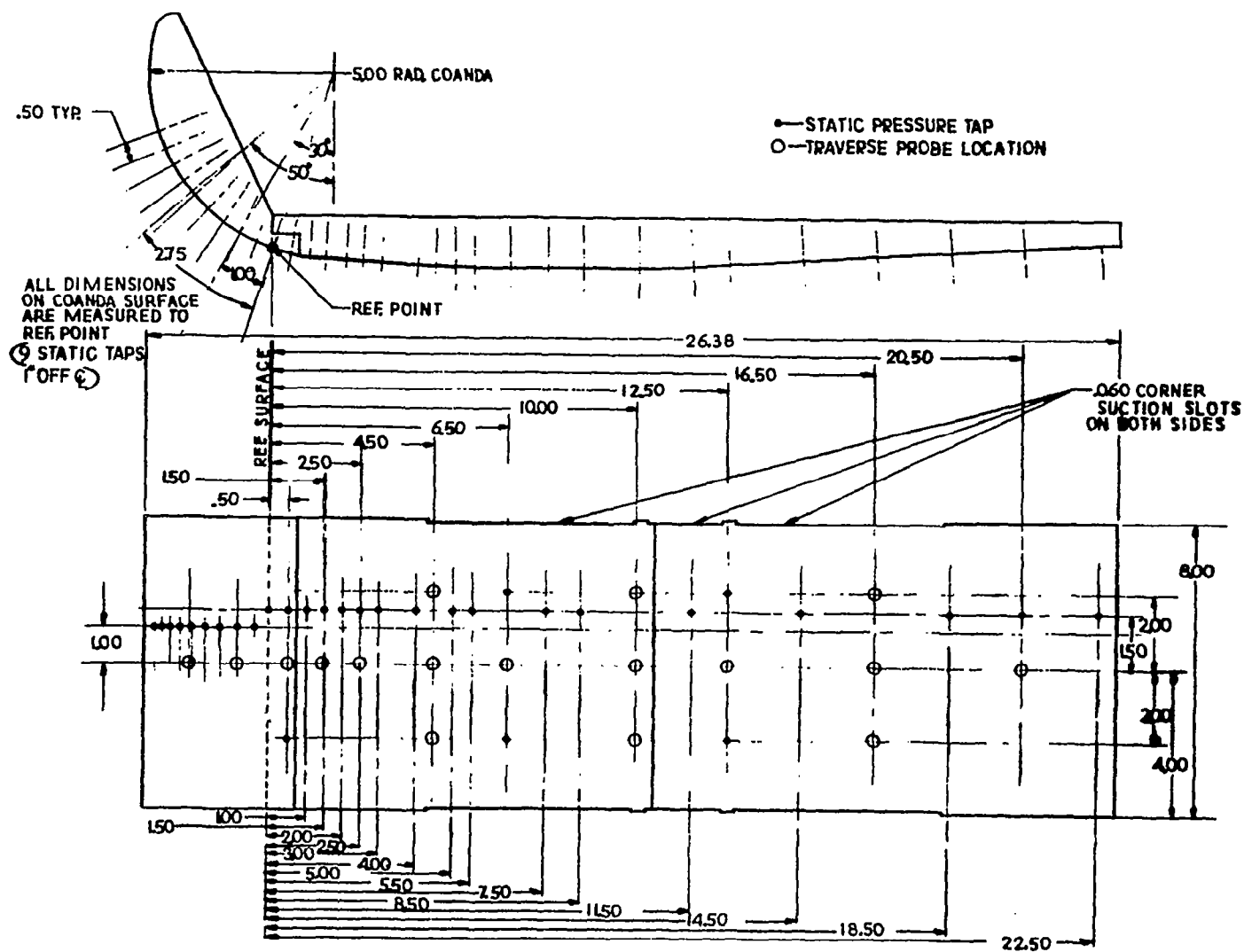


Figure 12

Mixing Section Traverse and Static Pressure Tap Locations

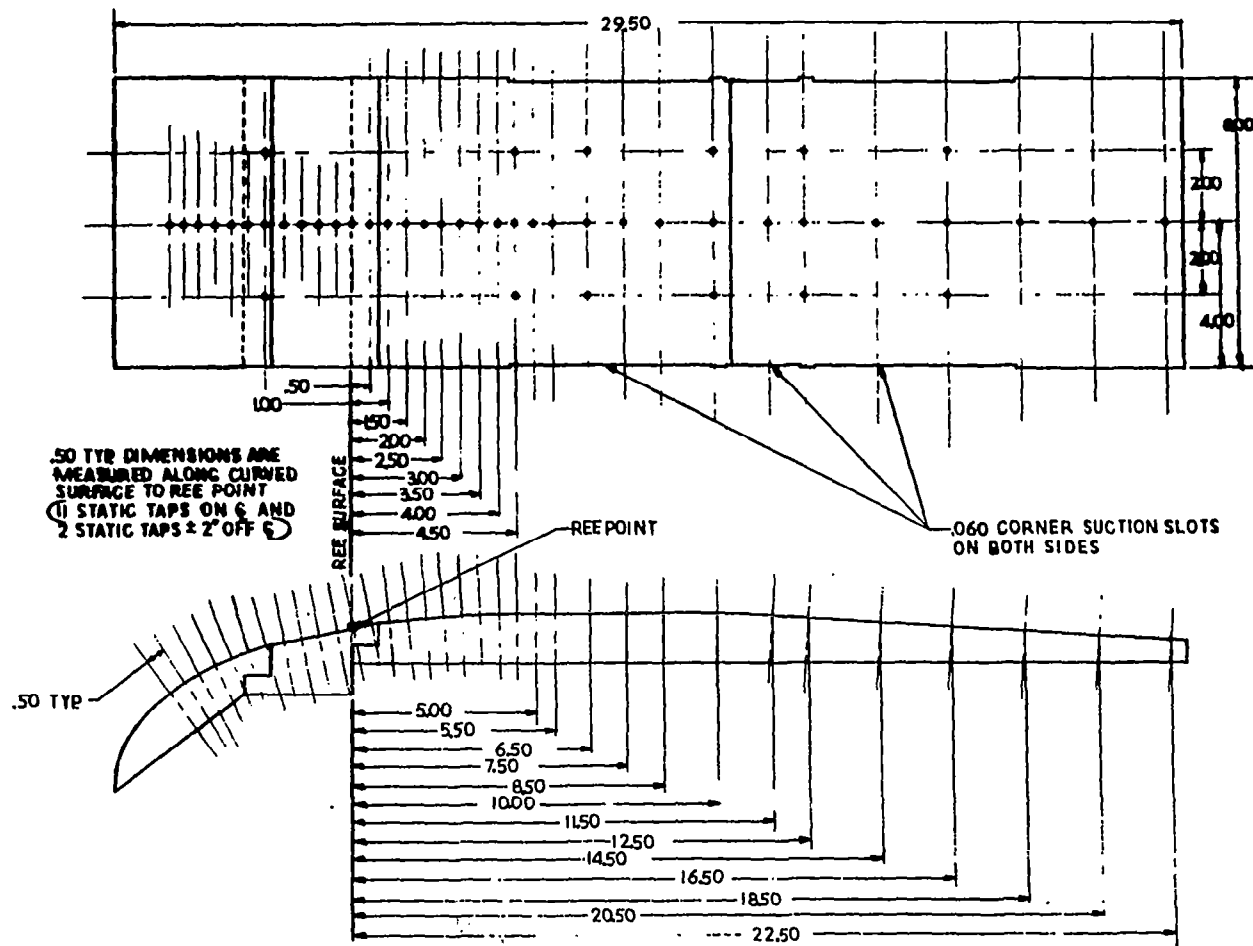


Figure 13

Mixing Section Static Pressure Tap Locations



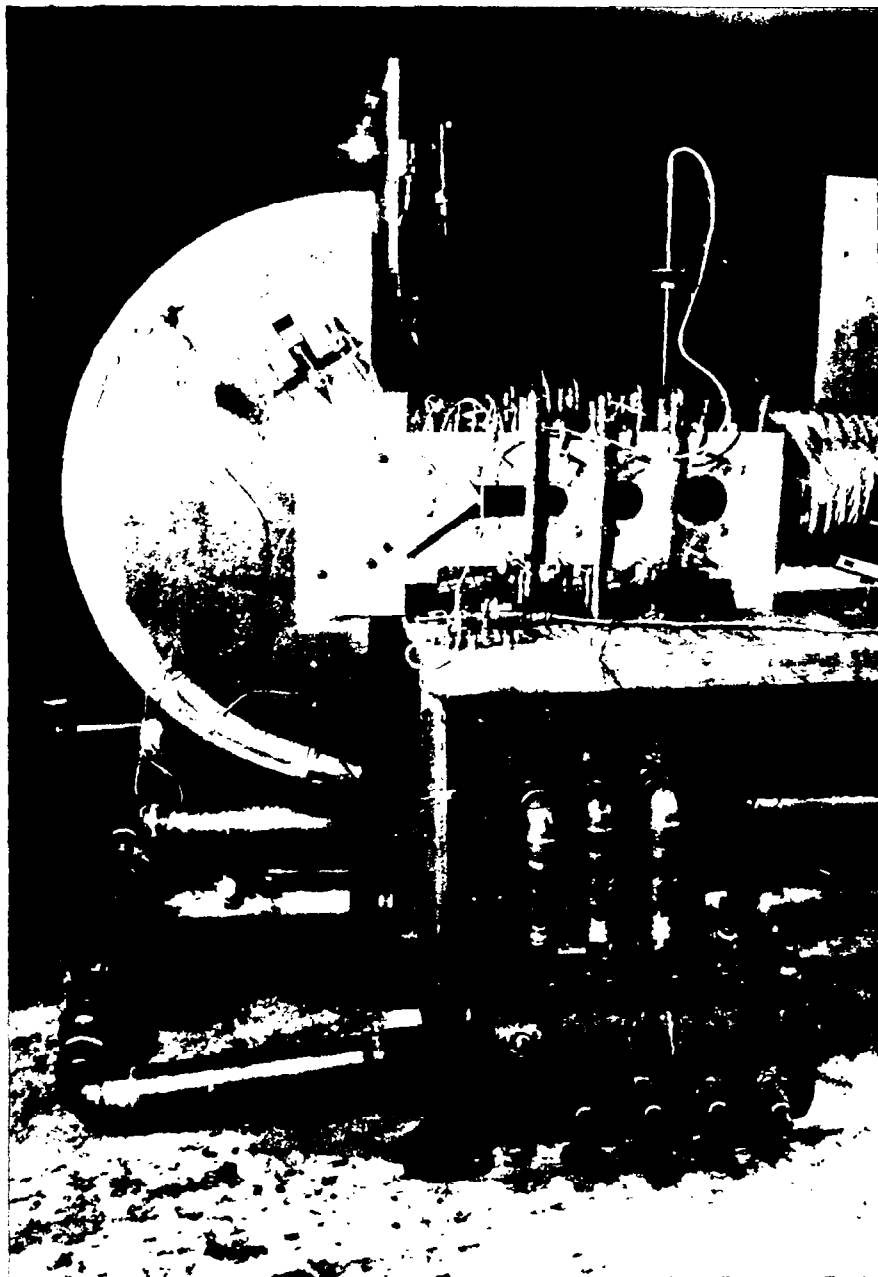


Figure 14

Side View of Test Section Showing  
Extended Inlet, Mixing Section and Suction  
System Piping and Manifold

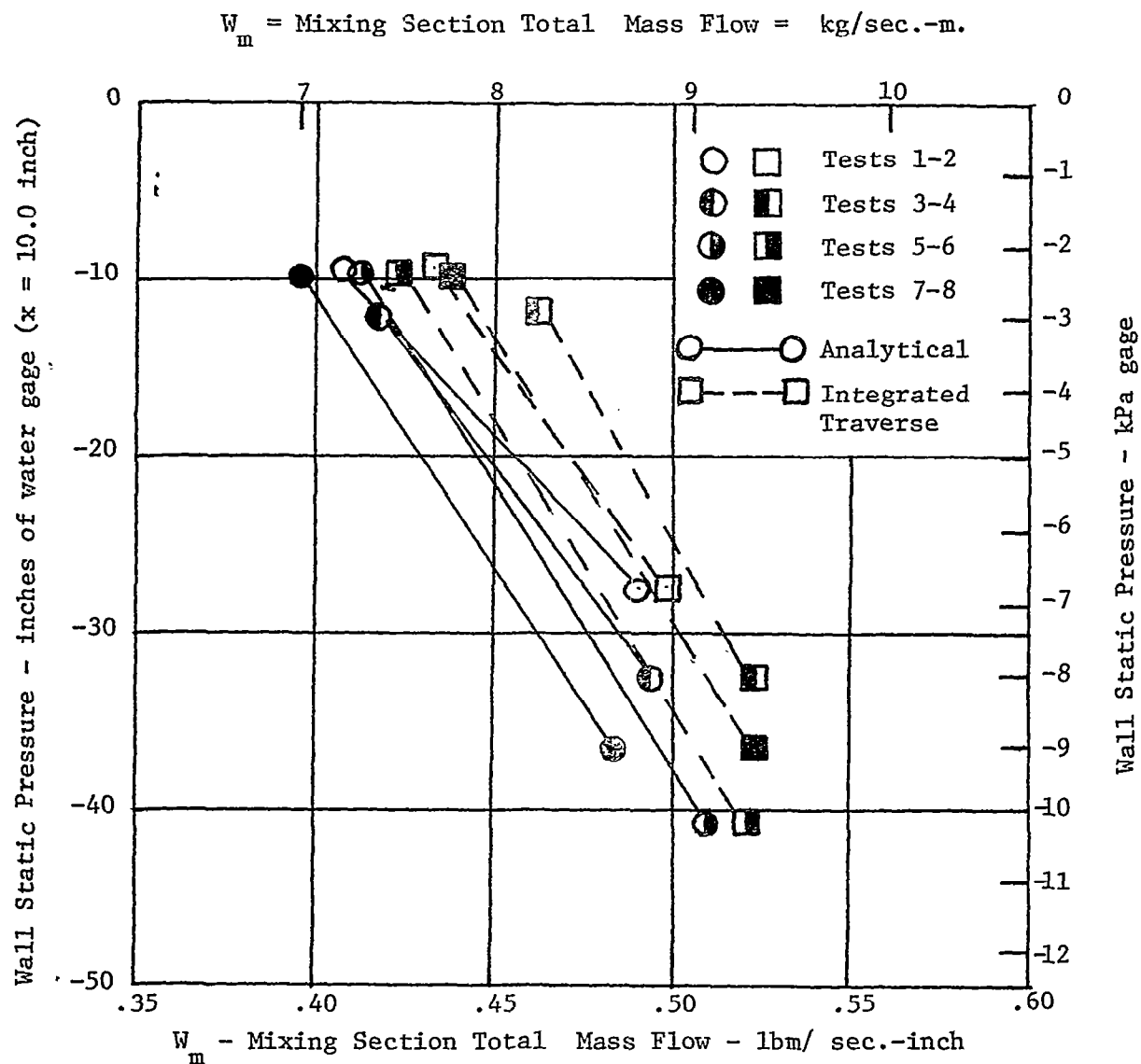


Figure 15

Comparison of Experimental and Analytical Flow Rates

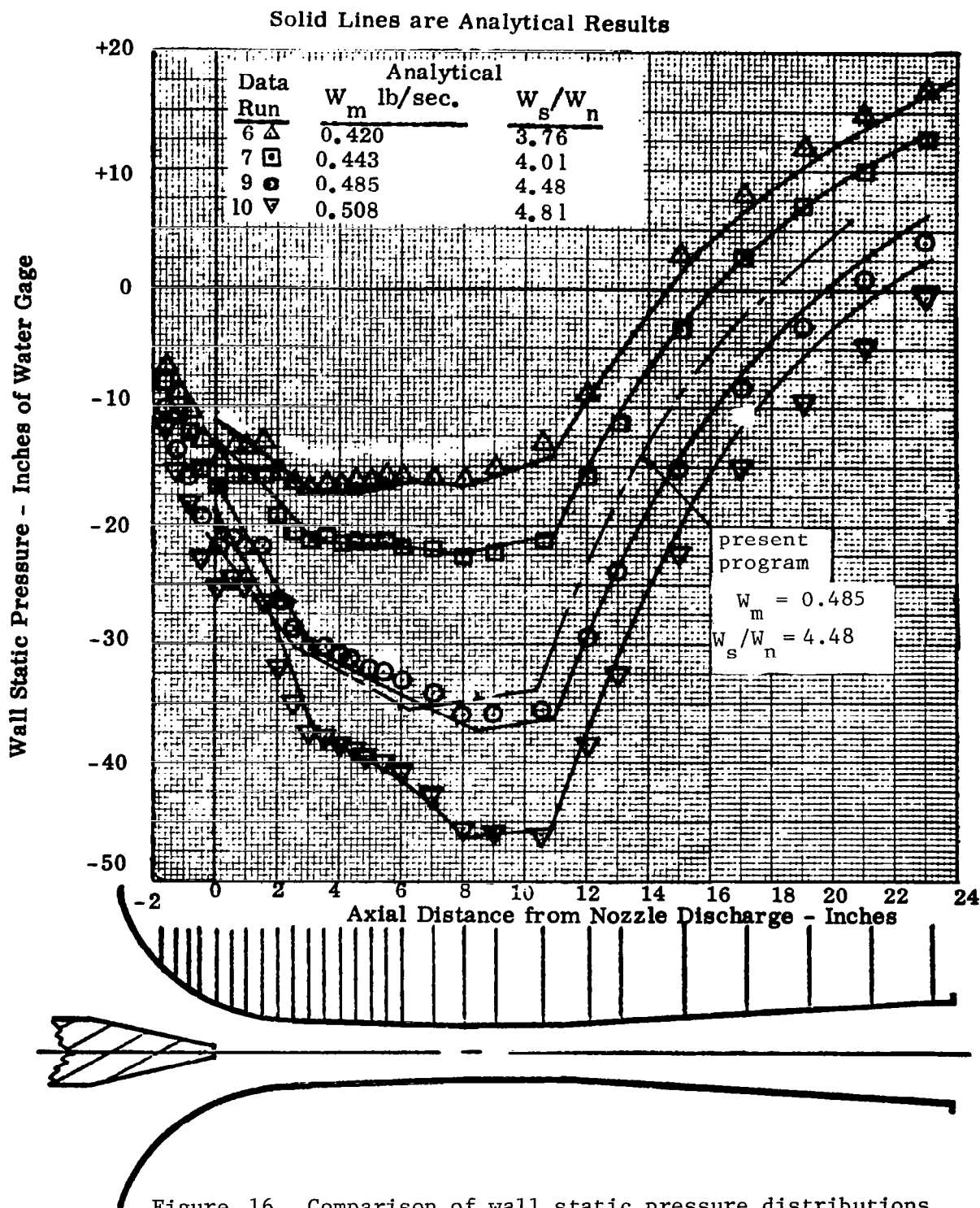


Figure 16 Comparison of wall static pressure distributions in a symmetrical mixing section with a 1.875" throat computed with the present streamline coordinate program and with the program from CR-2251

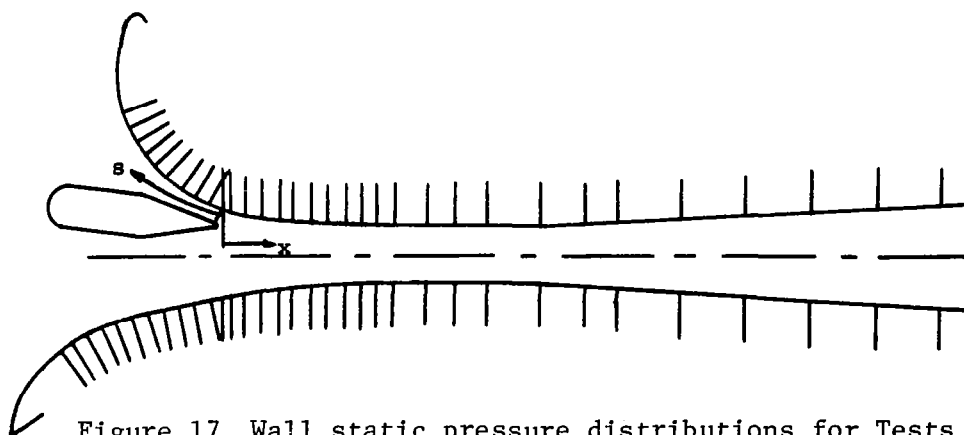
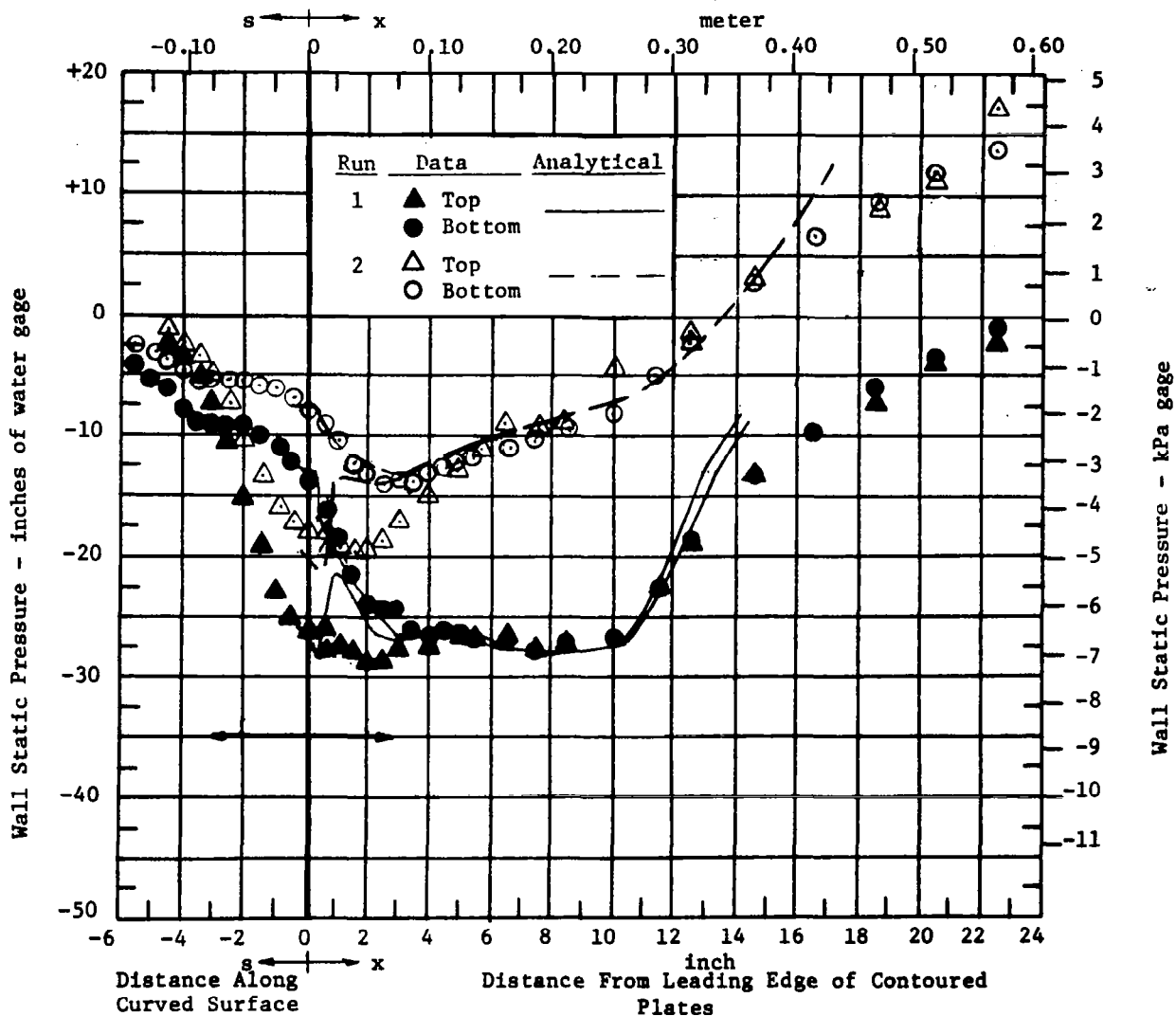


Figure 17 Wall static pressure distributions for Tests 1 and 2. Nozzle positioned at 22.5° and 0.020 meters (0.80 inches) spacing

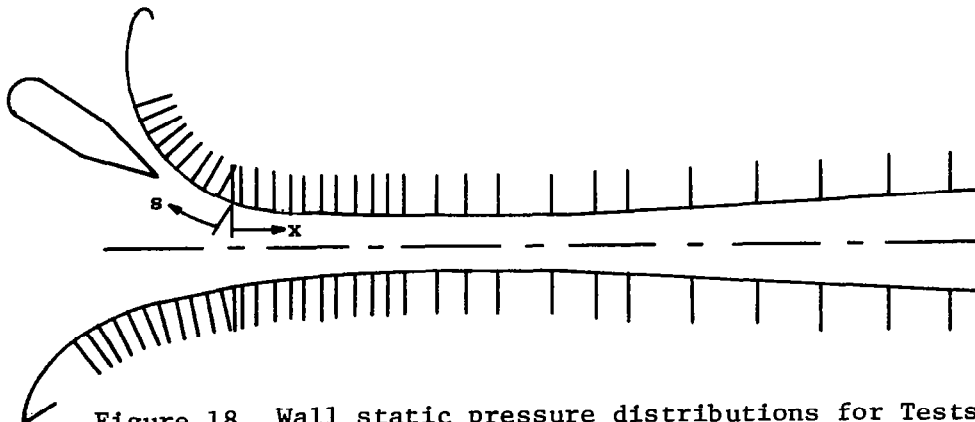
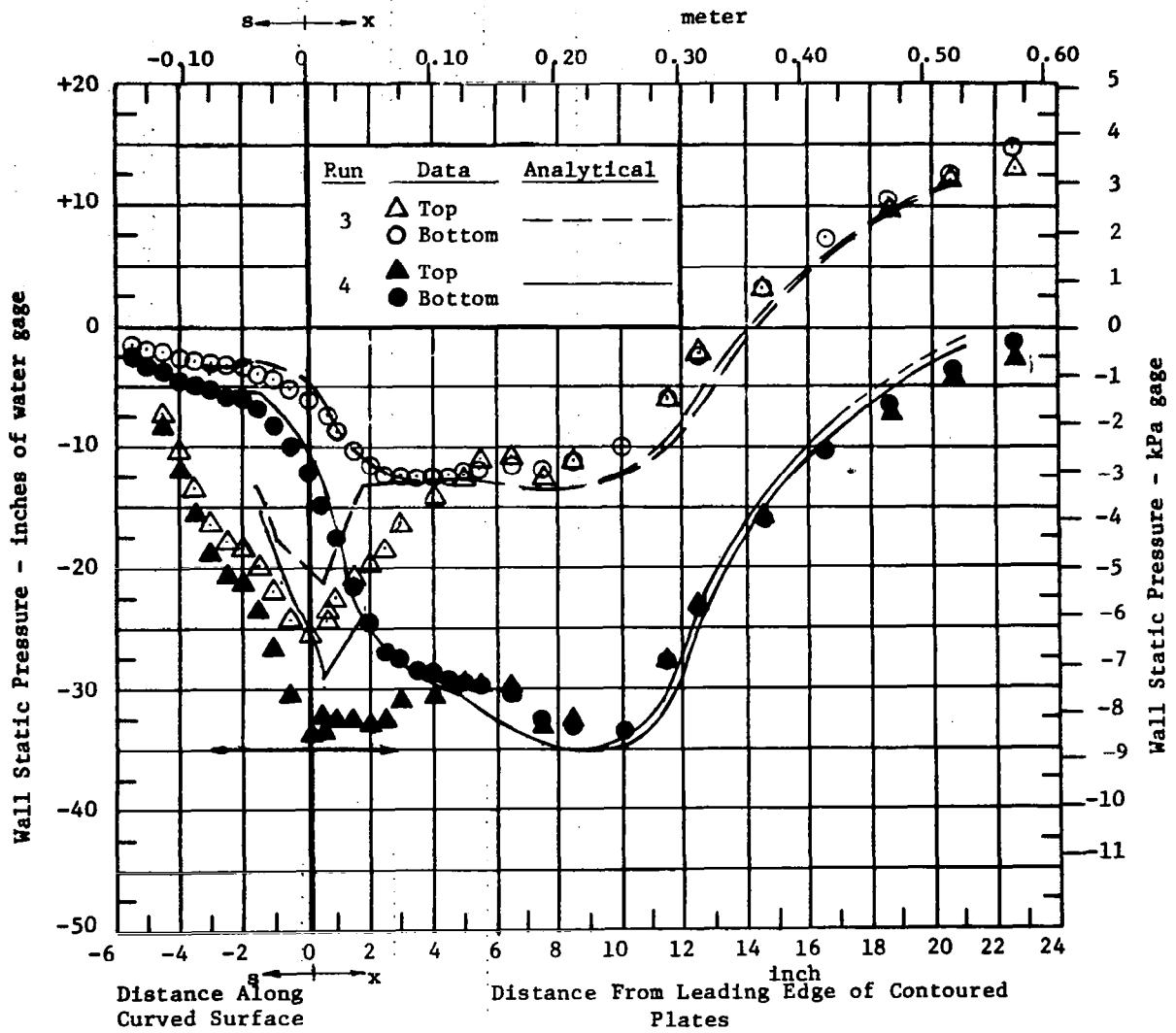


Figure 18 Wall static pressure distributions for Tests 3 and 4. Nozzle positioned at 45° and 0.020 meters (0.80 inches) spacing

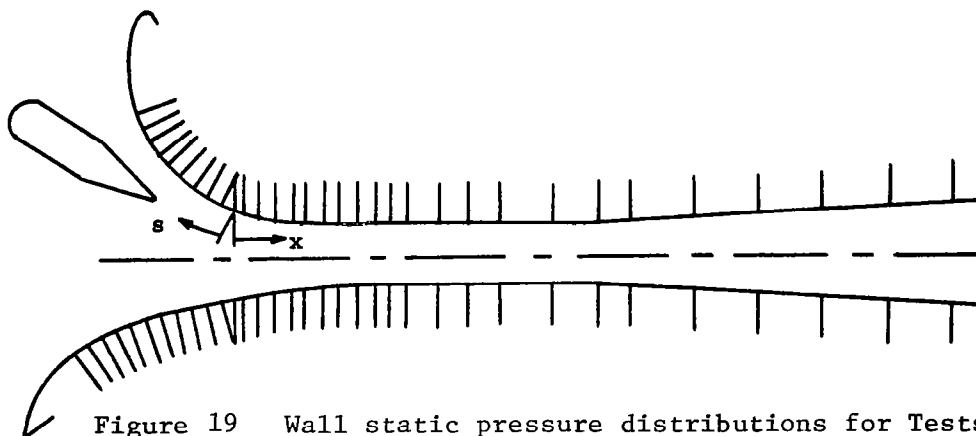
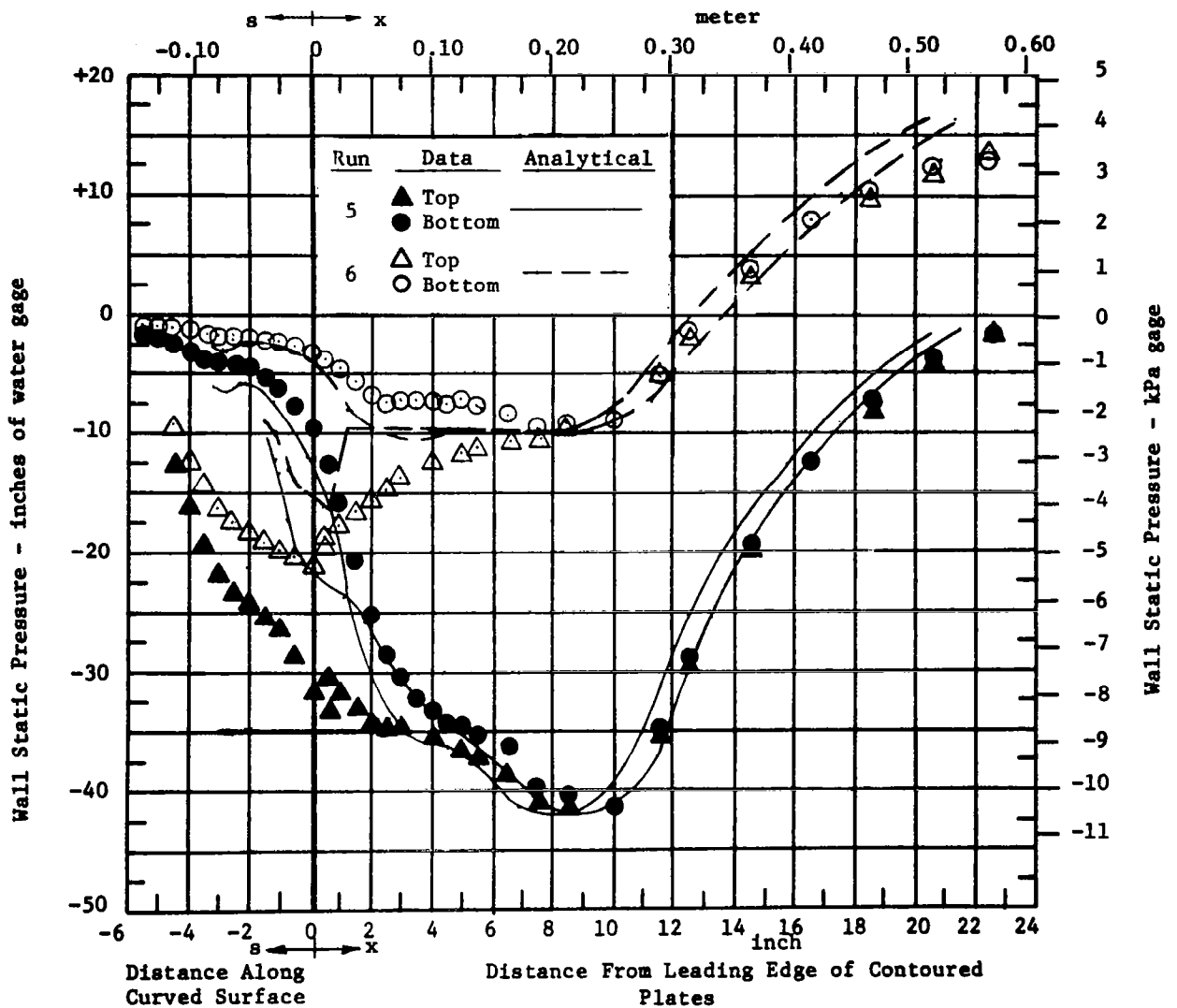


Figure 19 Wall static pressure distributions for Tests 5 and 6. Nozzle positioned at 45° and 0.034 meters (1.32 inches) spacing

Wall Static Pressure - inches of water gage

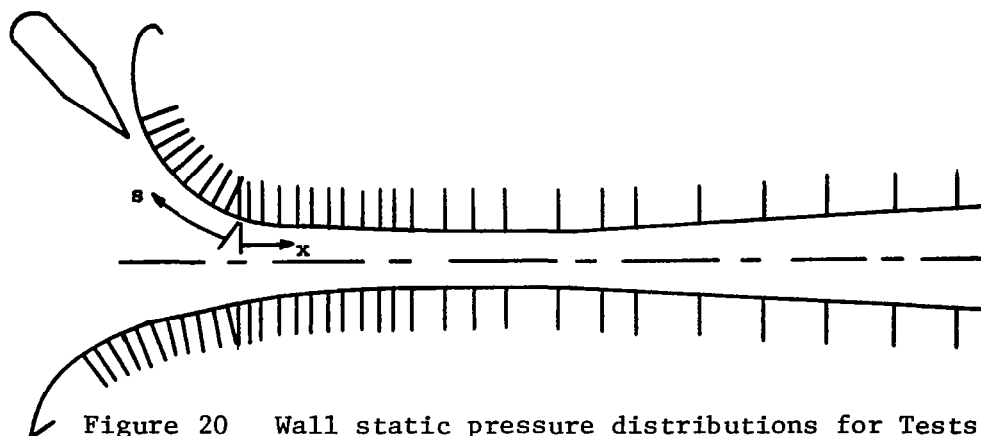
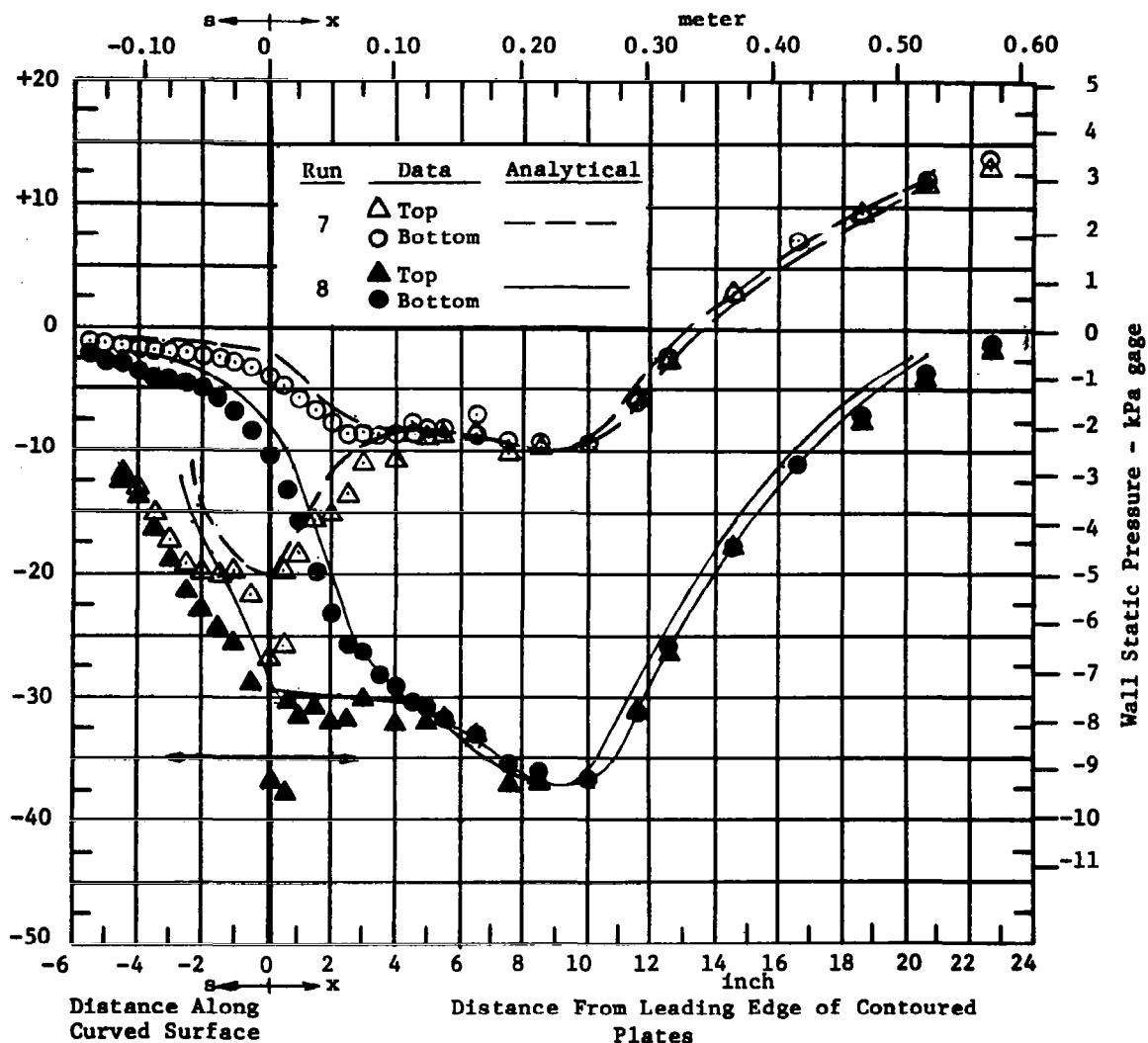


Figure 20 Wall static pressure distributions for Tests 7 and 8. Nozzle positioned at 67.5° and 0.018 meters (0.70 inches) spacing

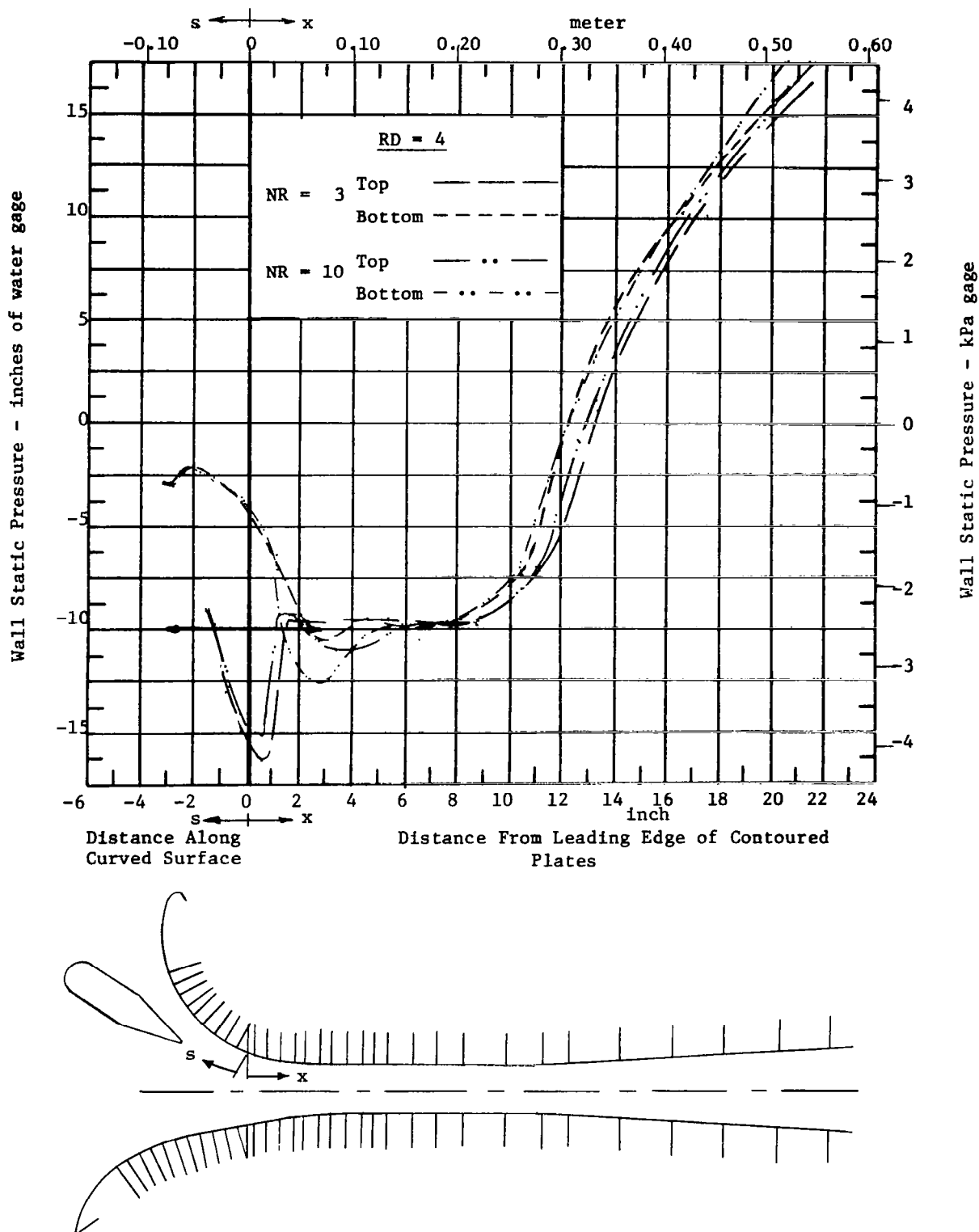


Figure 21 Wall static pressure sensitivity to the Richardson number coefficient (NR) for Test 6



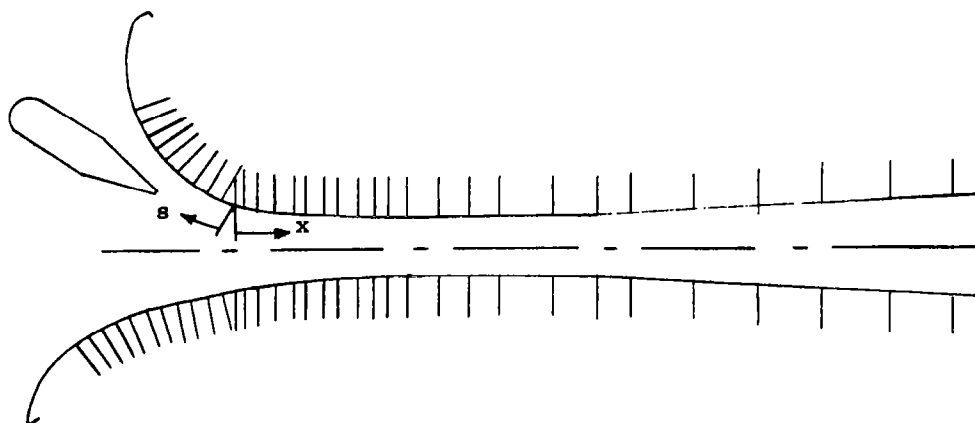
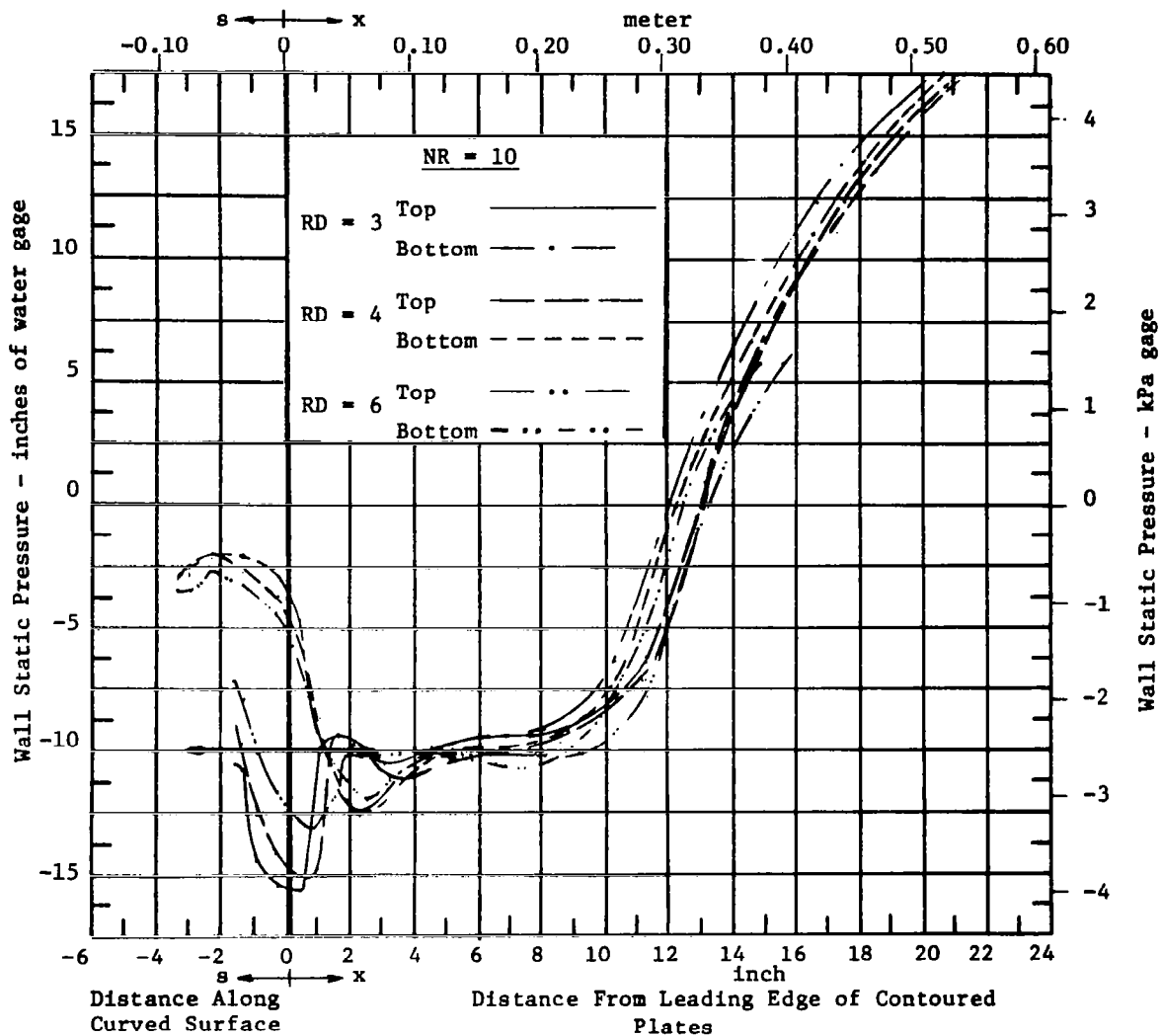


Figure 22 Wall static pressure sensitivity to the rate of streamline curvature decay (RD) for Test 6

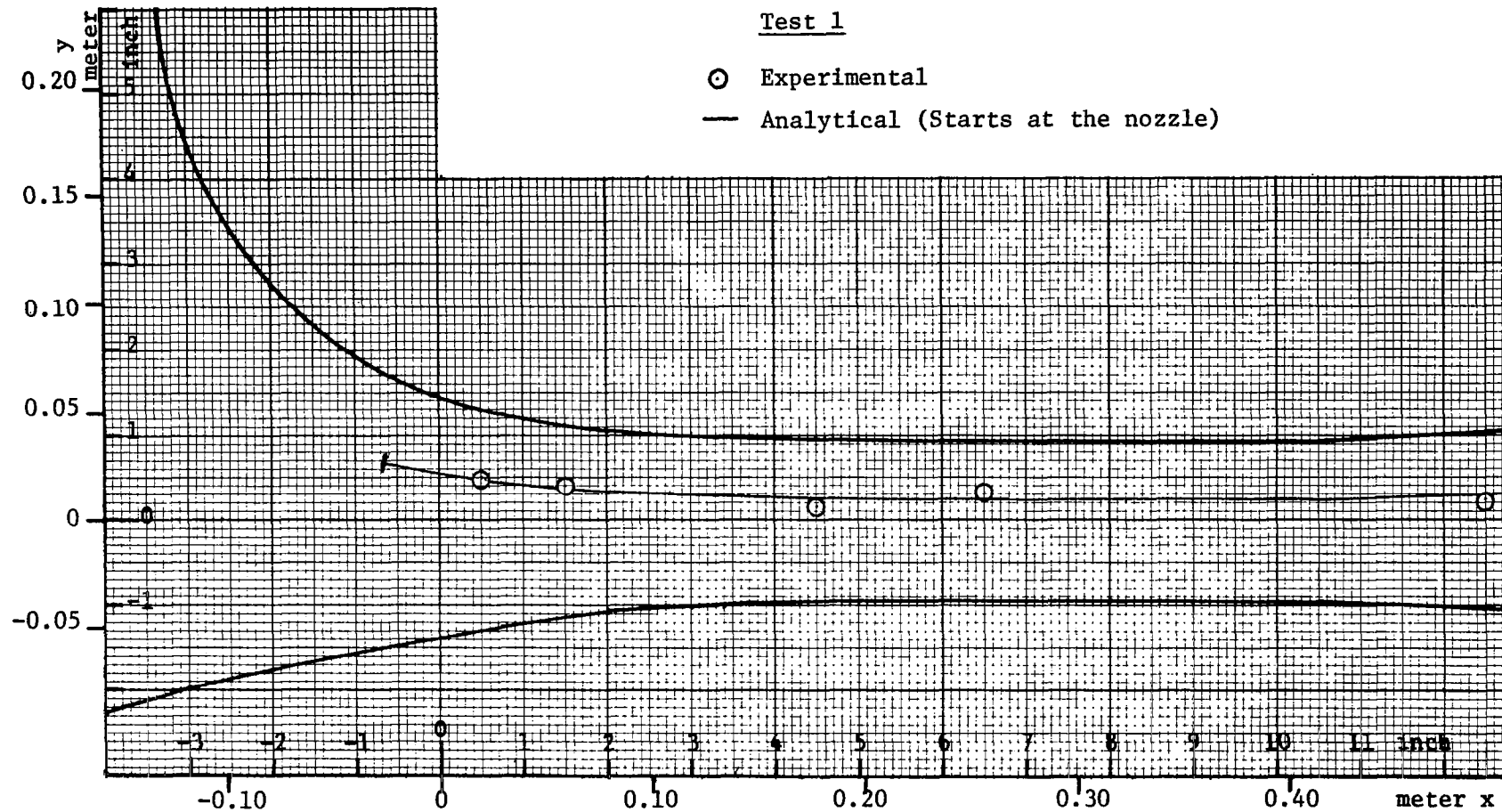


Figure 23

Loci of Maximum Stagnation Pressures ( $P_{o \max}$ ) in the Mixing Section  
 Nozzle Positioned at  $22.5^\circ$  and 0.020 meters (0.80 inches) spacing

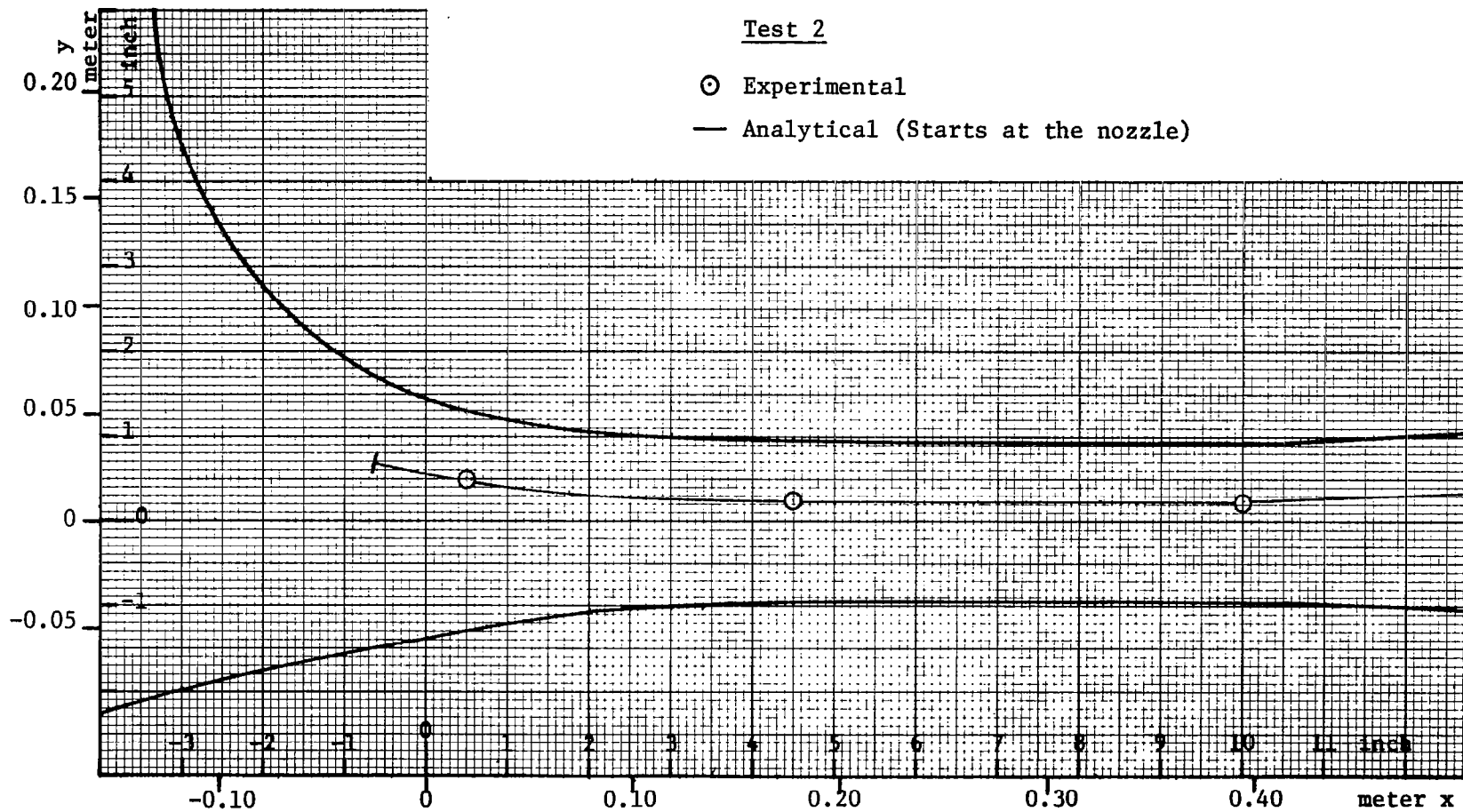


Figure 24

Loci of Maximum Stagnation Pressure ( $P_{o \max}$ ) in the Mixing Section  
Nozzle Positioned at  $22.5^\circ$  and 0.020 meters (0.80 inches) spacing

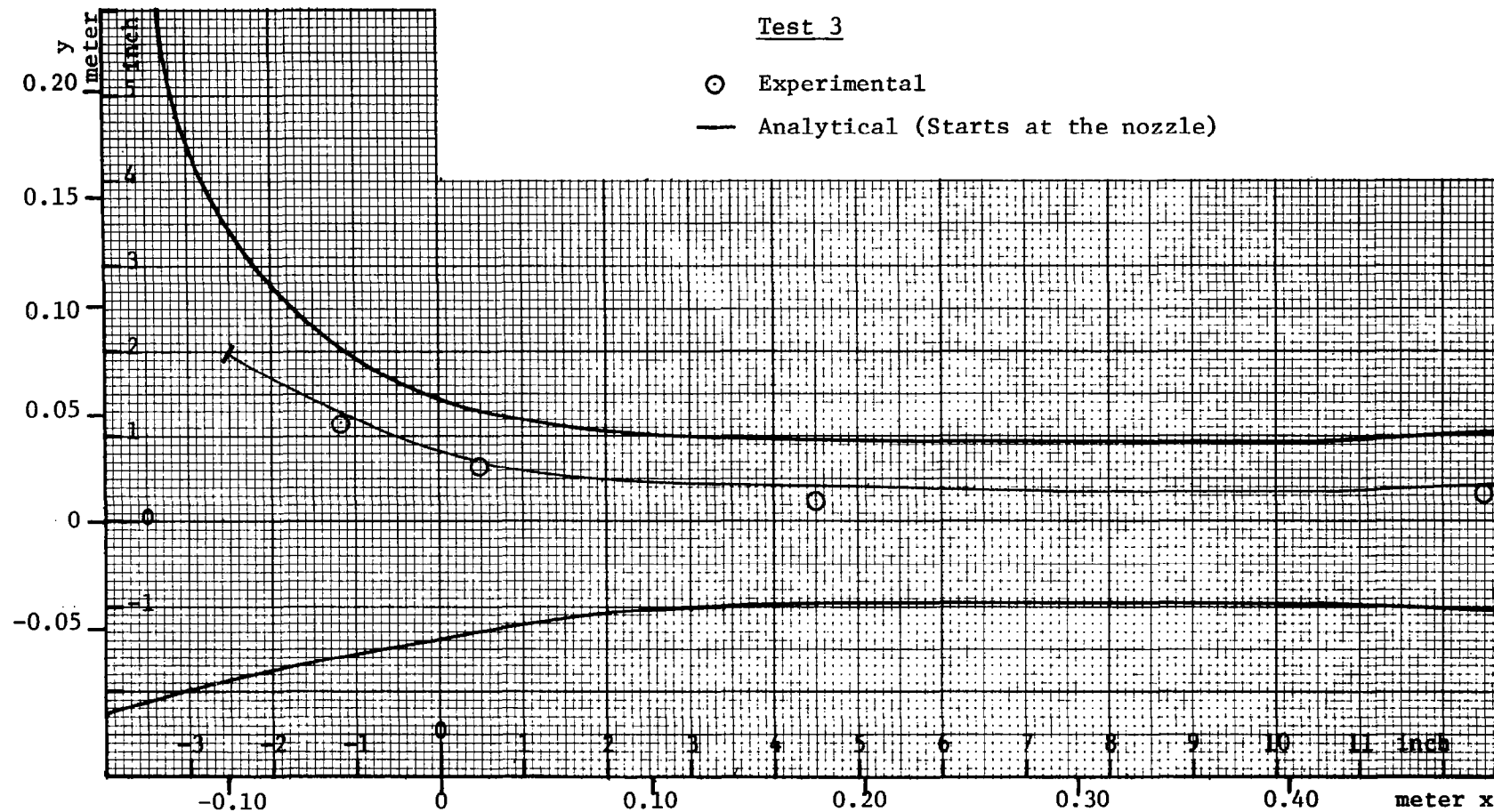
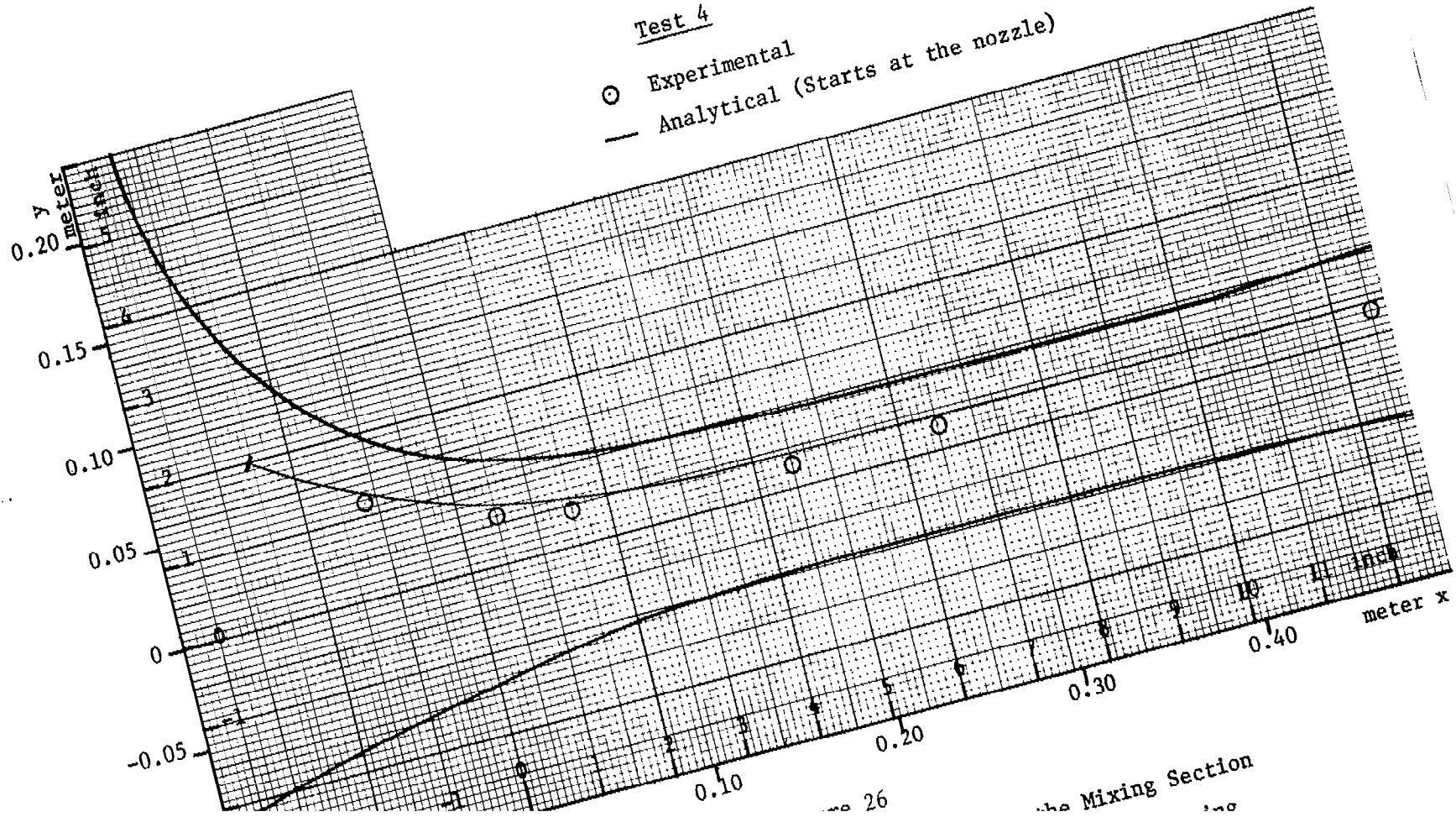


Figure 25

Loci of Maximum Stagnation Pressure ( $P_{o \max}$ ) in the Mixing Section  
 Nozzle Positioned at  $45^\circ$  and 0.020 meters (0.80 inches) spacing

Test 4

○ Experimental  
— Analytical (Starts at the nozzle)



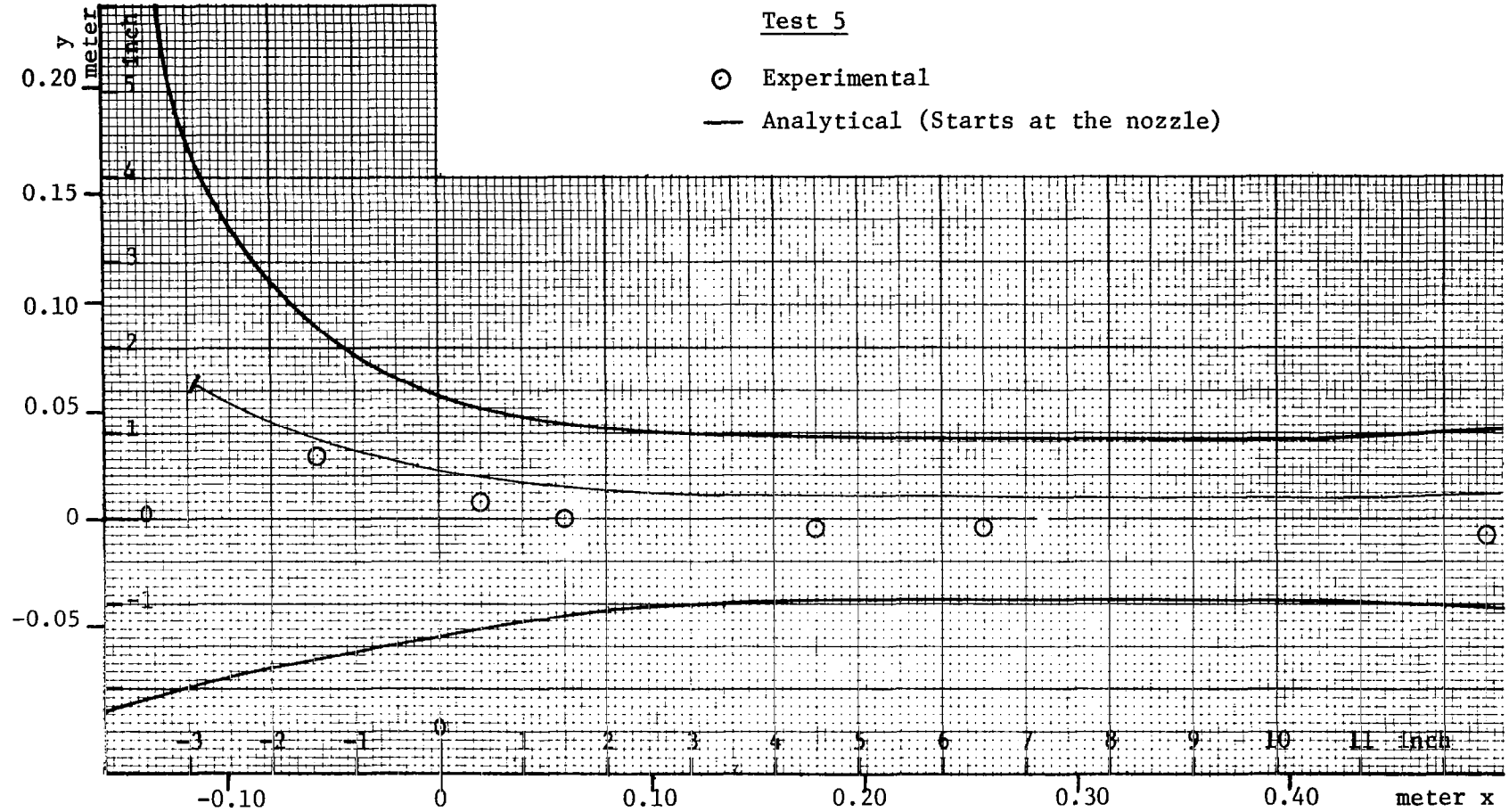


Figure 27

Locii of Maximum Stagnation Pressure ( $P_{o \max}$ ) in the Mixing Section  
 Nozzle Positioned at  $45^\circ$  and 0.034 meters (1.32 inches) spacing

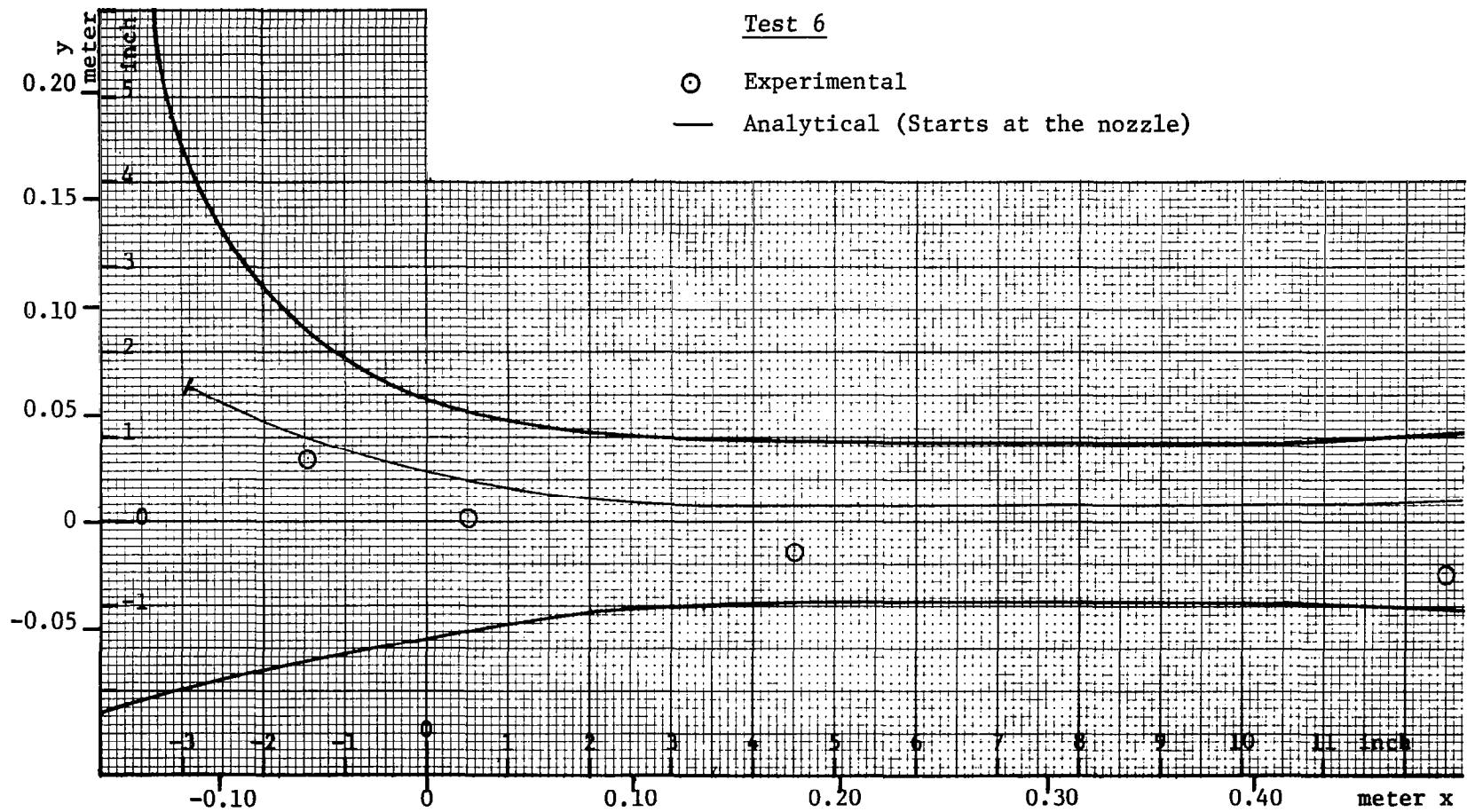


Figure 28

Loci of Maximum Stagnation Pressure ( $P_{o \max}$ ) in Mixing Section  
Nozzle positioned at  $45^\circ$  and 0.034 meters (1.32 inches) spacing

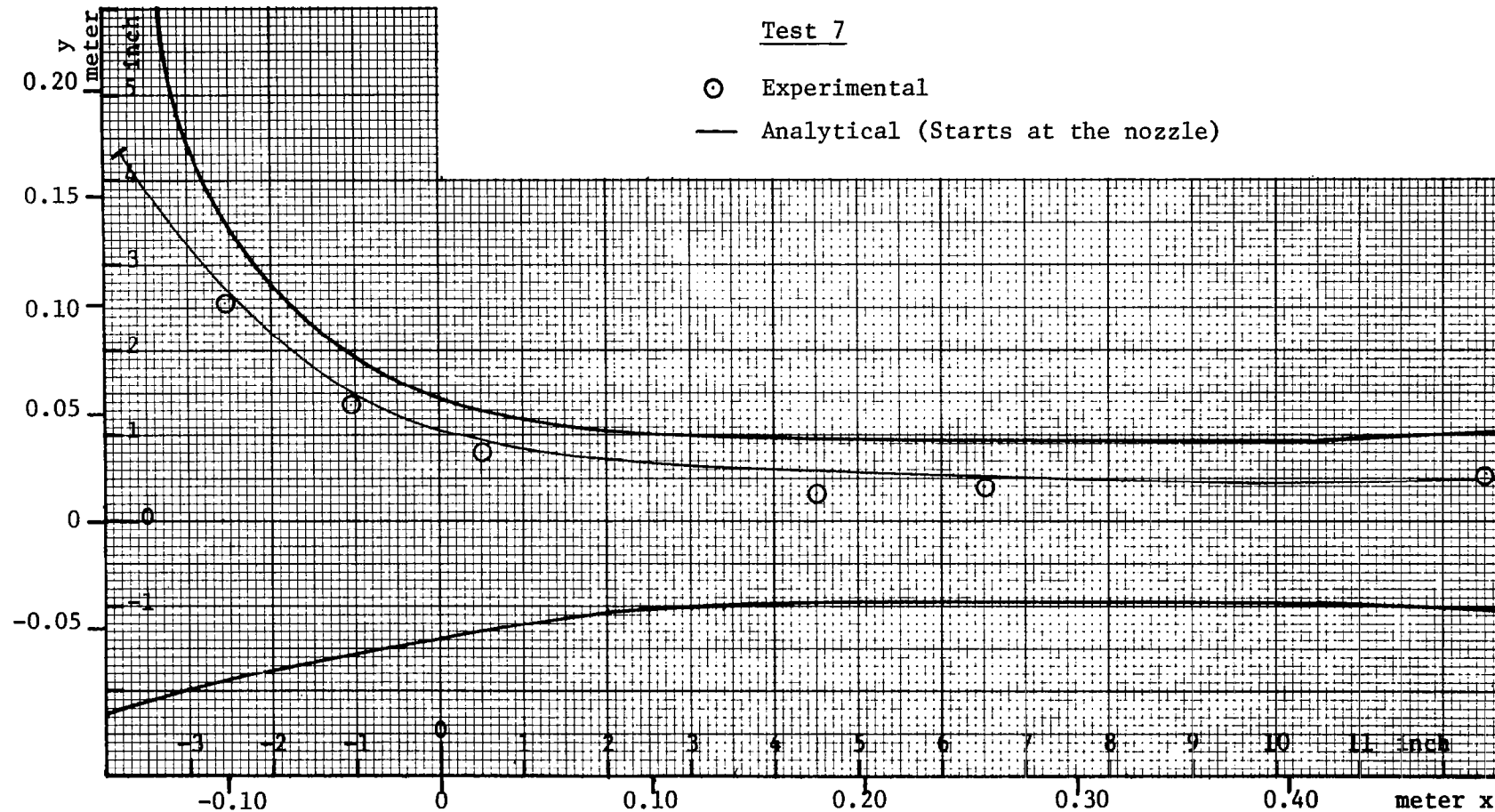


Figure 29

Loci of Maximum Stagnation Pressure ( $P_{o \max}$ ) in Mixing Section  
 Nozzle positioned at  $67.5^\circ$  and 0.018 meters (0.70 inches) spacing



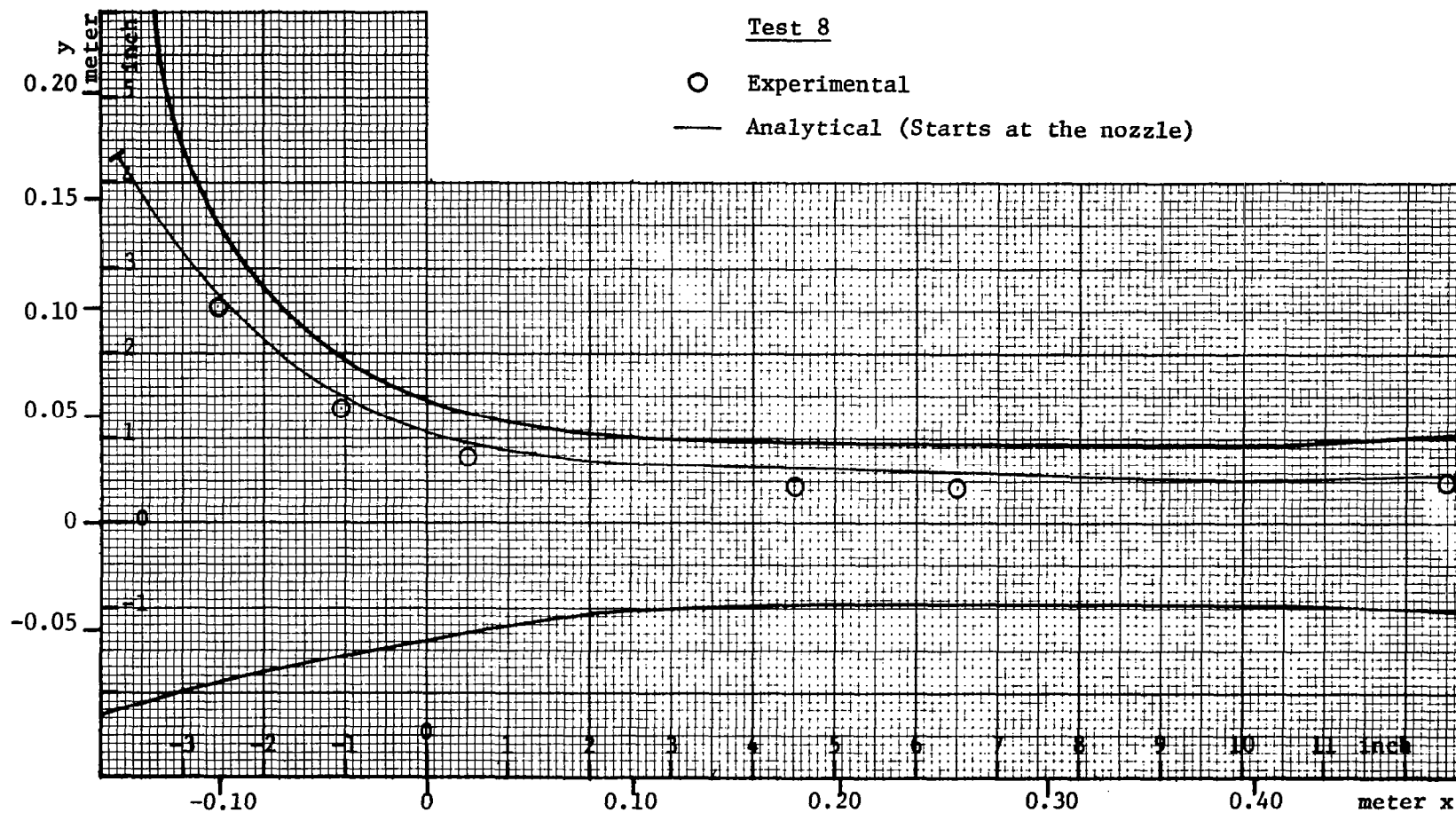


Figure 30

Loci of Maximum Stagnation Pressure ( $P_{o \max}$ ) in Mixing Section  
 Nozzle positioned at  $67.5^\circ$  and 0.018 meters (0.70 inches) spacing

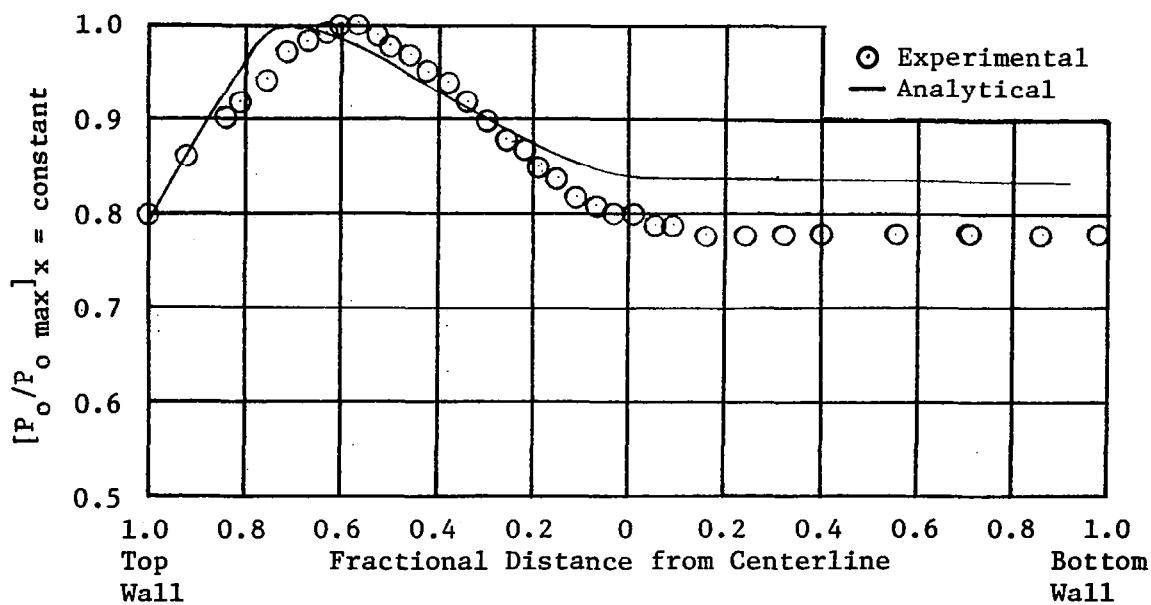


Figure 31 Total Pressure Profiles for Run 7 at  $x = +0.013$  meters (+0.50 inches), Nozzle Positioned at  $67.5^\circ$ , 0.018 meters (0.70 inches) spacing

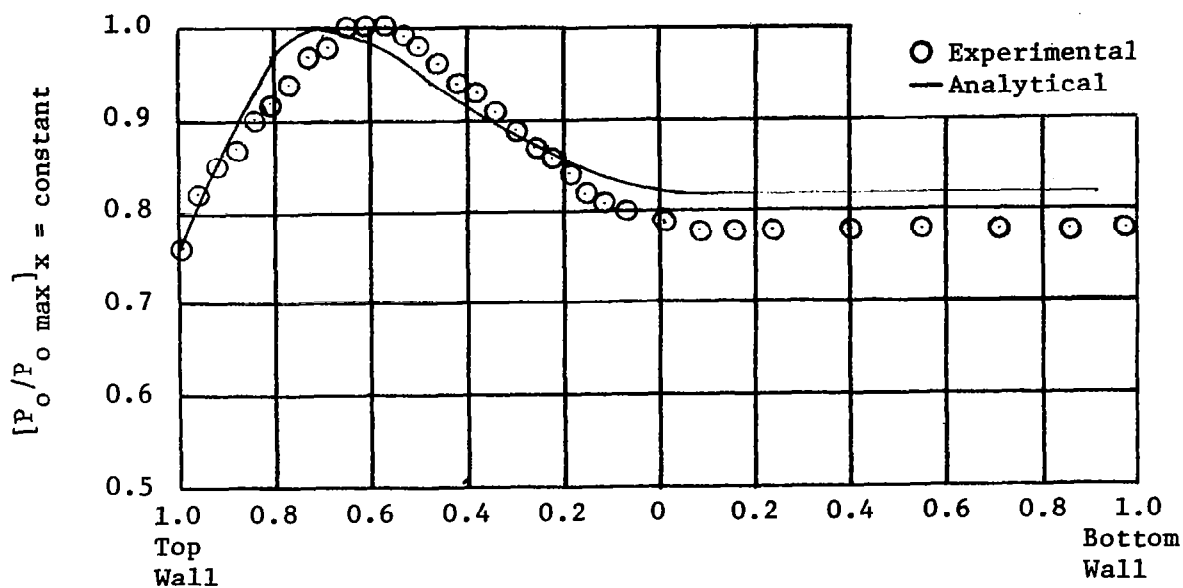


Figure 32 Total Pressure Profile for Run 8 at  $x = +0.013$  meters (+0.50 inches), Nozzle Positioned at  $67.5^\circ$ , 0.018 meters (0.70 inches) spacing

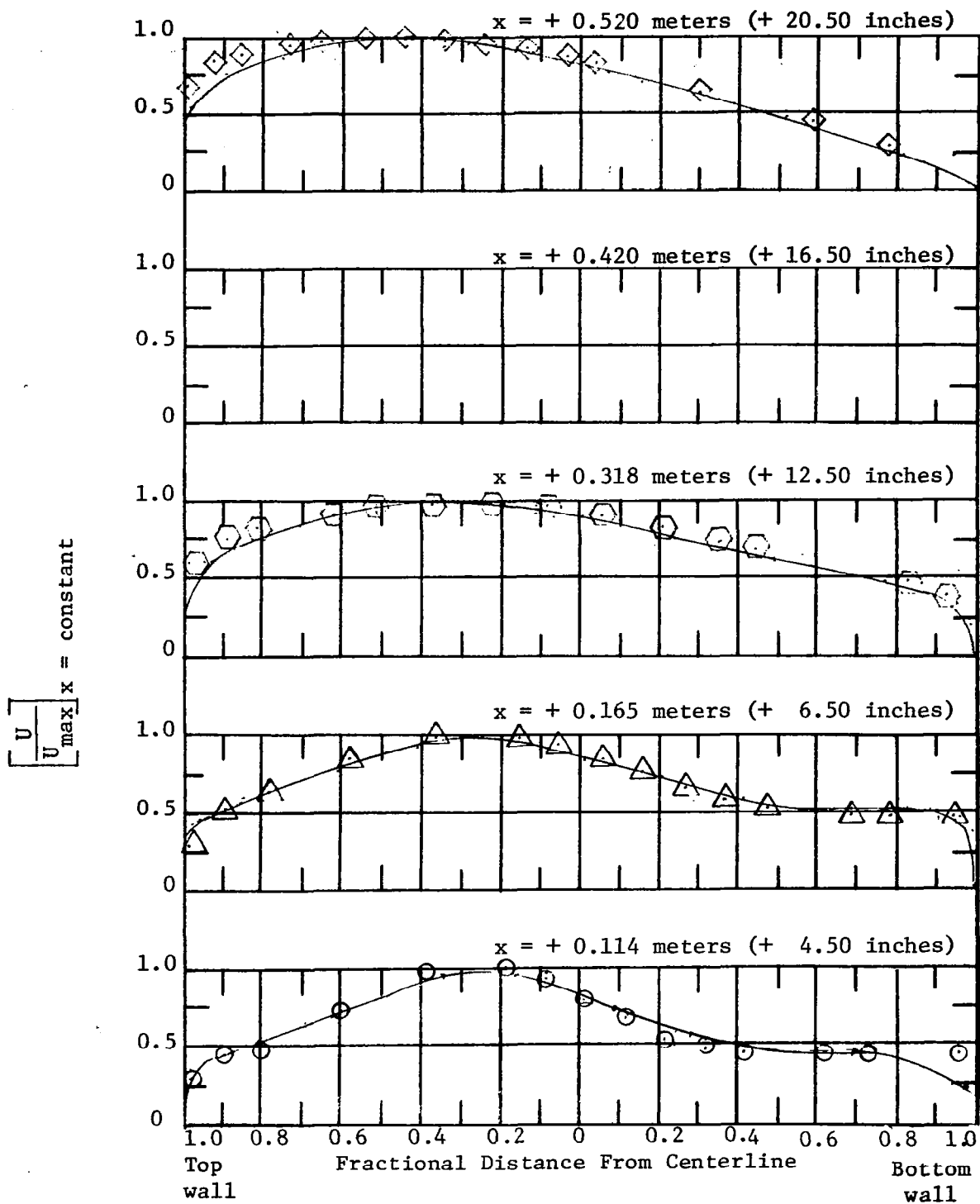


Figure 33 Velocity Profiles for Run 1, Nozzle Positioned at  $22.5^\circ$  and 0.020 meters (0.80 inches) spacing

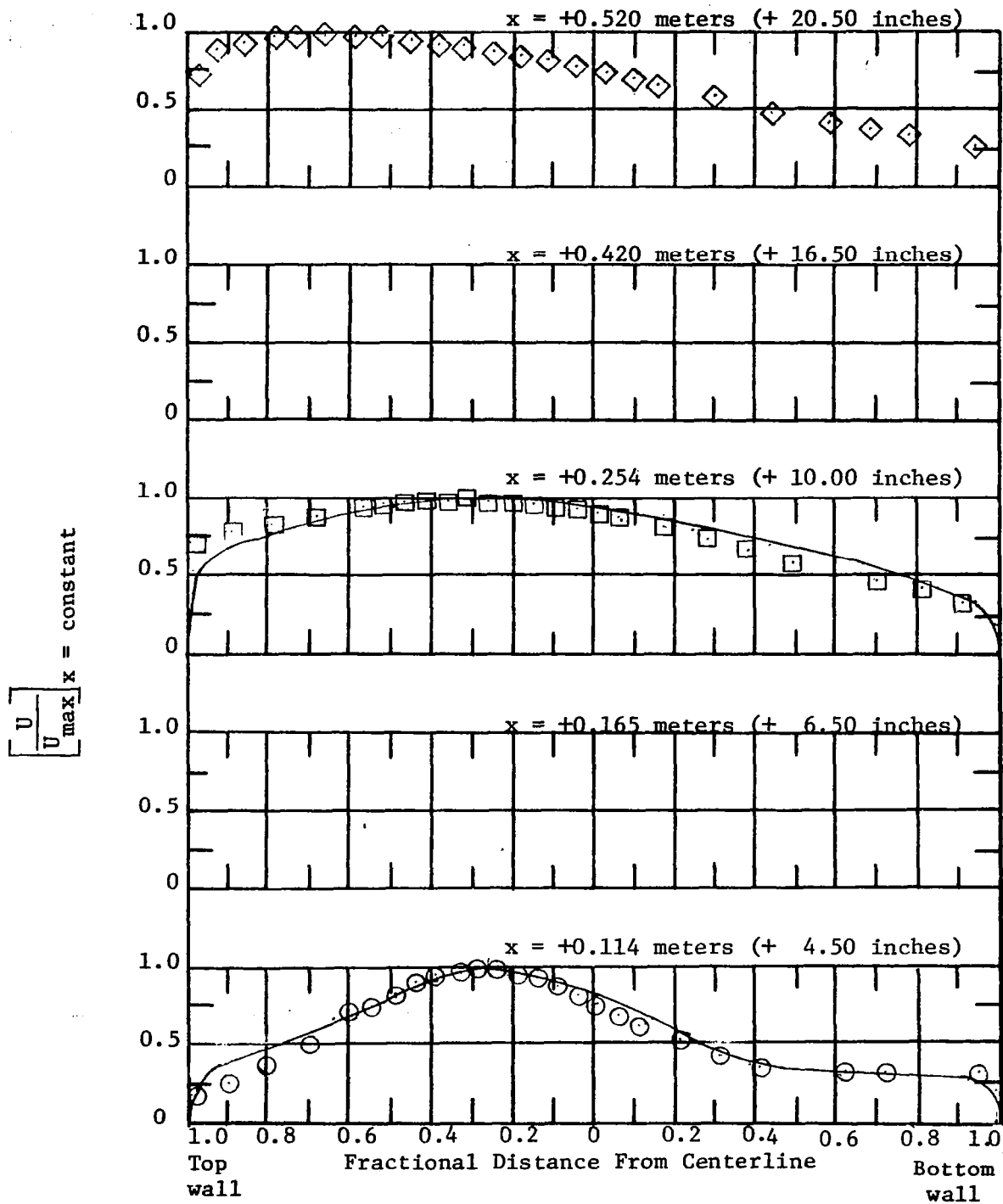


Figure 34 Velocity Profiles for Run 2, Nozzle Positioned at  $22.5^\circ$  and 0.020 meters (0.80 inches) spacing

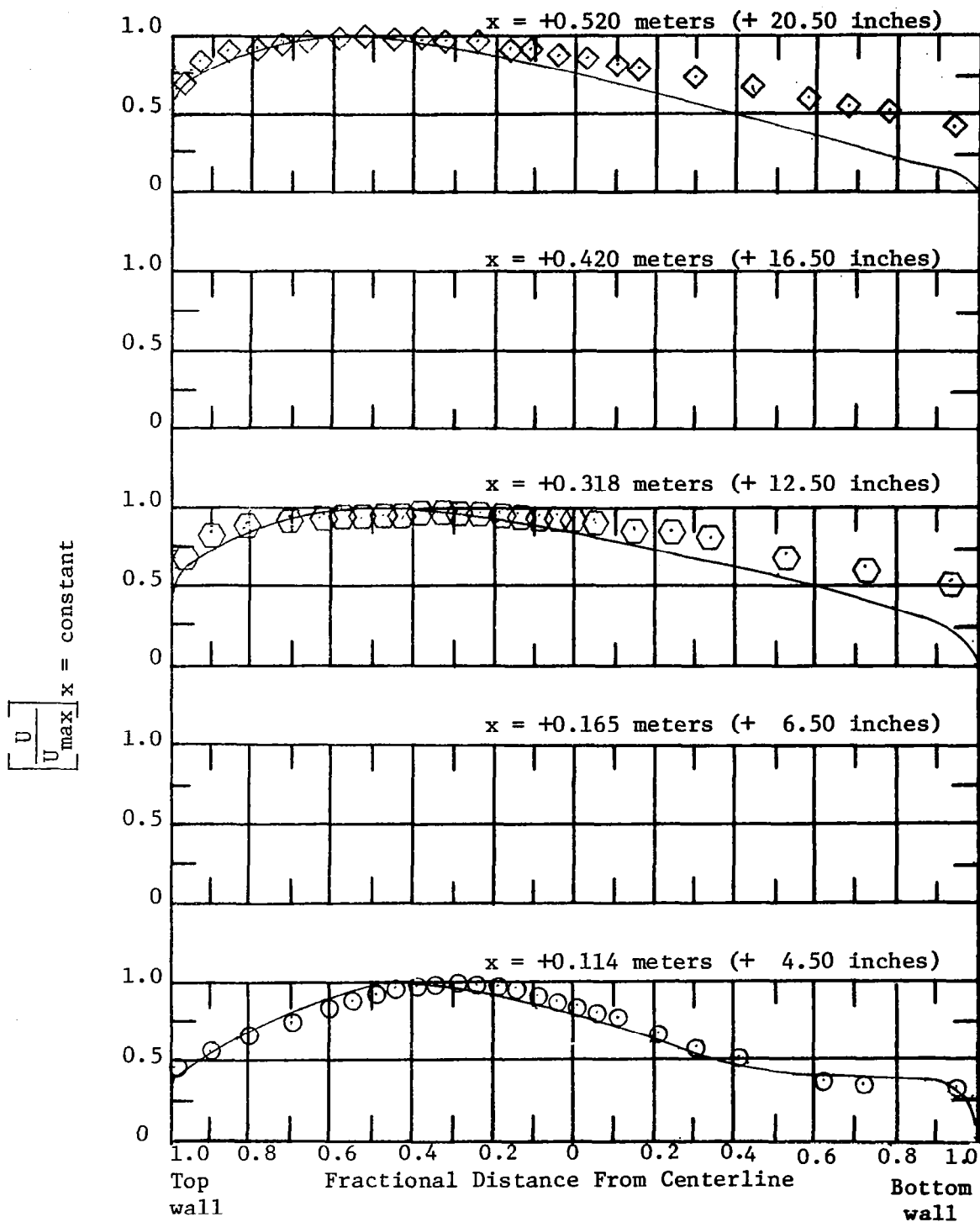


Figure 35 Velocity Profiles for Run 3, Nozzle Positioned at 45° and 0.020 meters (0.80 inches) spacing

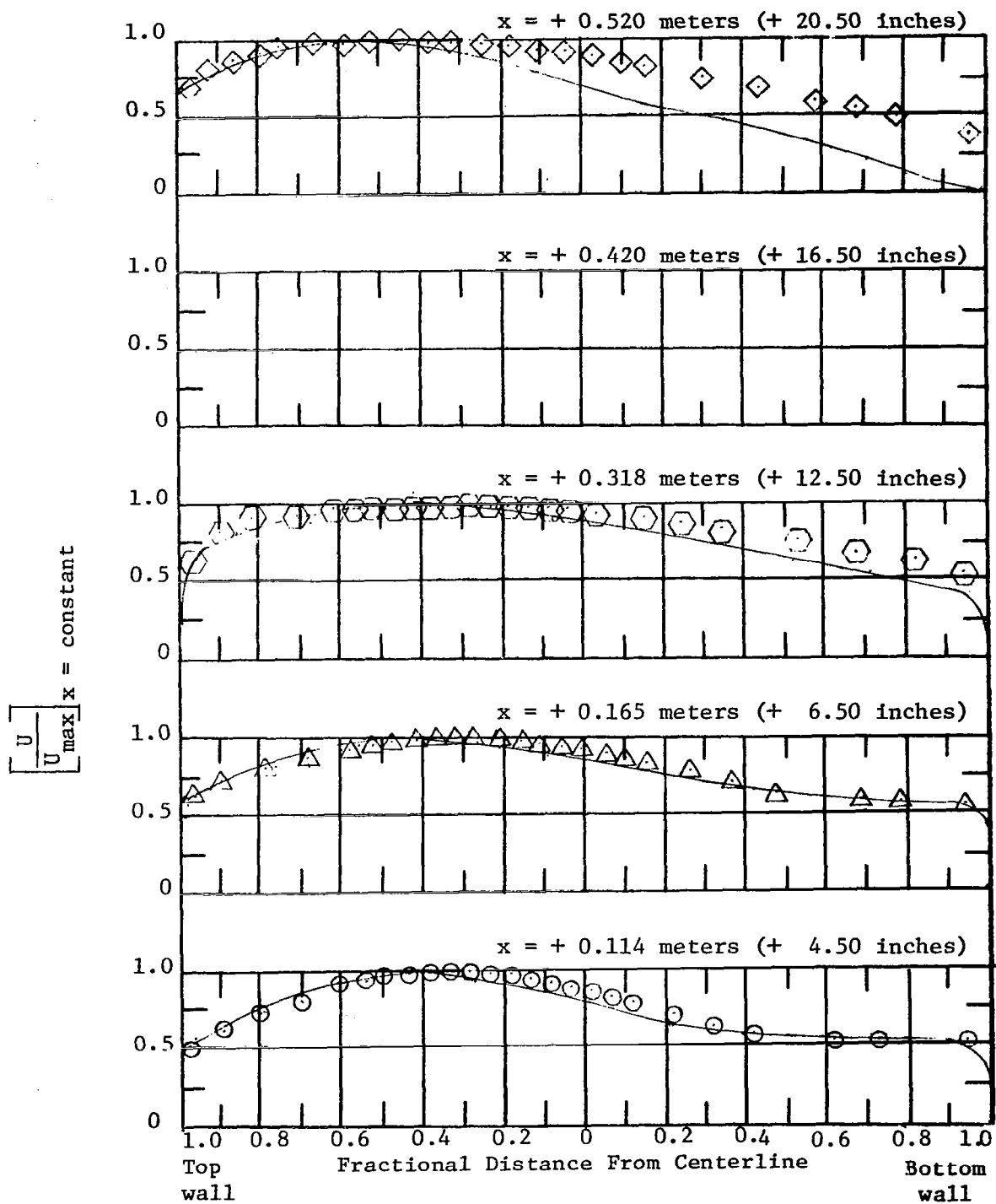


Figure 36 Velocity Profiles for Run 4, Nozzle Positioned at 45° and 0.020 meters (0.80 inches) spacing

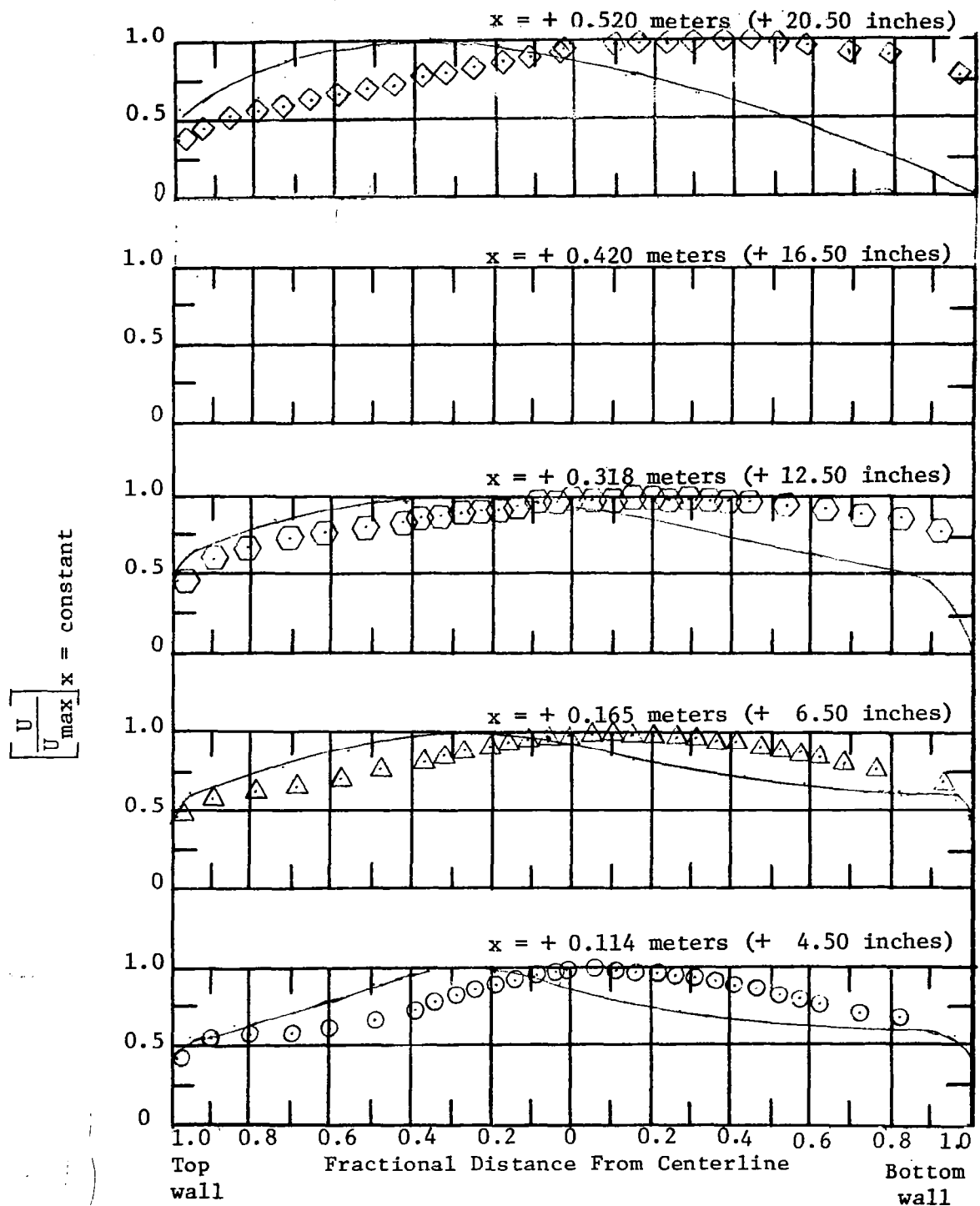


Figure 37 Velocity Profiles for Run 5, Nozzle Positioned at  $45^\circ$  and 0.034 meters (1.32 inches) spacing



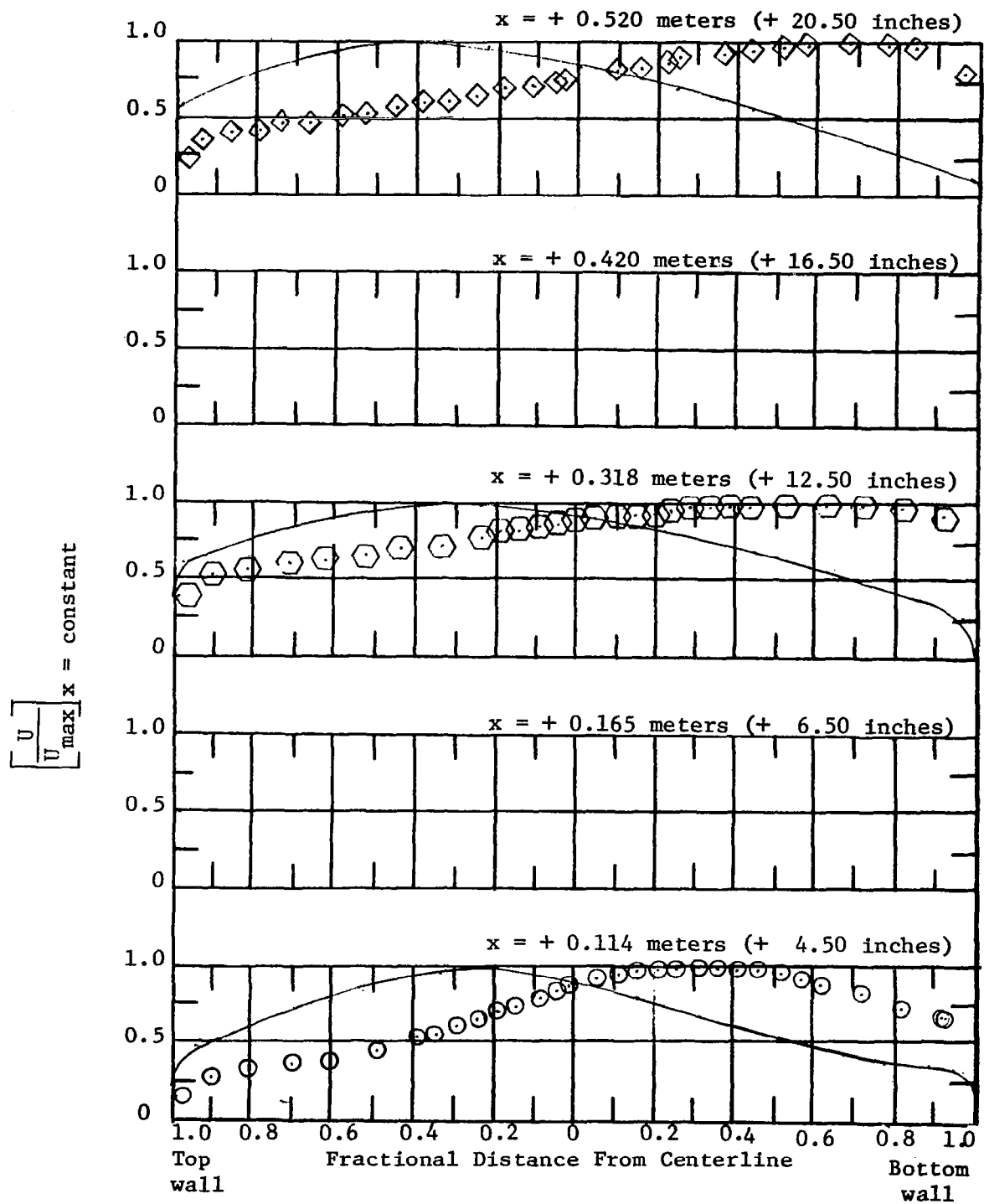


Figure 38 Velocity Profiles for Run 6, Nozzle Positioned at  $45^\circ$  and 0.034 meters (1.32 inches) spacing

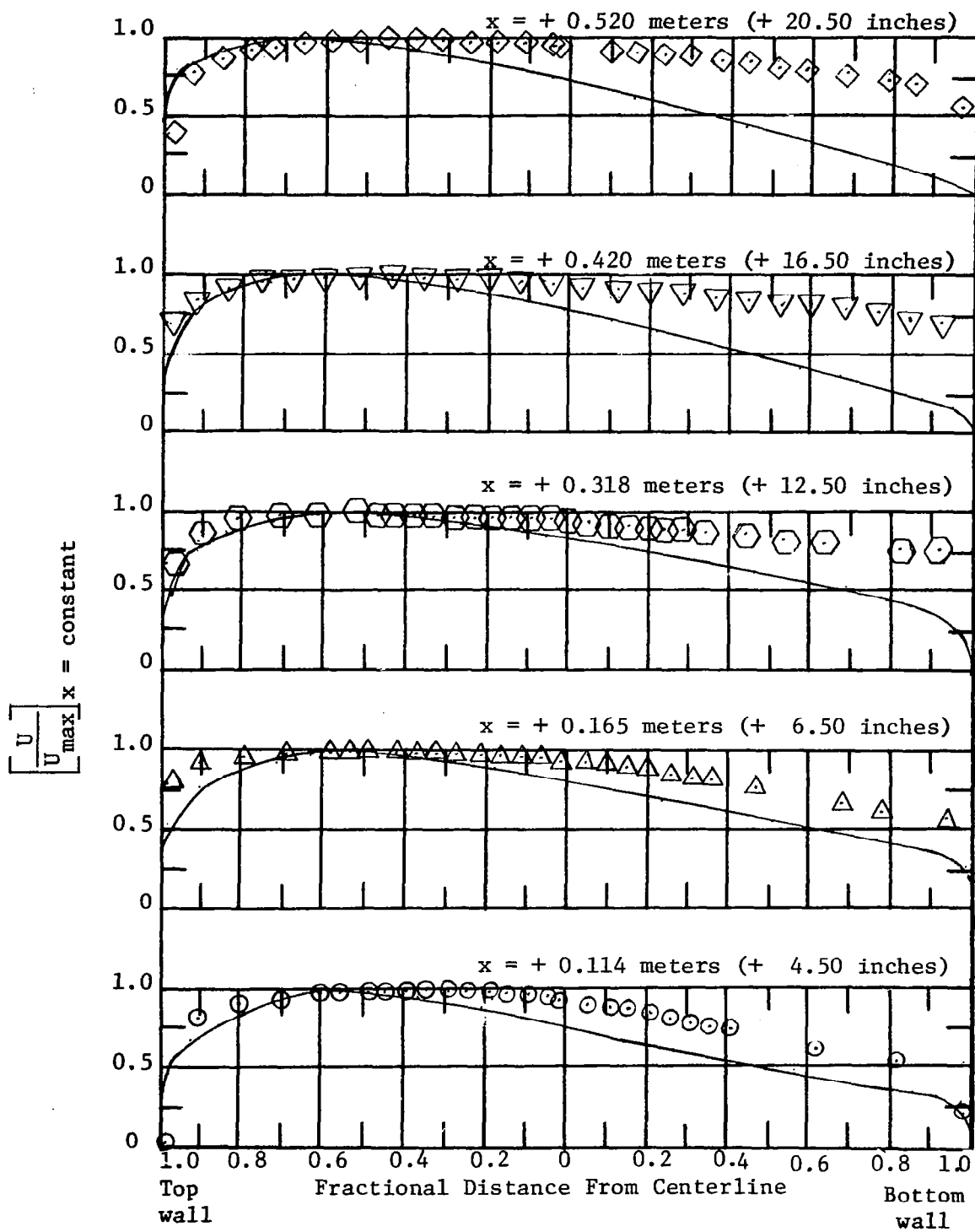


Figure 39 Velocity Profiles for Run 7, Nozzle Positioned at  $67.5^\circ$  and 0.018 meters (0.70 inches) spacing

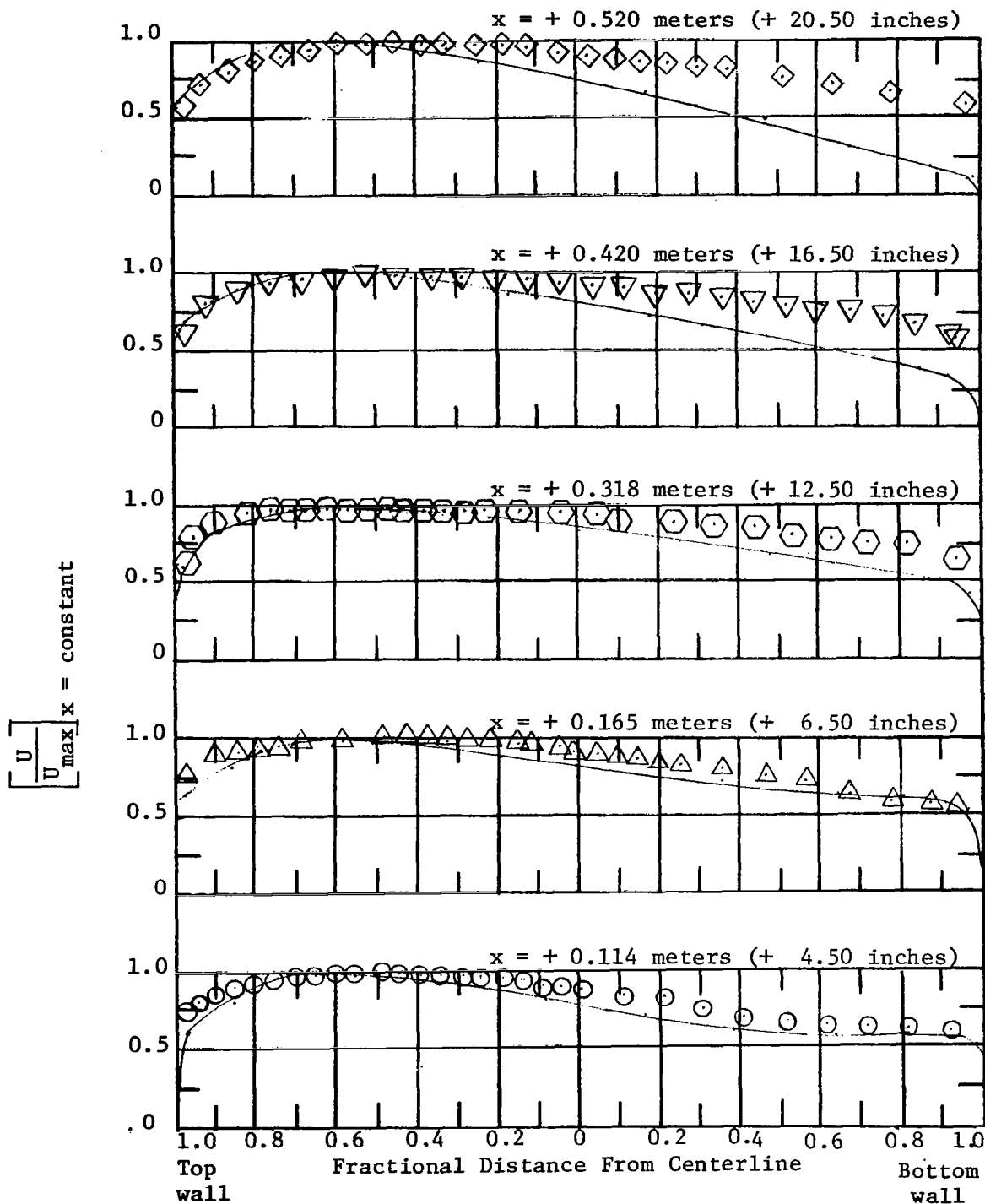


Figure 40 Velocity Profiles for Run 8, Nozzle Positioned at  $67.5^\circ$  and 0.018 meters (0.70 inches) spacing